Search for Technicolor Particles Produced in Association with $W^\pm$ Boson with 1.9 fb$^{-1}$ at CDF

The CDF Collaboration
http://www-cdf.fnal.gov
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We present a search for technicolor particles decaying into $b\bar{b}$, $b\bar{c}$ or $b\bar{u}$ and produced in association with $W$ bosons in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV. The search uses approximately 1.9 fb$^{-1}$ of the dataset accumulated in the CDF detector at the Fermilab Tevatron. We select events matching the $W + 2$-jets signature and require at least one jets to be identified as $b$-quark jets. In the case of exactly one vertex $b$-tagged events, we apply a neural network flavor separator to reject contamination from charm and light quark jets. The number of tagged events and the invariant mass distributions of $W + 2$ jets and dijets are consistent with the Standard Model expectations. We set a 95% confidence level excluded region on the production cross section times branching ratio as a function of the technicolor particle masses.

Preliminary results for Winter 2008 conferences
I. INTRODUCTION

The mechanism of electroweak symmetry breaking in the standard model is still unknown. Two of the most popular mechanisms to induce spontaneous symmetry breaking of the gauge theory, resulting in massive gauge bosons and fermions, are the Higgs mechanism and the dynamics of a new interaction such as technicolor [1, 2]. Both mechanisms predict the existence of new particles which could be produced at the Tevatron in association with a $W$ boson.

Here we present a search for technirho production in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV with subsequent decay to technipions (Fig 1). The signatures we focus on are $p\bar{p} \rightarrow \rho_T^{\pm/0} \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{b}$ and $p\bar{p} \rightarrow \rho_T^0 \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{c}$, $\ell \nu b\bar{u}$. The production cross section of these processes is the order picobarns (Fig. 2) [3].

![Feynman diagram for $p\bar{p} \rightarrow \rho_T^{\pm/0} \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{b}$, $\ell \nu b\bar{c}$ production.](image)

**FIG. 1:** Feynman diagram for $p\bar{p} \rightarrow \rho_T^{\pm/0} \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{b}$, $\ell \nu b\bar{c}$ production.

![Production cross section calculated in Pythia 6.216 with the Technicolor Straw Model (TCSM).](image)

**FIG. 2:** Production cross section calculated in Pythia 6.216 with the Technicolor Straw Model (TCSM). (Left) $p\bar{p} \rightarrow \rho_T^{\pm/0} \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{b}$ production cross section as a function of the technipion mass for various technirho masses. (Right) $p\bar{p} \rightarrow \rho_T^0 \rightarrow W^\pm \pi_T^{\mp} \rightarrow \ell \nu b\bar{c}$, $\ell \nu b\bar{u}$ production cross section as a function of the technipion mass for various technirho masses.

The final state from $W \pi_T$ production is high-$p_T$ lepton, large missing transverse energy ($E_T$) and one or two b-jets.

The previous search at CDF [4] was performed in a dataset with integrated luminosity equivalent to 955 pb$^{-1}$. Other searches at Tevatron have been published using data collected at the D detector [5]. The search for technicolor particle were also performed at LEP [6]. They only excluded technirho mass less than 206.7 GeV due to the LEP beam energy restriction.
II. DATA SAMPLE & EVENT SELECTION

We use the data collected through May 2007, corresponding to an integrated luminosity $1.9 \text{ fb}^{-1}$. The events are collected by the CDF II detector with high-$p_T$ electron or muon triggers [7], which can detect electrons or muons with $E_T$ or $p_T > 18 \text{ GeV}$. We further require the electron or muon that is isolated with $E_T$ or $p_T > 20 \text{ GeV}$ at offline level.

To select the $W + 2\text{jet}$ final state, we require that events have the large missing transverse energy and one or two $b$-jets. Therefore, we require that events have $E_T > 20 \text{ GeV}$ and exactly two jets, where jets are defined using a cone algorithm with radius 0.4. We count jets which have $E_T > 20 \text{ GeV}$ and $|\eta| < 2.0$.

To reduce the background, we require that at least one jet in the event are identified as $b$-jet by the Secondary Vertex Tagging Algorithm. The secondary vertex tagging algorithm identifies $b$-jets by fitting tracks displaced from the primary vertex. This method has been used by other Higgs searches and top analyses at CDF [8, 9]. In addition, we add the Jet Probability Tagging Algorithm that identifies $b$-jets by requiring a low probability that all tracks contained in a jet originated from the primary vertex, based on the track impact parameters [10]. To be considered for double tag category, we require that events have two secondary vertex tags or one secondary vertex tag and one jet probability tag.

To increase the signal acceptance, we also make use of the exactly one $b$-tagged events with secondary vertex tagging algorithm. For the exactly one $b$-tagged events, we apply the neural network $b$-tagging algorithm to improve signal-to-background ratio. This neural network $b$-tagging algorithm is used in previous analysis [4, 9]. With using this algorithm, we can improve the purity of $b$-jets while keeping about 90% signal.

The dijet mass is reconstructed from the 2 jets in selected events. To reconstruct the $W + 2\text{jets}$ invariant mass, we need to determine the $p_z$ of the neutrino from the $W$ boson. After using the $W$ mass constraint to solve for the kinematics of the $\ell \nu$ system, we take a smaller value of the two $p_z$ solutions. (If there is no real solution for $p_z$, we take the real part of the complex solution.)

III. SIGNAL ACCEPTANCE

The signal acceptance is calculated using technicolor events generated with the PYTHIA program [11]. PYTHIA version 6.216 implements the Technicolor Straw Man Model [3] in leading-order calculations. We set the mass parameters as $M_V = M_A = 200 \text{ GeV}/c^2$, the charge of up-type technifermion as $Q_U = 1$ and the mixing angle between isotriplet technipion interaction and mass eigenstates as $\sin \chi = 1/3$ in this model. For this study, we focus on the mass region as:

$$m(W) + m(\pi_T) < m(\rho_T) < 2 \times m(\pi_T) \quad (1)$$

$$180 \text{ GeV}/c^2 < m(\rho_T) < 250 \text{ GeV}/c^2$$

$$95 \text{ GeV}/c^2 < m(\pi_T) < 165 \text{ GeV}/c^2$$

because of the kinematical threshold of $W\pi_T$ production and pair $\pi_T$ production.

Figure 3-4 show the acceptance as a function of $m(\pi_T)$ for each $m(\rho_T)$ for neutral and charged technipi channel. The acceptance includes various calibration scale factors quantifying the difference between simulation and data, and it includes the efficiency of the high $p_T$-lepton triggers. No $K$-factor is applied to the leading-order calculation of the acceptance and cross section.
FIG. 3: Technicolor signal acceptance for neutral technipi channel as a function of technipion mass $m(\pi_T)$ for each technirho mass $m(\rho_T)$. The three different selections – exactly one secondary vertex tag with NN $b$-tag, double secondary vertex tag or one secondary vertex tag plus one jet probability tag – are shown in magenta, red and blue, respectively.
FIG. 4: Technicolor signal acceptance for charged technipi channel as a function of technipion mass $m(\pi_T)$ for each technirho mass $m(\rho_T)$. The three different selections—exactly 1 one secondary vertex tag with NN $b$-tag, double secondary vertex tag or one secondary vertex tag plus one jet probability tag—are shown in magenta, red and blue, respectively.
Systematic uncertainties from the $b$-tagging efficiency, initial and final state radiation effects (ISR/FSR), parton distribution function (PDF) and the jet energy scale (JES) are also considered. The summarized systematic uncertainties are shown in Table I-II.

### IV. BACKGROUND ESTIMATES

This analysis builds on the method of background estimation detailed in Ref. [8]. In particular, the contributions from the following individual backgrounds are calculated: falsely $b$-tagged events (mistag), $W$ production with heavy flavor quarks ($W + bb$, $W + cc$, $W + c$), QCD events with false $W$ signatures, top quark pair or single production and diboson production ($WW$, $WZ$, $ZZ$).

We estimate the mistag events by using the mistag probability that measured from the inclusive jets sample. Such mistag rate are obtained using negative tags, which are the tags that appear to travel back toward the primary vertex. The mistag rate derived from negative tags is due to tracking resolution limitations, but they provide a reasonable estimate of the number of false positive tags after a correction for material interactions and long-lived light flavor particles.

The number of events from $W +$ heavy flavor is calculated using information from both data and Monte Carlo samples. We calculate the fraction of $W$ events with associated heavy flavor production in the ALPGEN Monte Carlo program interfaced with the PYTHIA parton shower code [11, 12]. This fraction and the tagging efficiency for such events are applied to the number of events in the original $W$+jets sample after correcting for the $t\bar{t}$ and electroweak contributions.

We constrain the number of QCD events with false $W$ signatures by assuming the lepton isolation is independent of $E_T$ and measuring the ratio of isolated to non-isolated leptons in a $E_T$ sideband region. The result in the tagged sample can be calculated in two ways: by applying the method directly to the tagged sample, or by estimating the number of non-$W$ QCD events in the pretag sample and applying an average $b$-tagging rate.

The summary of the background contributions to the double secondary vertex tagged events, the one secondary vertex plus one jet probability tagged events and the one secondary vertex tagged with neural network $b$-tag events are shown in Table III-V. Because the expected number of technicolor signal events is small in the 1-, 3- and 4-jet bins, the reasonable agreement between predicted backgrounds and observed data gives us confidence in our overall background estimate.
\begin{table}[h]
\centering
\begin{tabular}{lccc}
\hline
\hline
Njet & 2jet & 3jet & \geq 4jet \\
\hline
Pretag Events & 32242 & 5496 & 1494 \\
Mistag & 3.88\pm0.35 & 2.41\pm0.24 & 1.62\pm0.14 \\
$Wb\bar{b}$ & 37.93\pm16.92 & 14.05\pm5.49 & 7.39\pm2.93 \\
$Wc\bar{c}$ & 2.88\pm1.25 & 1.52\pm0.61 & 1.15\pm0.47 \\
$t\bar{t}(6.7\text{pb})$ & 19.05\pm2.92 & 54.67\pm8.38 & 94.93\pm14.56 \\
Single top(s-ch) & 6.90\pm1.00 & 2.28\pm0.33 & 0.61\pm0.088 \\
Single top(t-ch) & 1.60\pm0.23 & 1.43\pm0.21 & 0.50\pm0.07 \\
WW & 0.17\pm0.02 & 0.15\pm0.02 & 0.16\pm0.02 \\
WZ & 2.41\pm0.26 & 0.68\pm0.07 & 0.16\pm0.02 \\
ZZ & 0.06\pm0.01 & 0.06\pm0.01 & 0.02\pm0.001 \\
$Z^- > \tau\tau$ & 0.25\pm0.04 & 0.19\pm0.03 & 0.06\pm0.01 \\
nonW QCD & 5.50\pm1.00 & 2.56\pm0.48 & 1.02\pm0.22 \\
\hline
Total Bkg & 80.62\pm18.75 & 79.99\pm10.92 & 107.63\pm15.15 \\
\hline
\end{tabular}
\caption{Background summary table for double secondary vertex tag category.}
\end{table}

\begin{table}[h]
\centering
\begin{tabular}{lccc}
\hline
\hline
Njet & 2jet & 3jet & \geq 4jet \\
\hline
Pretag Events & 32242 & 5496 & 1494 \\
Mistag & 11.73\pm0.92 & 8.11\pm0.64 & 8.39\pm0.58 \\
$Wb\bar{b}$ & 31.15\pm14.03 & 11.47\pm4.55 & 6.55\pm2.63 \\
$Wc\bar{c}$ & 7.87\pm3.43 & 4.38\pm1.76 & 3.09\pm1.27 \\
$t\bar{t}(6.7\text{pb})$ & 15.56\pm2.39 & 47.48\pm7.28 & 79.81\pm12.24 \\
Single top(s-ch) & 5.14\pm0.75 & 1.90\pm0.27 & 0.53\pm0.07 \\
Single top(t-ch) & 1.87\pm0.27 & 1.49\pm0.22 & 0.44\pm0.06 \\
WW & 0.93\pm0.11 & 0.63\pm0.08 & 0.47\pm0.06 \\
WZ & 1.84\pm0.20 & 0.59\pm0.06 & 0.19\pm0.02 \\
ZZ & 0.08\pm0.01 & 0.04\pm0.003 & 0.02\pm0.002 \\
$Z^- > \tau\tau$ & 1.29\pm0.20 & 0.53\pm0.08 & 0.20\pm0.03 \\
nonW QCD & 9.55\pm1.73 & 4.87\pm0.93 & 1.80\pm0.40 \\
\hline
Total Bkg & 86.99\pm17.99 & 81.46\pm10.22 & 101.49\pm13.08 \\
\hline
\end{tabular}
\caption{Background summary table for one secondary vertex tag + one jet probability tag category.}
\end{table}
V. RESULTS

We perform a direct search for a resonant mass peak in the reconstructed $W + 2$-jet and dijet invariant mass distributions from the single-tagged and double-tagged events. In fact, we consider the dijet mass distribution and the $Q$-value defined as $Q = m(T^+) - m(T^-) - m(W)$. The observed data spectra for these variables are consistent with the background estimate, as shown in Figs. 5-7.

![Distribution of dijet invariant masses for the predicted and observed data. Results are shown for the double secondary vertex tag (left), the one secondary vertex tag + one jet probability tag (middle) and the one secondary vertex tag with NN tag (right).](image-url)
Since there is no significant excess of events in the data compared to the predicted background, we set the 95% C.L. excluded region on technicolor production as a function of the technicolor particles mass. A 2-dimensional binned maximum likelihood technique is used on the 2-dimensional distribution of dijet invariant mass vs $Q$-value by constraining the number of background events within the uncertainties. Figs. 8-10 show the 2-dimensional distribution of data, backgrounds and signal for each tag category.
FIG. 8: Dijet mass vs Q value 2-dimensional distribution of signal MC (left, $m(\rho_T) = 200$GeV/$c^2$, $m(\tau_T) = 115$GeV/$c^2$), data (middle) and backgrounds (right) for the double secondary vertex tag category.

FIG. 9: Dijet mass vs Q value 2-dimensional distribution of signal MC (left, $m(\rho_T) = 200$GeV/$c^2$, $m(\tau_T) = 115$GeV/$c^2$), data (middle) and backgrounds (right) for the one secondary vertex tag + one jet probability tag category.

FIG. 10: Dijet mass vs Q value 2-dimensional distribution of signal MC (left, $m(\rho_T) = 200$GeV/$c^2$, $m(\tau_T) = 115$GeV/$c^2$), data (middle) and backgrounds (right) for the one secondary vertex tag with NN tag category.
The final expected and observed excluded region at 95% C.L. are shown in Fig. 11. A region of $m(\rho_T) = 180 - 250\text{GeV}$ are excluded at 95% C.L. based on the Technicolor Straw Man Model, except the region which are near the $W \pi_T$ production threshold with $m(\rho_T) \geq 220\text{GeV}$ and $m(\pi_T) \geq 125\text{GeV}$.

![CDF Run II Preliminary (1.9 fb⁻¹)](image)

**FIG. 11:** The expected and observed excluded region at 95% C.L. as a function of technirho mass and technipi mass. Kinematical threshold of $W\pi_T$ and $\rho_T \rightarrow \pi_T\pi_T$ are shown in the figure.
<table>
<thead>
<tr>
<th>$m(\rho_T)$</th>
<th>$m(\pi_T)$ exclude region at 95% C.L.</th>
</tr>
</thead>
<tbody>
<tr>
<td>180 GeV</td>
<td>$m(\pi_T) = 95$ GeV</td>
</tr>
<tr>
<td>190 GeV</td>
<td>$m(\pi_T) = 105$ GeV</td>
</tr>
<tr>
<td>200 GeV</td>
<td>$105 &lt; m(\pi_T) &lt; 115$ GeV</td>
</tr>
<tr>
<td>210 GeV</td>
<td>$115 &lt; m(\pi_T) &lt; 125$ GeV</td>
</tr>
<tr>
<td>220 GeV</td>
<td>$115 &lt; m(\pi_T) &lt; 125$ GeV</td>
</tr>
<tr>
<td>230 GeV</td>
<td>$125 &lt; m(\pi_T) &lt; 135$ GeV</td>
</tr>
<tr>
<td>240 GeV</td>
<td>$125 &lt; m(\pi_T) &lt; 145$ GeV</td>
</tr>
<tr>
<td>250 GeV</td>
<td>$135 &lt; m(\pi_T) &lt; 145$ GeV</td>
</tr>
</tbody>
</table>

TABLE VI: The excluded region at 95% C.L in the $(m(\rho_T)-m(\pi_T))$ plane for $\rho_T \to W \pi_T \to \ell \nu b \bar{q}$ production

VI. CONCLUSIONS

We have performed a search for technicolor production $p \bar{p} \to \rho_T^{\pm}/0 \to W^{\pm}, \pi_T^{0}/\pi^0 \to \ell \nu b \bar{b}$ or $\ell \nu b \bar{u}$ with the CDF II detector. We find that the dataset corresponding to 1.9 fb$^{-1}$ agrees with the background predictions within uncertainties. A large region of $m(\rho_T) = 180 - 250$GeV and $m(\pi_T) = 95 - 145$GeV are excluded at 95% CL, based on the Technicolor Straw Man model (Table VI).

VII. ACKNOWLEDGEMENTS

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