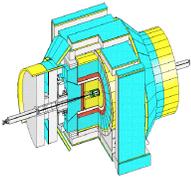


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# Some jet theory to set up Gene's talk

J. Huston

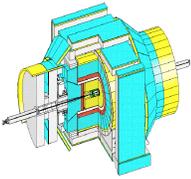
Michigan State University



# Global pdf fits



- Calculation of production cross sections at the Tevatron and LHC relies upon knowledge of pdfs in relevant kinematic range
- pdfs are determined by global analyses of data from DIS, DY and jet and direct  $\chi$  production
- Two major groups that provide semi-regular updates to parton distributions when new data/theory becomes available
  - ◆ MRS->MRST98->MRST99->MRST2001->MRST2002
  - ◆ CTEQ->CTEQ5->CTEQ5(1)->CTEQ6->CTEQ6.1(new)
- CTEQ6 is based on series of previous CTEQ distributions, but represents more than an evolutionary advance
  - ◆ update to new data sets
  - ◆ incorporation of correlated systematic errors for all experiments in the fit
  - ◆ new methodology enables full characterization of parton parametrization space in neighborhood of global minimum
    - ▲ Hessian method
    - ▲ Lagrange Multiplier
  - ◆ results available both in conventional formalism and in Les Houches accord format (more on this later)

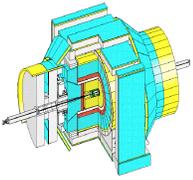


# Uncertainties in pdf's



- What's unknown about PDF's
  - ◆ the gluon distribution
  - ◆ strange and anti-strange quarks
  - ◆ details in the {u,d} quark sector; up/down differences and ratios
  - ◆ heavy quark distributions

- $\Sigma$  of quark distributions ( $q + qbar$ ) is well-determined over wide range of  $x$  and  $Q^2$ 
  - ◆ Quark distributions primarily determined from DIS and DY data sets which have large statistics and systematic errors in few percent range ( $\pm 3\%$  for  $10^{-4} < x < 0.75$ )
  - ◆ Individual quark flavors, though may have uncertainties larger than that on the sum; important, for example, for  $W$  asymmetry
- information on  $dbar$  and  $ubar$  comes at small  $x$  from HERA and at medium  $x$  from fixed target DY production on  $H_2$  and  $D_2$  targets
  - ◆ Note  $dbar \neq ubar$
- strange quark sea determined from dimuon production in  $\Sigma$  DIS (CCFR)
- $d/u$  at large  $x$  comes from FT DY production on  $H_2$  and  $D_2$  and lepton asymmetry in  $W$  production



# Uncertainties in pdf fits



- Two sources

- ◆ Experimental errors

- ▲ Hessian/Lagrange multiplier techniques designed to address estimate of these effects

- question is what  $\Delta\chi^2$  change best represents estimate of uncertainty (we use  $\Delta\chi^2$  of 100 (out of 2000))

- ◆ Theoretical

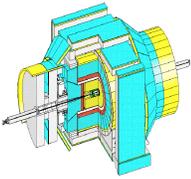
- ▲ higher twist/non-perturbative effects

- choose  $Q^2$  and  $W$  cuts to try to avoid

- ▲ higher order effects

- is NNLO necessary yet?

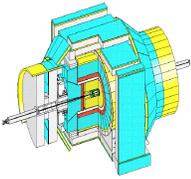
- ▲ edge of phase space effects



# Nuts/bolts of fits



- Functional form used is:
  - ◆  $xf(x, Q_0) = A_0 x^{A_1} (1-x)^{A_2} e^{A_3 x} (1 + A_4 x)^{A_5}$ 
    - ▲  $Q_0 = 1$  GeV (below any data used in fit)
      - easier to do forward evolution than backward
    - ▲ functional form arrived at by adding a 1:1 Pade expansion to quantity  $d(\log xf)/dx$
    - ▲ more versatile than form used in CTEQ5 or MRST
    - ▲ there are 20 free parameters used in the global fit
- Light quarks treated as massless; evolution kernels of PDFs are mass-independent
- Zero mass Wilson coefficients used in DIS structure functions



# D0 jet cross section



- CTEQ4 and CTEQ5 had CDF and D0 central jet cross sections in fit
- Statistical power not great enough to strongly influence high  $x$  gluon
  - ◆ CTEQ4HJ/5HJ required a special emphasis to be given to high  $E_T$  data points
- Central fit for CTEQ6 is naturally *HJ-like*
- $\chi^2$  for CDF+D0 jet data is 113 for 123 data points
- Note the power of having *search for new physics regions and control regions*
  - ◆ pdf explanation should work for all regions
  - ◆ new physics should be central

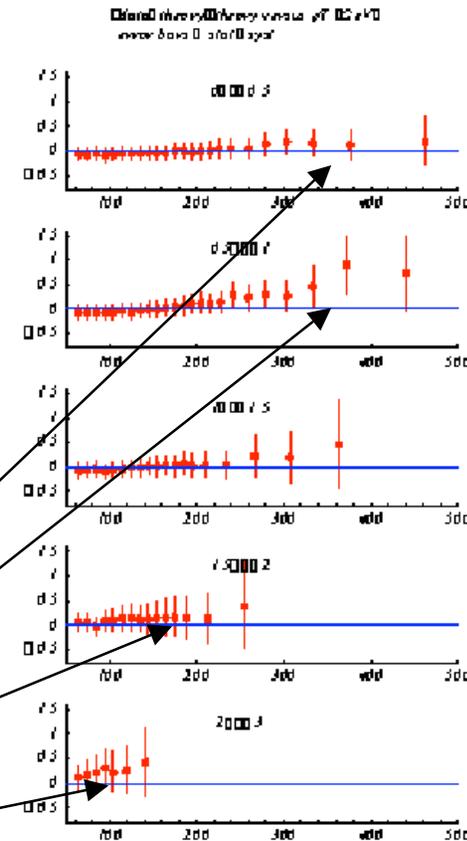
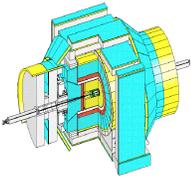


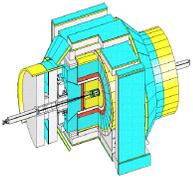
Figure 10: The ratio between the central jet cross section to the CTEQ4 prediction (with all systematic errors included)



# Remaining questions/discussion



- HJ-like behavior for gluon motivated to describe CDF Run 1 data now seems a natural consequence of global fitting
- ...but need to be careful not to hide any possible new physics in the gluon uncertainty
- How much room is left over for new physics in Run 1 data?
- What is uncertainty on Run 2 jet cross sections?
- Is NLO QCD valid for description of jet cross sections in full range of  $E_T/p_T$  and rapidity, for both Run 1 and Run 2?
- These and other questions to be answered in CTEQ paper in preparation
- Here are a few excerpts

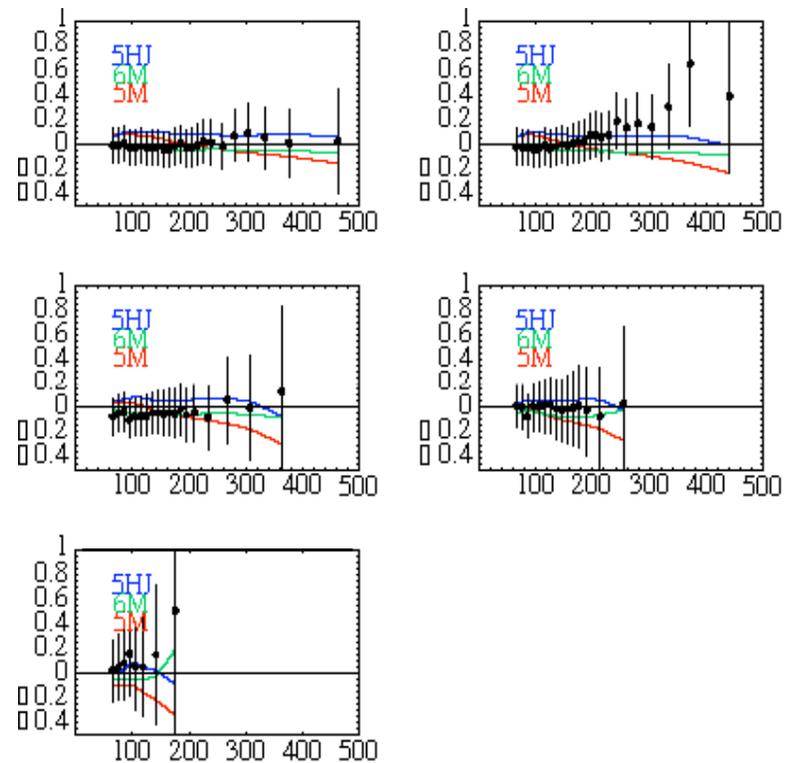
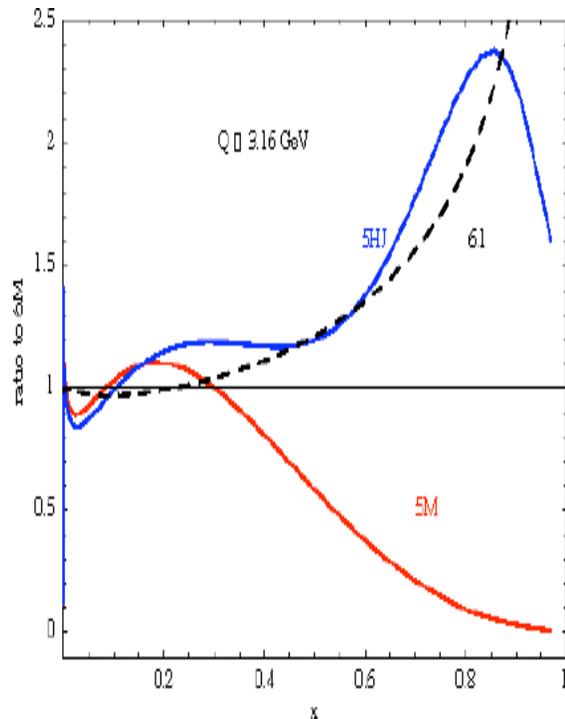


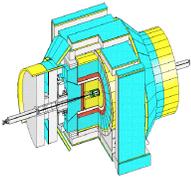
# 6.1 gluon compared to 5M/5HJ/6M



- In the course of investigations for this paper, some improvements to the analysis were made that changed the gluon distribution: cteq6m->cteq6.1

- small changes in jet cross sections, as for example the D0 jet cross section below

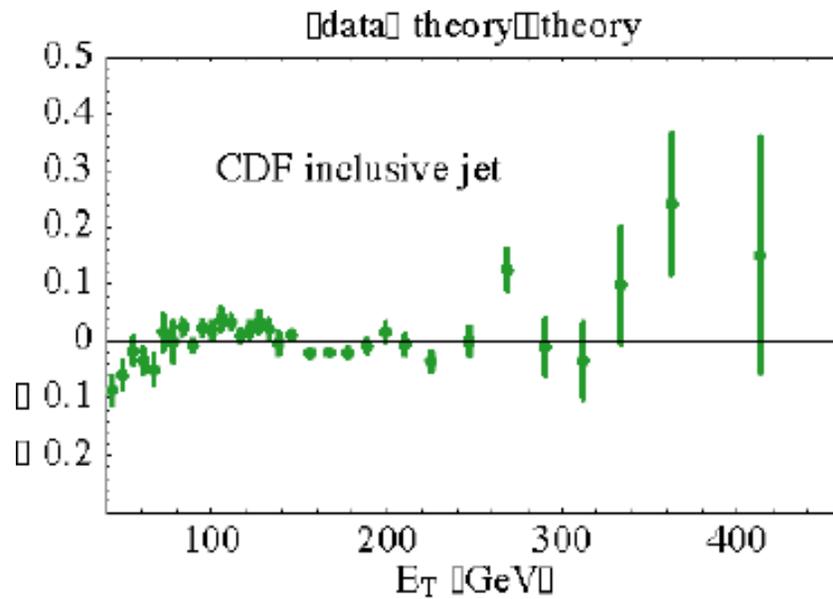




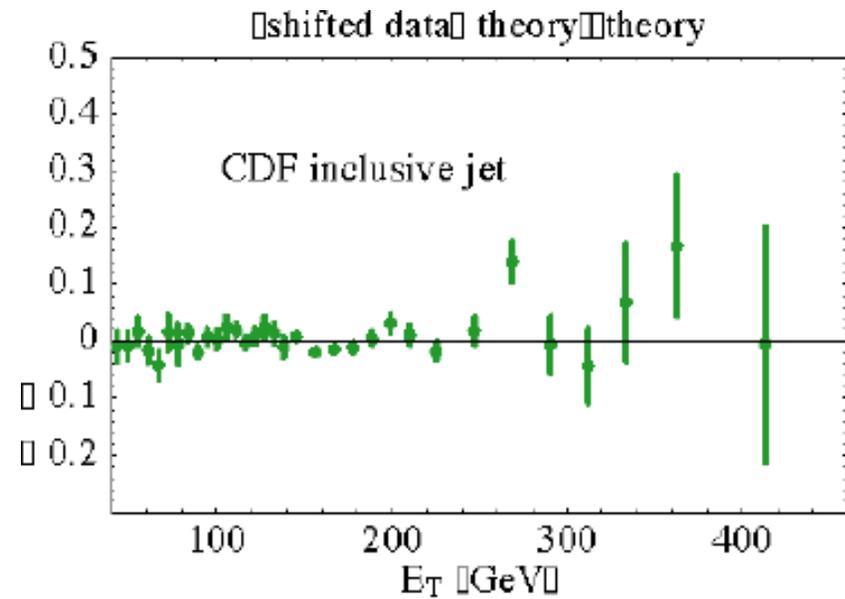
# CDF Run 1 data and CTEQ6.1

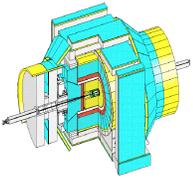


before systematic error shifts



after systematic error shifts

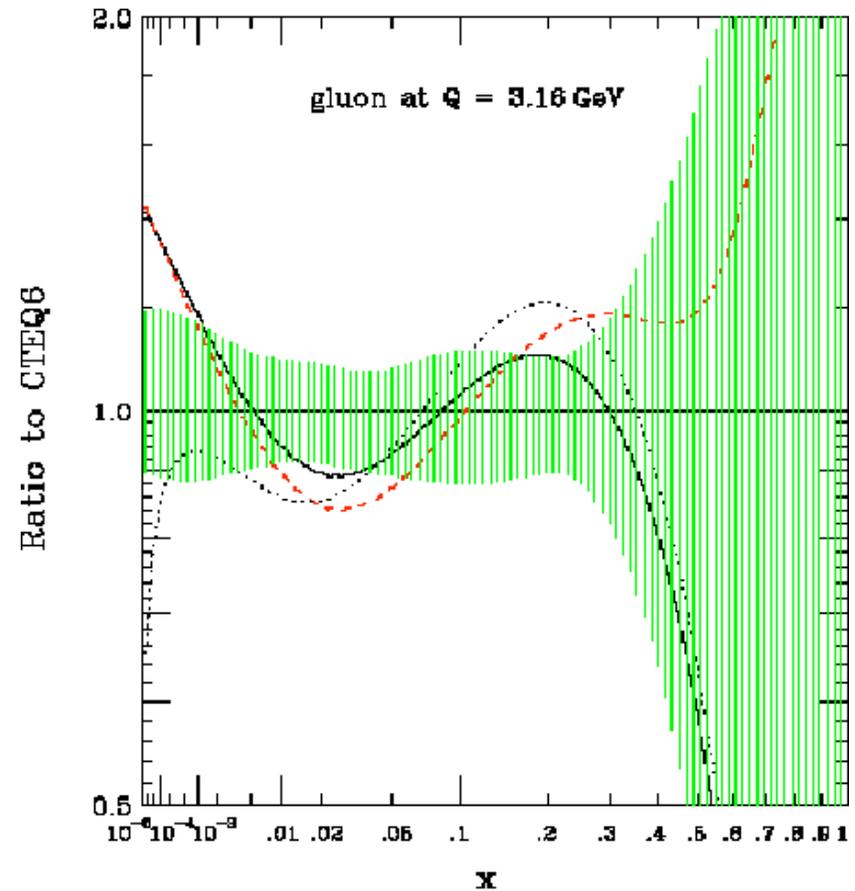
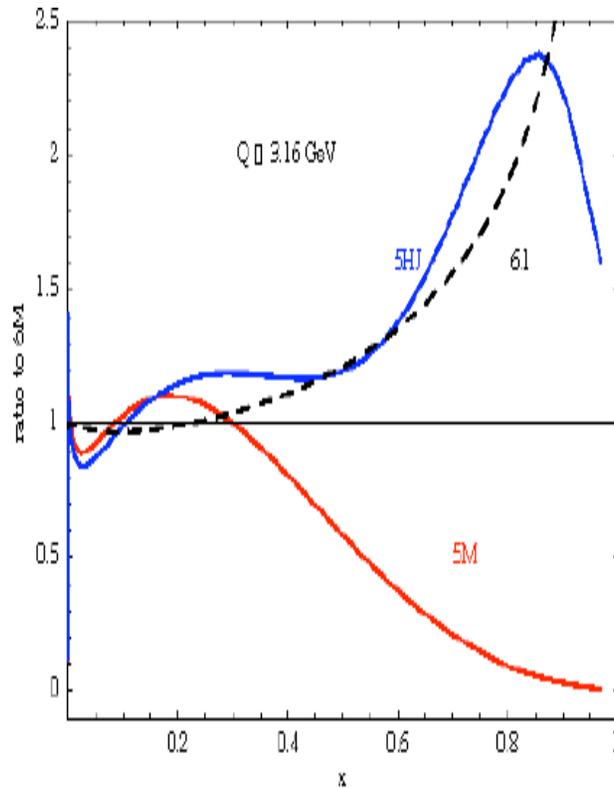


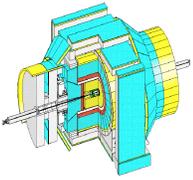


# Remaining gluon uncertainties



New cteq6.1 well within uncertainty band for gluon distribution





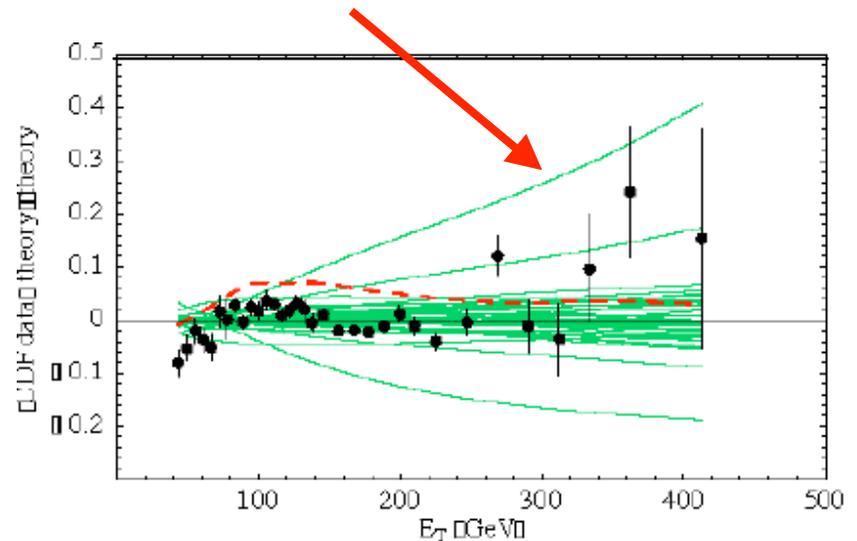
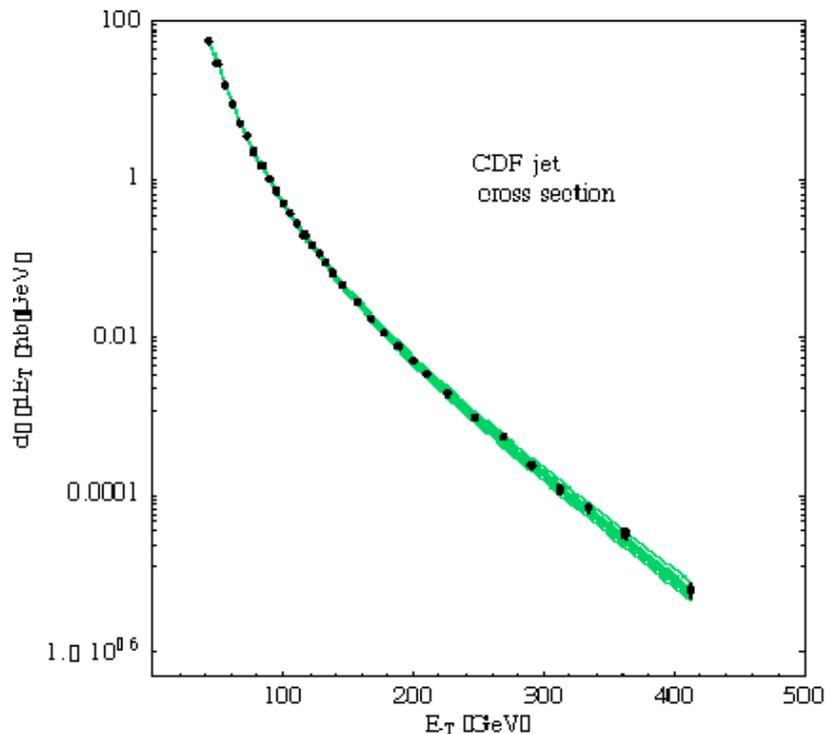
# PDF uncertainties for Run 1 cross section

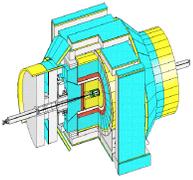


- 20 free parameters in the fit
- In the Hessian method, a 20X20 matrix is diagonalized and 20 orthogonal eigenvector directions in parameter space are determined

- Each eigenvector direction corresponds to some linear combination of pdf parameters
- Large eigenvalues correspond to highly determined directions (e.g. valence quarks)
- Small eigenvalues correspond to poorly determined directions (high x gluon)
- Result is 40 pdf's (go along + and - direction  $\pm \sqrt{\lambda}$  of 100 for each eigenvalue)

Note 1 eigenvector(15+) leads to noticeably larger prediction than the others





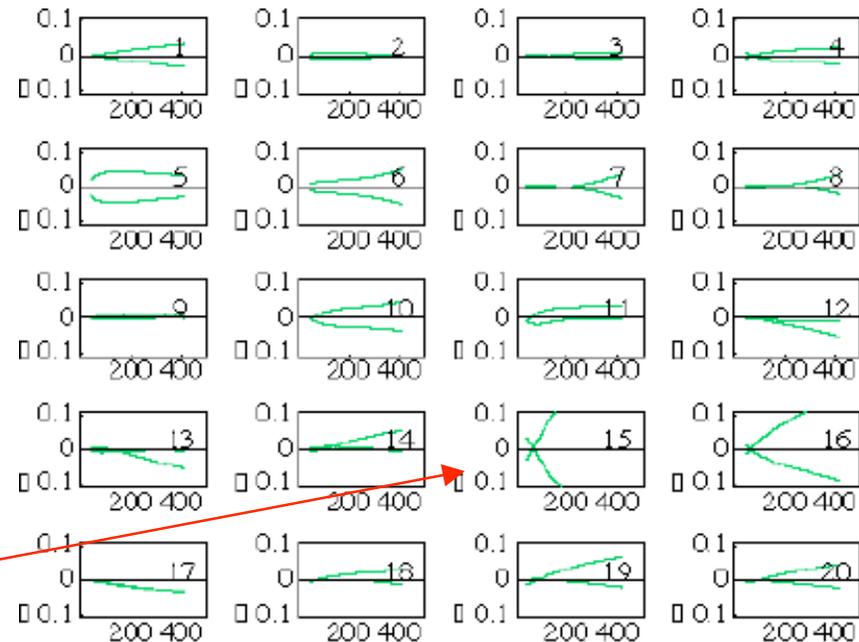
# CDF jet cross section uncertainties

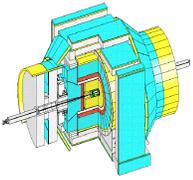


- On the right are shown the uncertainties for the CDF jet cross section along each eigenvector ( $\Delta\sigma^2 = 100$ )

- ◆ jet cross section most sensitive to eigenvector 15

- ▲ which mainly contains parameters relating to behavior of high  $x$  gluon

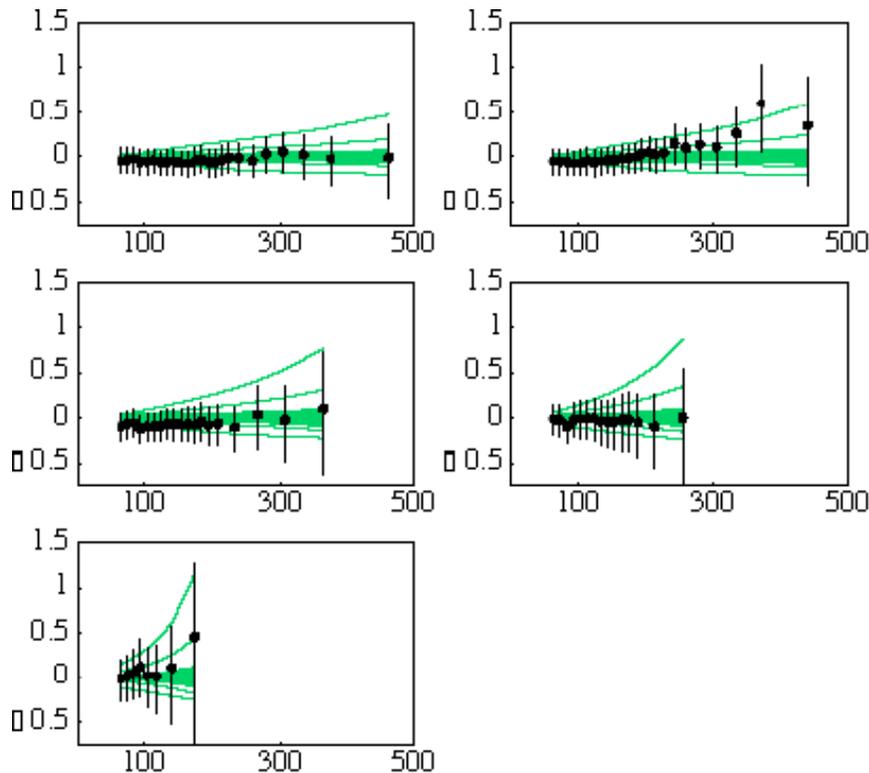




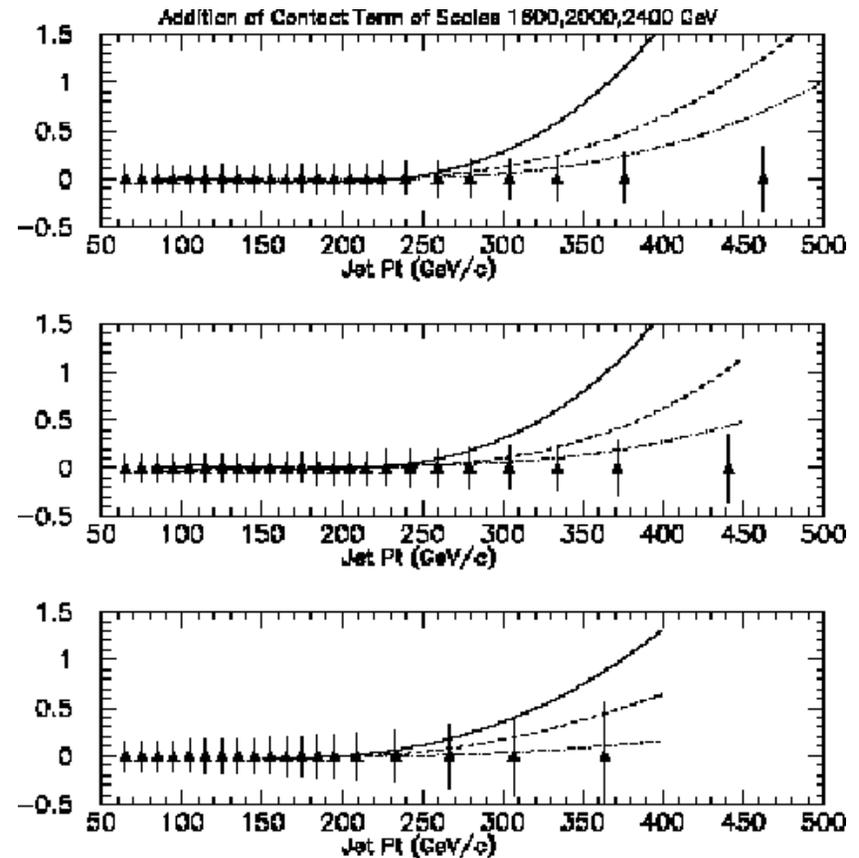
# Room for new physics

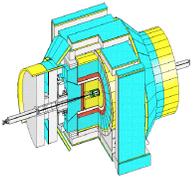


- Uncertainties for D0 Run 1 jet cross section as a function of rapidity



Effect of new physics (compositeness) for the first 3 rapidity bins (effects on last bins are negligible)

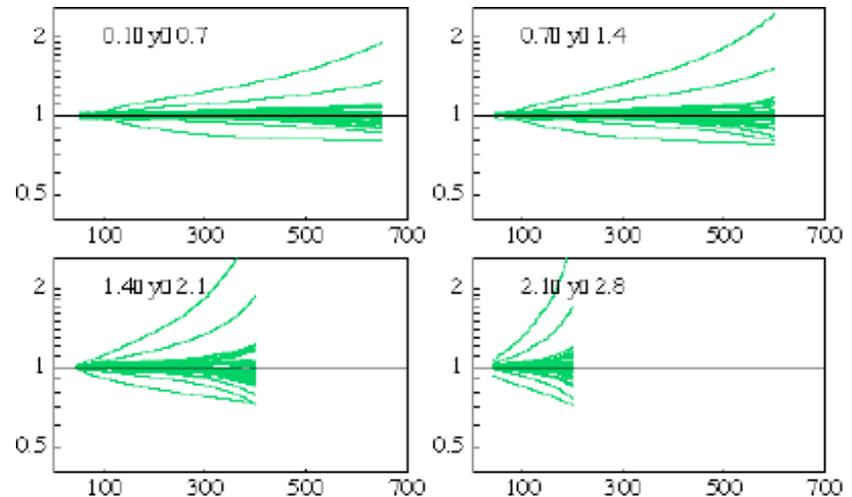
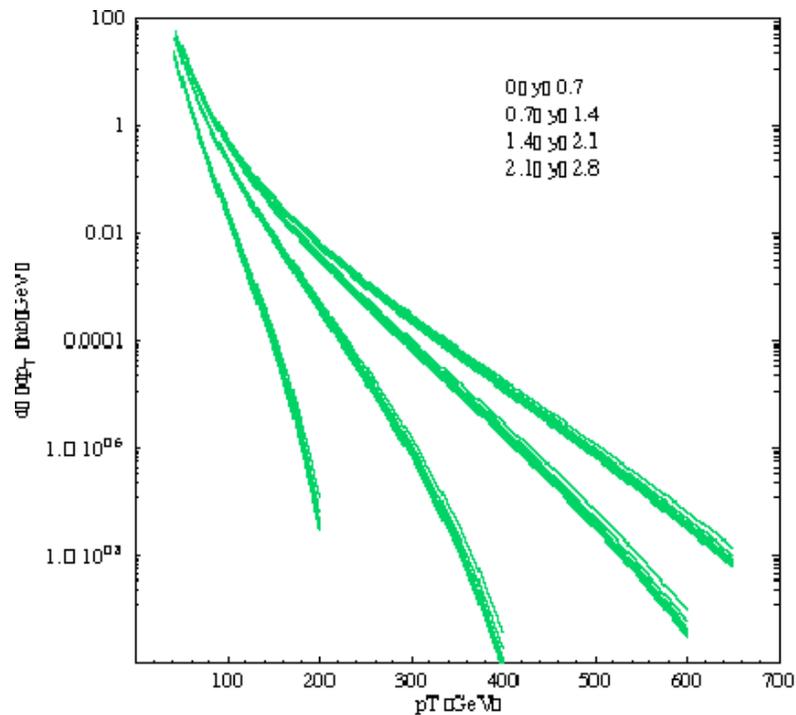


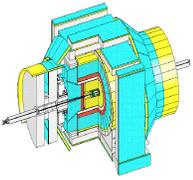


# Uncertainties on Run 2 predictions for CDF



- CDF will measure the inclusive jet cross section in the forward regions as well



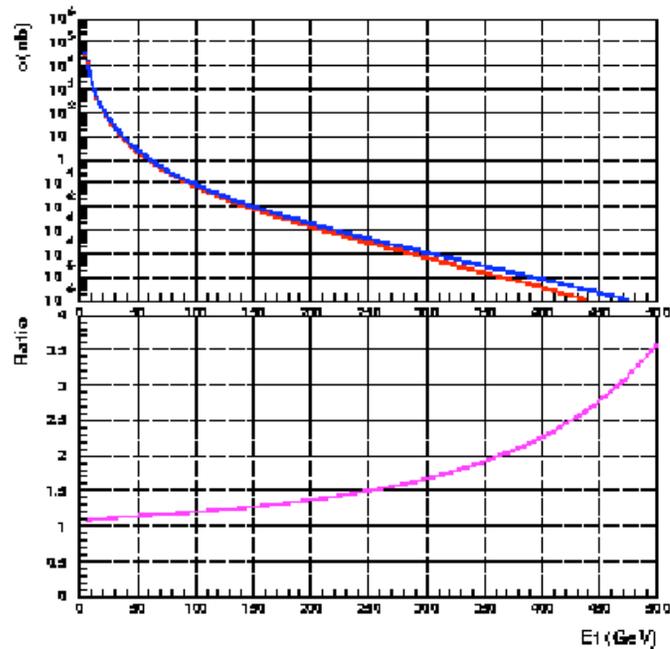


# Ratio of Run 2/Run 1

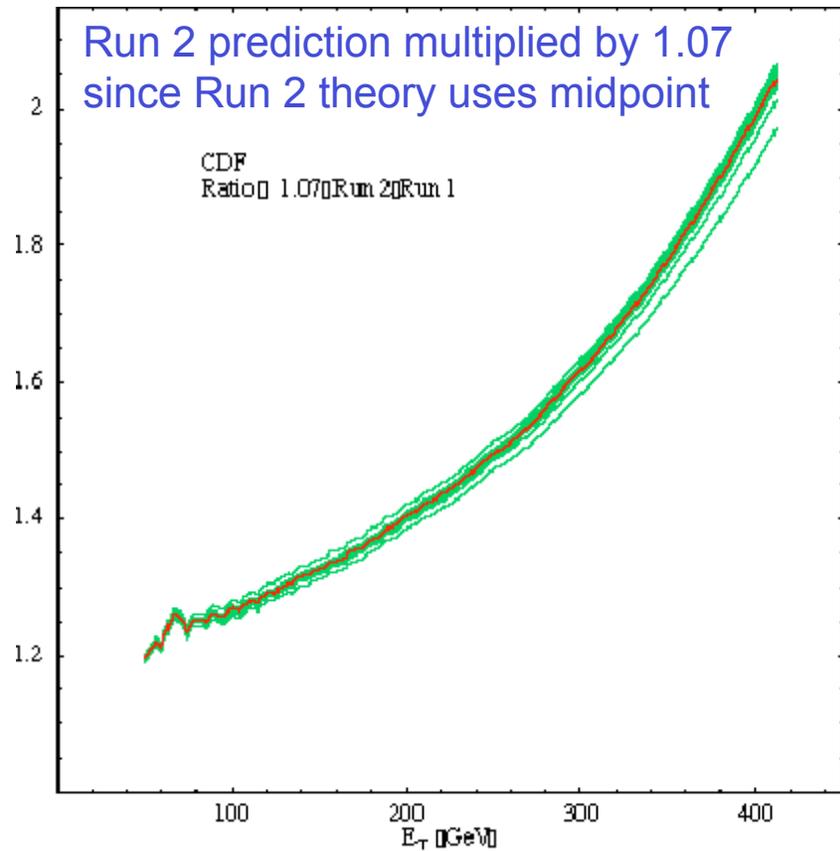


The increase in the center-of-mass energy from 1.8 to 1.96 TeV has a large effect on the high  $E_T$  jet rate.

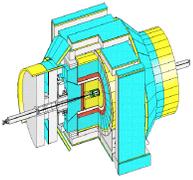
Inclusive jet cross section at 1.8 and 2.0 TeV (CTEQ4HJ)



For the full Run IIa sample the number of jets above 100 GeV will increase from 11 to  $\sim 500$ .



Small theoretical uncertainties on the ratio; experimental uncertainties should partially cancel as well

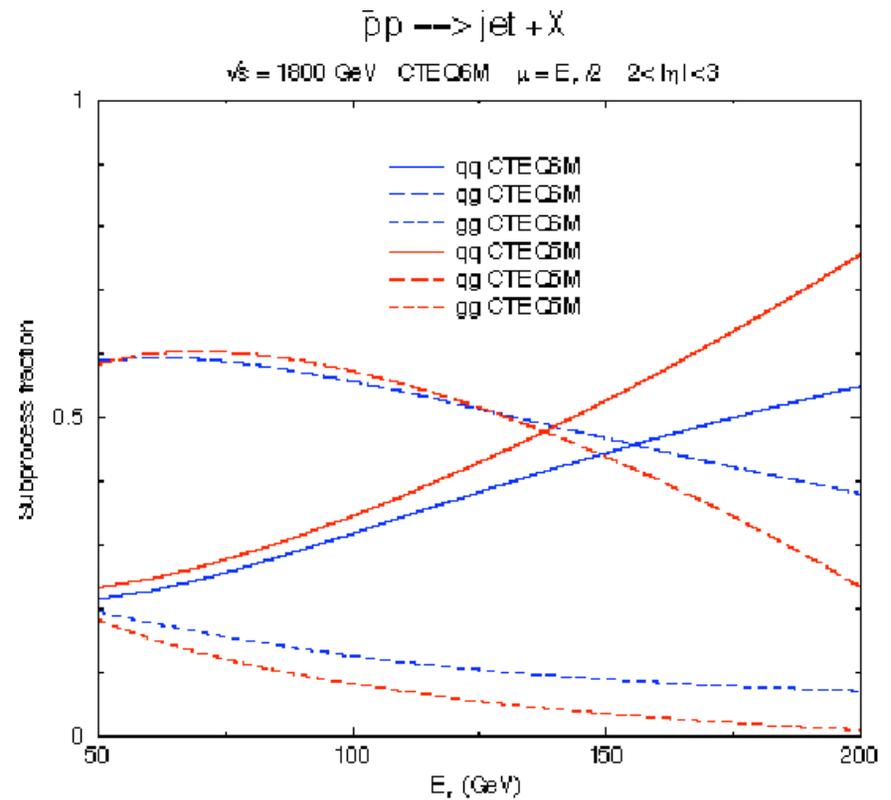
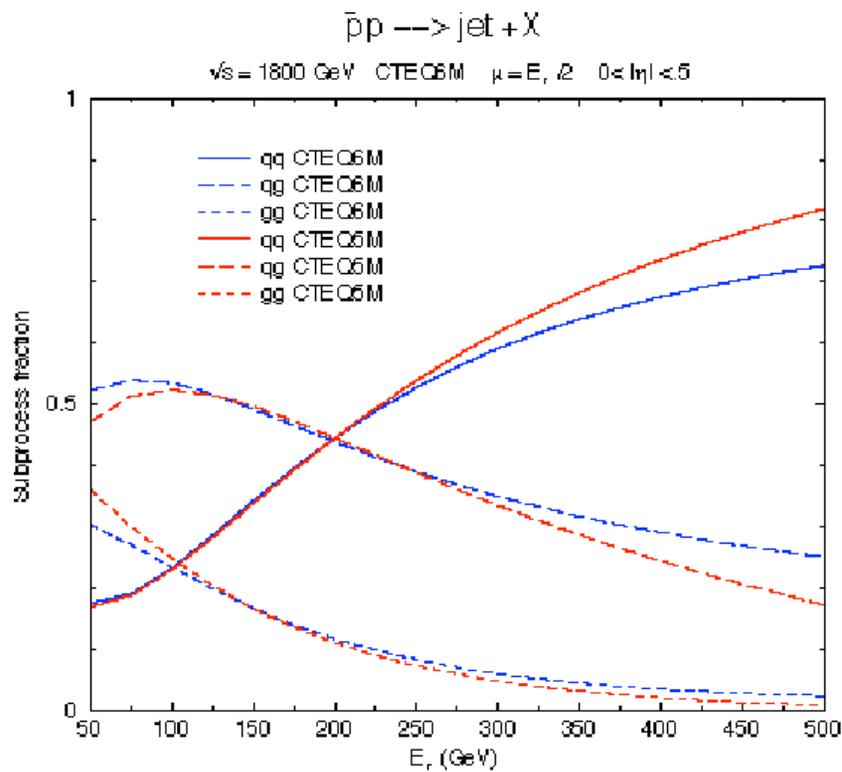


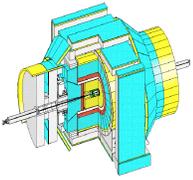
# Subprocess plots



- Enhanced gluon-quark scattering in CTEQ6

Larger contribution for forward rapidities

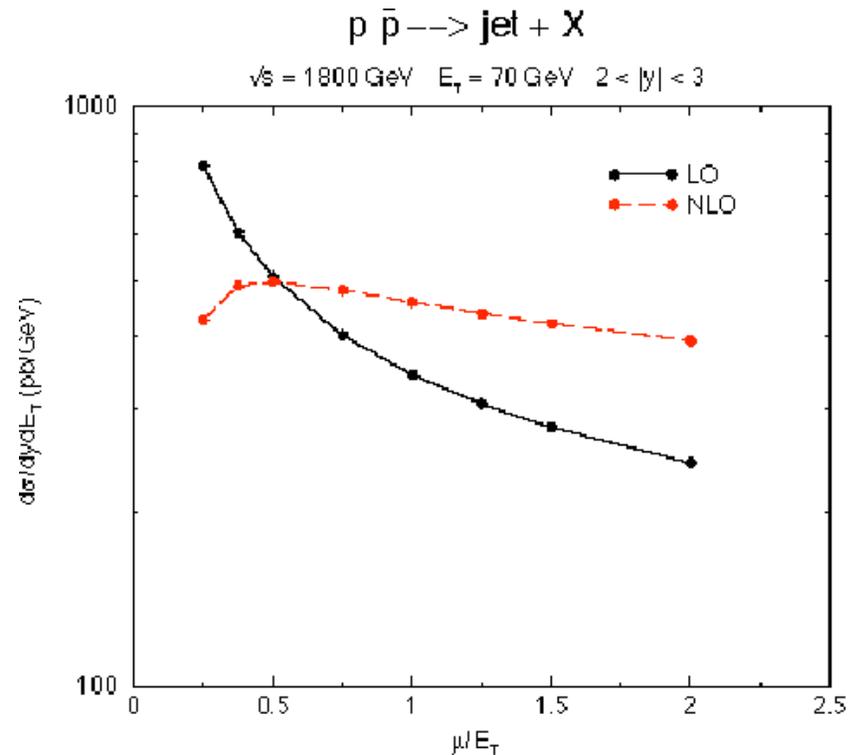


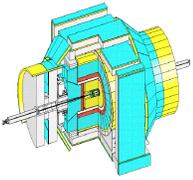


# Reliability of NLO QCD



- One way of looking for possible effects of higher order terms is to examine the scale dependence of the jet cross section
- Also, look at the K-factor (NLO/LO) at scale cross section evaluated at
- For  $\mu/E_T=0.5$ , K-factor is of the order of 1 for most of Tevatron jet data
- Is NLO QCD reliable for all data taken in Run 1?

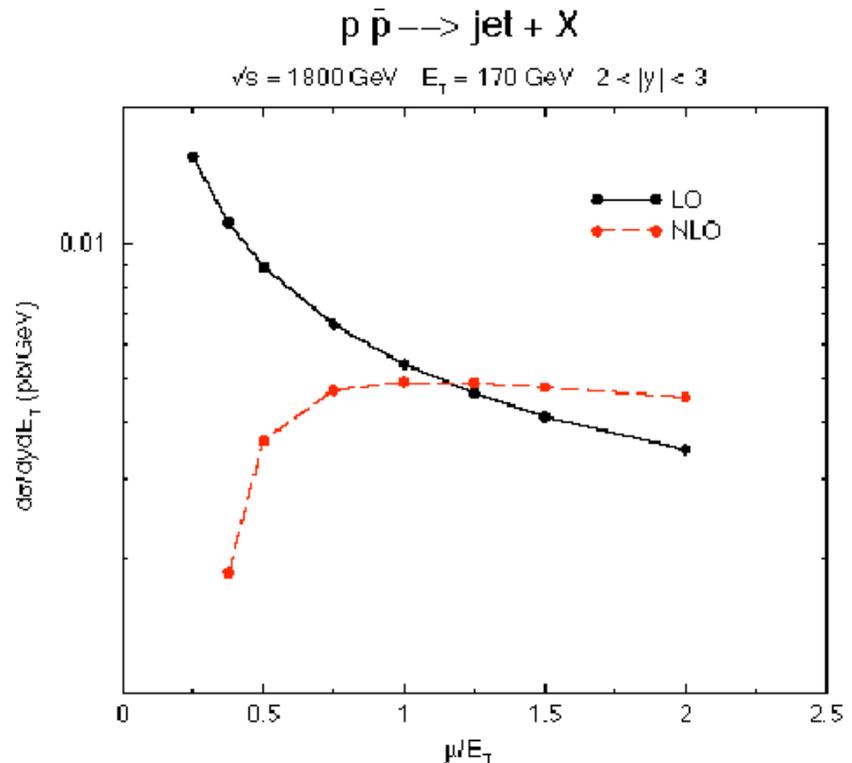


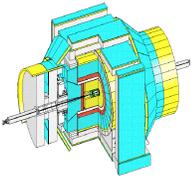


# NLO and LO scale dependence



- Leading log calculations of the jet cross section at high  $E_T$  generally decrease as renorm/factorization scales are increased
- At NLO, scale dependence is reduced and cross section typically has parabolic shape as function of scale
- For large  $y$ , high  $E_T$ , point where NLO curve crosses LO curve moves out in  $\mu/E_T$ , due to kinematic reasons
- So, K-factor at  $\mu/E_T$  can be significantly less than 1
- Not a cause for alarm per se, but scale dependence of cross section does become steep at scales lower than  $0.5 E_T$



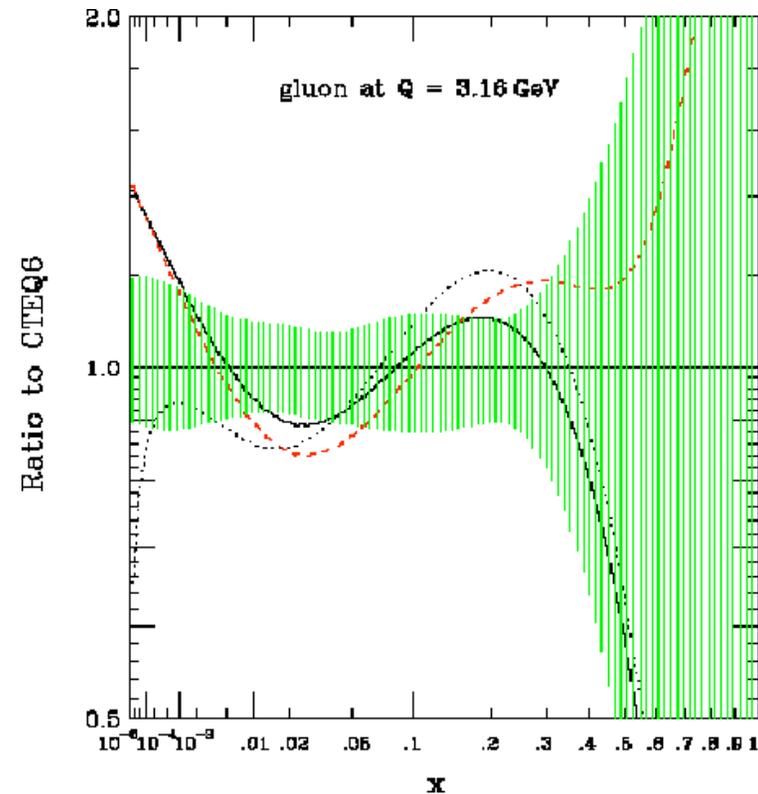
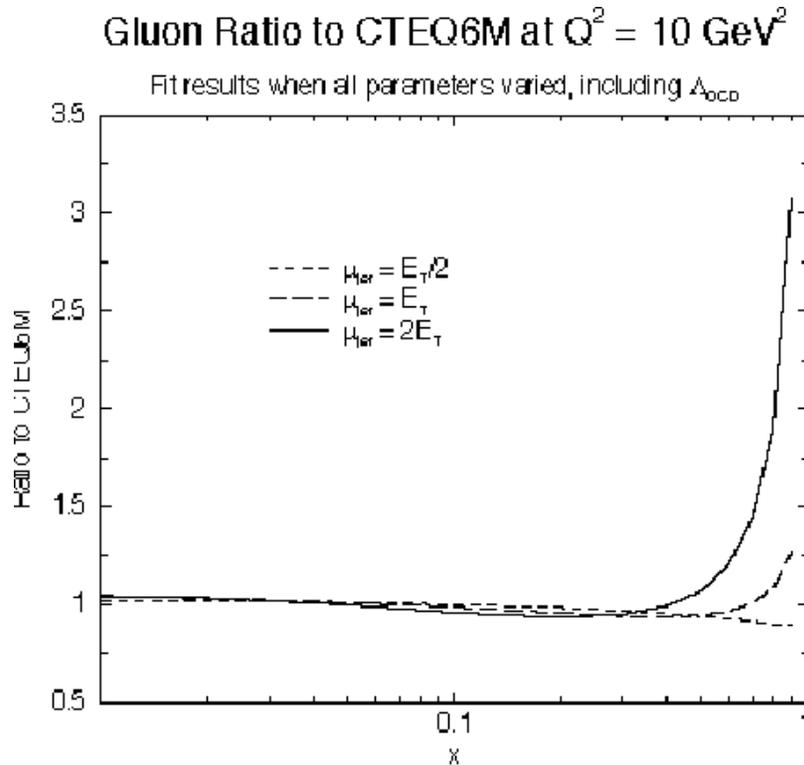


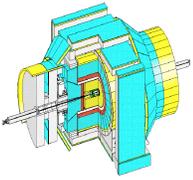
# Scale dependent fits



- Repeat CTEQ6 fit using scales of  $E_T$ ,  $2E_T$  for jet cross section
- Compare gluon distribution to CTEQ6M ( $\alpha = E_T/2$ )

Variations within CTEQ6 pdf uncertainty band

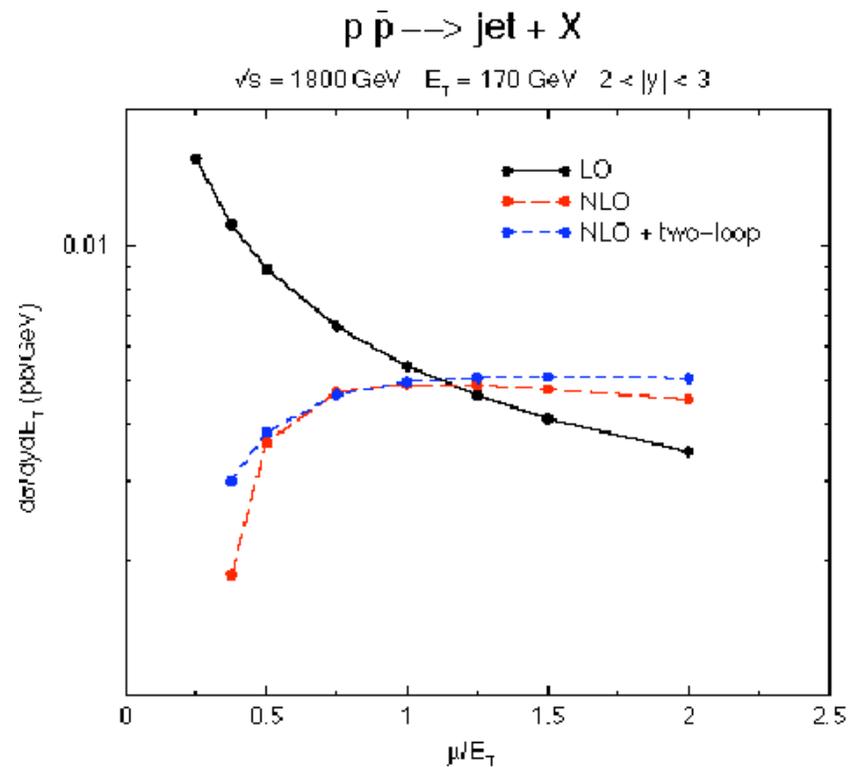


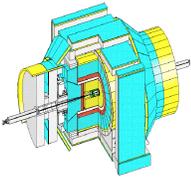


# Edge of phase space



- Potentially significant problems from phase space limitations at large  $E_T$  and rapidity
- Such corrections can be treated by threshold resummation techniques
- Estimated at two loop level and found to be small for central rapidity region
- Repeated for forward rapidity region and again are not significant at scale used in fit

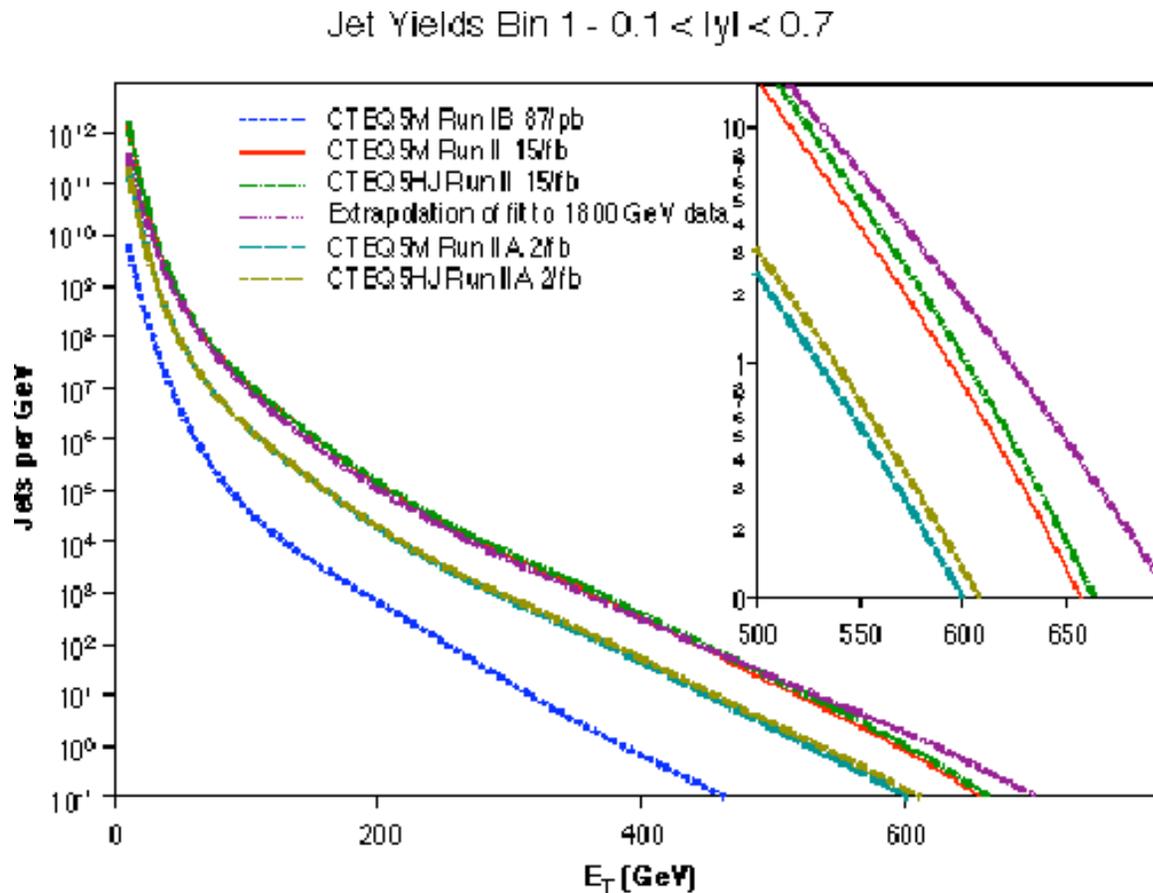


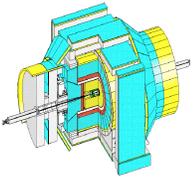


# Jet Yields in Run II (from Steve Ellis)



See [http://www.pa.msu.edu/~huston/run2btdr\\_qcd/tdr.ps](http://www.pa.msu.edu/~huston/run2btdr_qcd/tdr.ps).

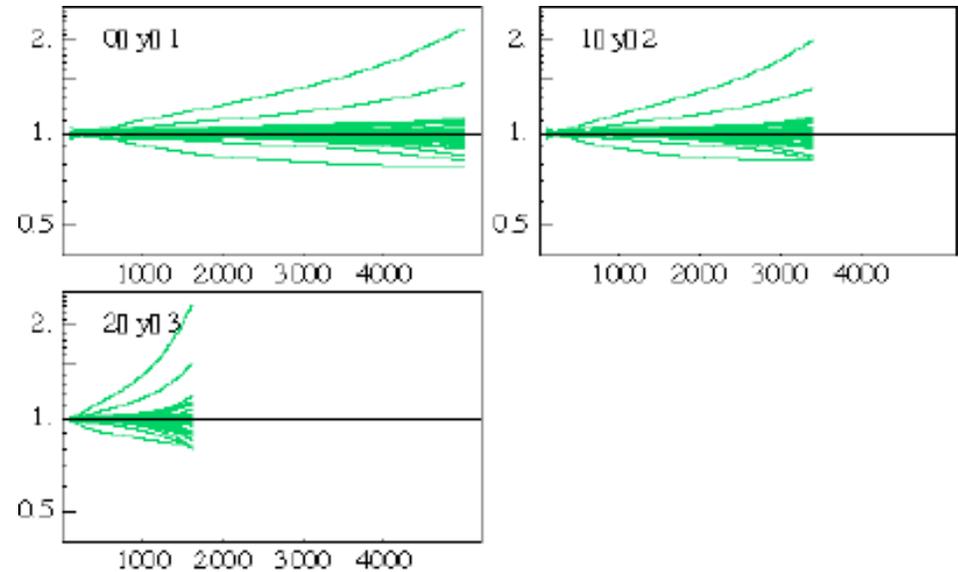
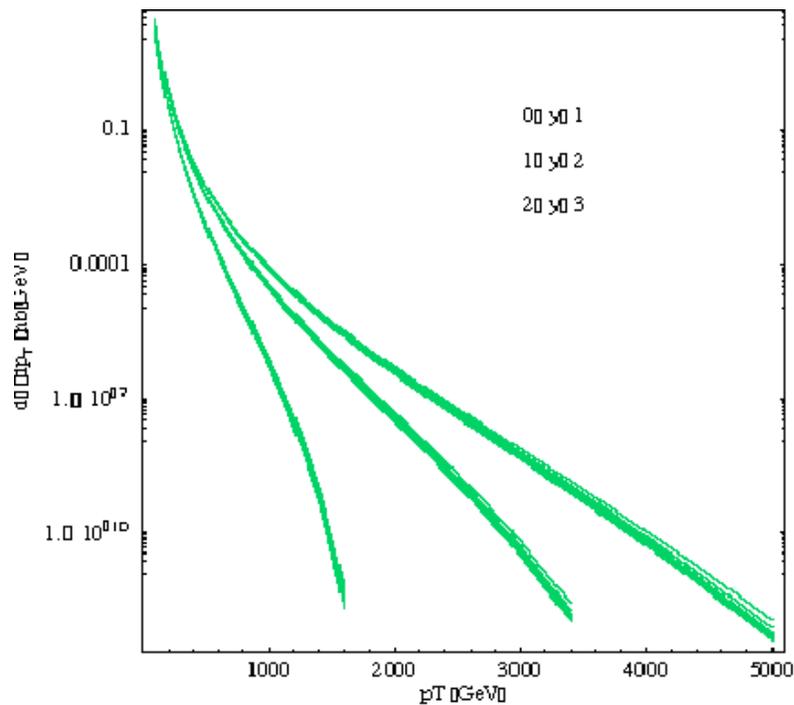


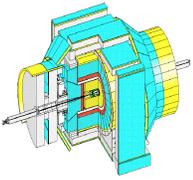


# Jet cross sections at the LHC



- Apply the same exercise to the LHC





# Conclusions regarding jet uncertainties



- Paper should be finished in  $<1$  month
  - ◆ uncertainties on jet cross sections in both Run 1 and Run 2
  - ◆ room for new physics in both