

## Search for top+jet resonances in $t\bar{t}$ +jet(s) at CDF.

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We present a search for a heavy new particle  $M$  produced in association with a top quark ( $p\bar{p} \rightarrow M t/M\bar{t}$ ) and decaying via  $M \rightarrow \bar{t}q/tq$ , leading to a resonance in the  $\bar{t}/t$ +jet system of  $t\bar{t}$ +extra jet events. We use events with exactly one lepton, missing transverse energy and at least five jets in data with an integrated luminosity of  $8.7 \text{ fb}^{-1}$ . We find the data to be consistent with the Standard Model and set cross-section upper limits from 0.61 pb to 0.02 pb for resonances ranging from 200 GeV to 800 GeV. We reinterpret these cross-section limits for specific physics models as exclusions in mass-coupling space.

### I. INTRODUCTION

We present a search for top+jet resonances in  $t\bar{t}$ +jet events using the CDF detector [1]. Recently, CDF reported a measurement of the top-quark production forward-backward asymmetry ( $A_{FB}$ ) that is significantly larger than predicted by the SM [2];  $D\bar{O}$  sees a consistent result [3].

A wide class of models [4] have been built to explain such a discrepancy, most involving the production of a new heavy mediating particle  $M$  that enhances  $A_{FB}$ . Such new particles may also be singly produced in association with a top(antitop)-quark and further decay to a antitop(top)-quark and an additional jet.

$$p\bar{p} \rightarrow M t(\bar{t}) \rightarrow \bar{t}jt(tj\bar{t})$$

Such a decay would look like a top+jet resonance in  $t\bar{t} + \text{jet}$  events.

### II. SIGNAL AND SELECTION

Our signal is similar to  $t\bar{t}$  with an additional jet, in the lepton + jets channel the detector signature is  $\ell + \nu + qq' + b\bar{b}' + q$ . To isolate such a final state we require:

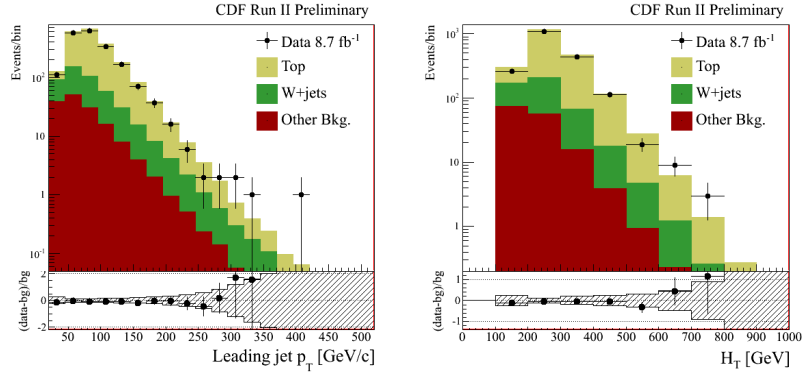


FIG. 1: Validation in the  $t\bar{t}$  region, with exactly four jets and at least one  $b$ -tag.

- Exactly one tight electron or muon with  $p_T > 20\text{GeV}$
- At least 5 jets with  $E_T > 20\text{GeV}$  and  $|\eta| < 2.0$
- At least 1 SECVTX [5] tag
- $\cancel{E}_T \geq 20\text{ GeV}$ .

### III. MODELING AND BACKGROUNDS

We model the  $M$  resonance signal using MADGRAPH [6] to describe the hard process and PYTHIA [7] for the showering. The dominant backgrounds with our selection are  $t\bar{t}$  and  $W$ +jets. The  $t\bar{t}$  and diboson background samples are generated using PYTHIA,  $W$ +jets and  $Z$ +jets using ALPGEN [8]+PYTHIA, and the QCD sample is modeled from jet events in the data.

We validate our background in three control regions:

- $t\bar{t}$  region: Exactly four jets and at least one  $b$ -tag. This validates the overall background normalization, as well as the modeling of  $m_{tj}$ ; the requirement of exactly four jets depletes the sample of potential signal contamination, Figure 1.
- $t\bar{t} + j$  region: At least five jets and exactly zero  $b$ -tags. This validates the modeling of additional jet radiation, which is the source of SM events with at least five jets. The requirement of exactly zero  $b$ -tags suppresses the signal, Figure 2.
- Low  $H_T$  region: At least five jets and small  $H_T$  ( $< 225\text{ GeV}$ ). This region directly probes the  $t\bar{t} + \text{jet}$  contribution, which is partially suppressed in the  $t\bar{t} + \text{jet}$  region by the low  $H_T$  requirement, Figure 3.

Our signal region is then defined as  $N_{jets} \geq 5$  and at least 1 tag.

We consider several sources of systematic uncertainty, including jet energy scale, contributions from additional interactions, uncertainty in descriptions of initial and final state radiation, differences in Monte Carlo generators for  $t\bar{t}$  background and the  $Q^2$  systematics for  $W$ +jets backgrounds. The impact of these systematics are listed in Table I.

### IV. RECONSTRUCTING THE RESONANCE MASS.

We construct the resonance mass,  $m_{tj}$ , using the top kinematic fitter. From the  $N$  ( $\geq 5$ ) jets we pick the 4 jets that have highest likelihood match to a  $t\bar{t}$  topology, the remaining  $N - 4$  are paired with the  $t/\bar{t}$ , the highest invariant mass of a jet +  $t/\bar{t}$  is chosen as our analysis variable,  $m_{tj}$ . Figure 4 shows this reconstruction for the backgrounds, data and example signal resonances of 300, 500 and 800  $\text{GeV}/c^2$  in the signal region.

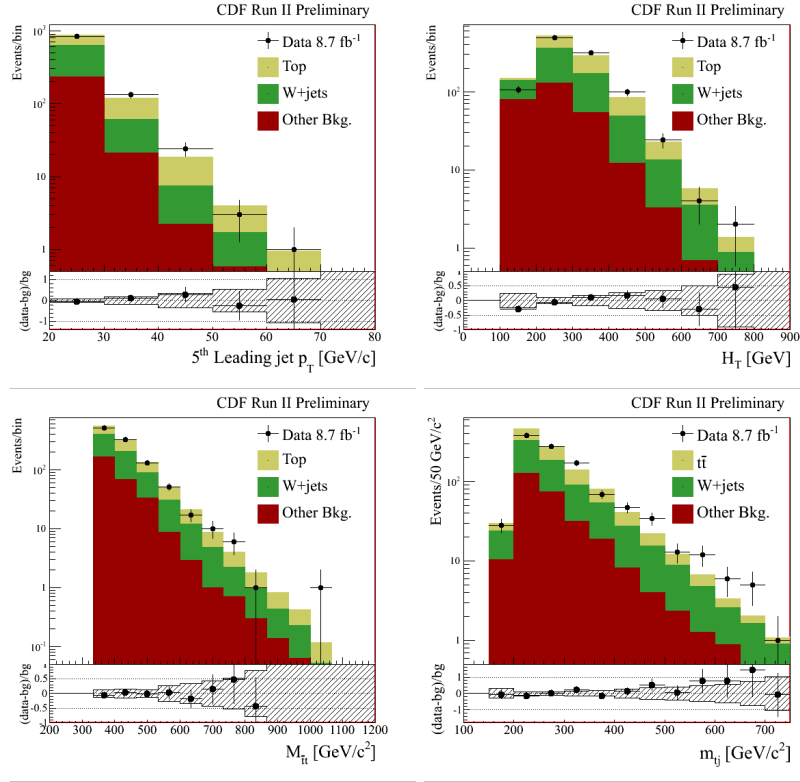


FIG. 2: Validation in the  $W$ +jets region, with at least 5 jets and exactly 0  $b$ -tag.

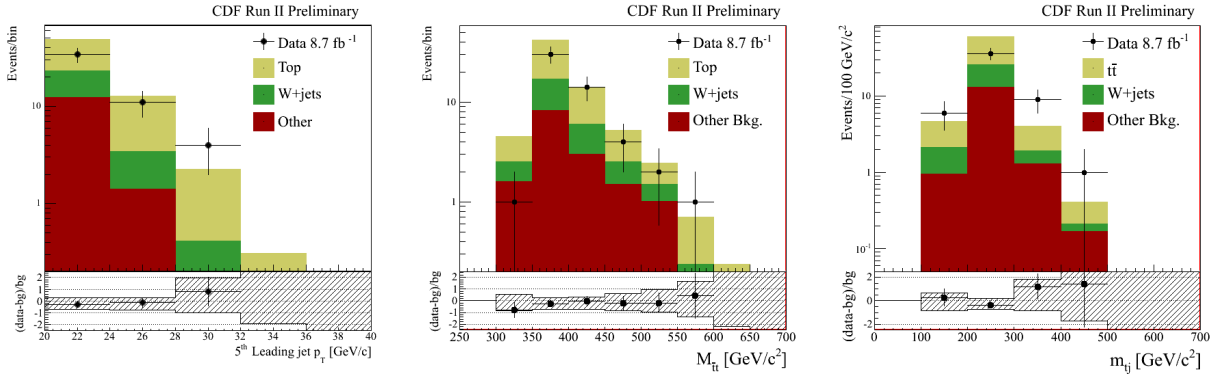


FIG. 3: Validation in the  $t\bar{t} + j$  region, with at least 5 jets, at least 1  $b$ -tag, and  $H_T < 225 \text{ GeV}$ .

TABLE I: Impact of systematic uncertainties on each background source and an example signal of 500 GeV in the signal region.

Systematic	$t\bar{t}$	$W$ +jets	Total	$M$ (500 GeV)
Nominal	550.55	78.64	669.17	339.69
JES	17%	15%	16%	9%
Radiation	6%	-	5%	4%
$Q^2$	-	19%	2%	-
Nvtx	3%	2%	3%	2%
$t\bar{t}$ Generator	6%	-	5%	-
Normalization	10%	30%	12%	-
Total Uncertainty	22%	38%	21%	10%

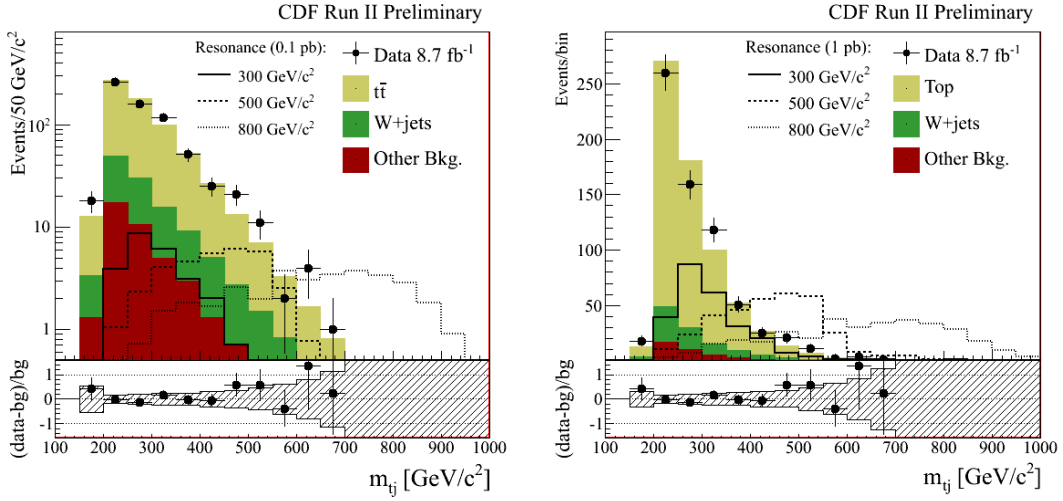


FIG. 4: Resonance mass reconstruction,  $m_{tj}$ , for backgrounds, data and example signal resonances of 300, 500 and 800  $GeV/c^2$  in the signal region scaled to a cross-section of (left) 0.1  $pb$  with log-scale and (right) 1  $pb$  on a linear scale.

## V. RESULTS: LIMITS ON CROSS SECTION

To extract the most likely signal cross-section, we perform a binned maximal-likelihood fit to the  $m_{tj}$  distribution, varying each background rate within uncertainties, allowing shape and rate variation due to systematic uncertainties described above. The signal and background rates are fitted simultaneously. The CLs method [9] is used to set 95% cross-section upper limits. The median expected upper limit is extracted in the background-only hypothesis.

We find that the observed limits are consistent with what we would expect if the data were drawn from the standard model. We estimate 95% C.L. upper limits on the cross-section of such resonances, see Figure 5.

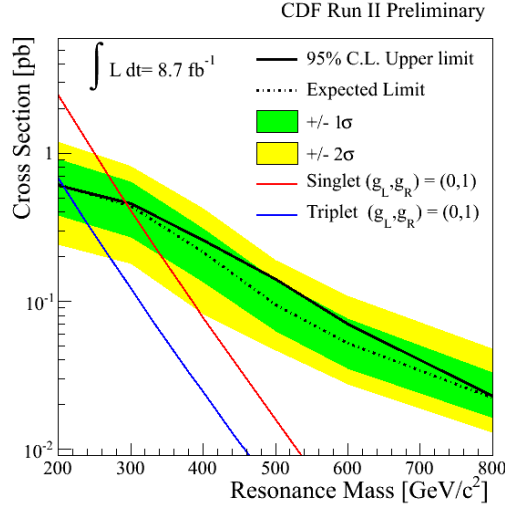


FIG. 5: Upper limits at 95% CL on  $t\bar{t} + j$  production via a heavy new mediator  $M$ , as a function of the mediator mass. Also shown are theoretical predictions, assuming a unit coupling.

## VI. RESULTS

We convert limits on top+jet resonance to exclusion of specific models in mass-coupling space, see Fig. 6

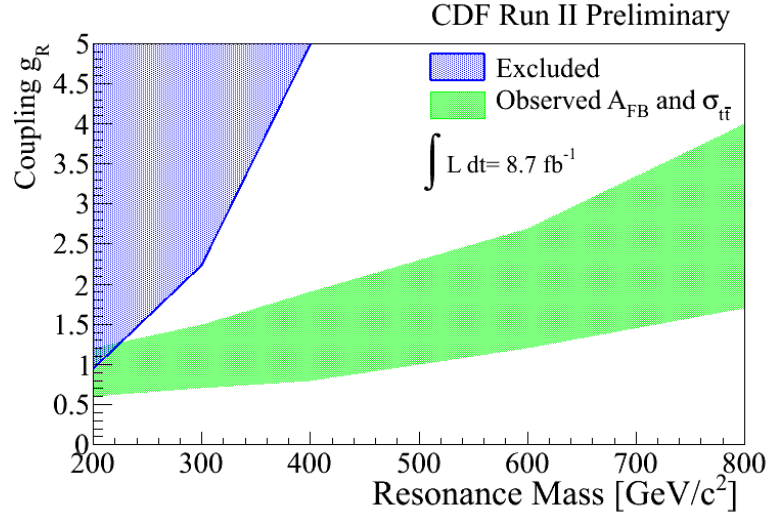
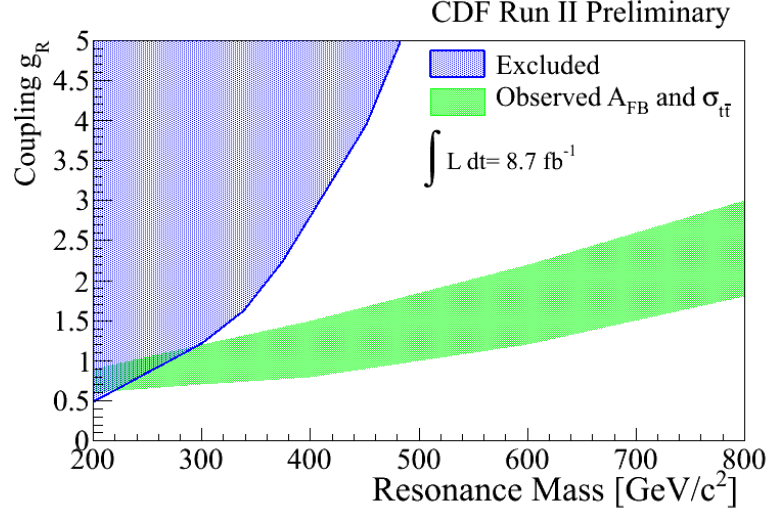


FIG. 6: Excluded region in mass-coupling space for two specific models, where the  $M$  particle is part of a new singlet (a) or colored triplet [4] (b). Also shown are regions [10] which are consistent with the observed anomalous  $A_{\text{FB}}$  and constraints from top-quark pair production and single-top production cross-section measurements.

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  - [10] This region simultaneously satisfies the observed high- $m_{t\bar{t}}$   $A_{\text{FB}}$ , low- $m_{t\bar{t}}$   $A_{\text{FB}}$ , and the  $t\bar{t}$  cross-section better than the Standard Model. Mathematically, it is defined as the region with  $\chi^2 < 2.8$ , where  $\chi^2$  is defined in Equation 22 in Ref. [4].  $\chi^2$  for the Standard Model is 2.8.