CDF note 8802

Top Dilepton Cross Section in 1.2 fb^{-1} using the DIL Selection

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(Dated: May 8, 2007)

A measurement of the $t\bar{t}$ production cross section in $p\overline{p}$ collisions at $\sqrt{s}$=1.96 TeV using events with two leptons is reported. The data were collected by the CDF II Detector. The result in a dataset corresponding to an integrated luminosity 1.2 fb^{-1} is:

$$\sigma_{t\bar{t}} = 6.16 \pm 1.05_{\rm stat} \pm 0.72_{\rm syst} \pm 0.37_{\rm lumi} \text{ pb}$$

Preliminary Results for Summer 2007 Conferences

I. INTRODUCTION

This note describes a measurement of the $t\bar{t}$ production cross section in $\bar{p}p$ collisions at $\sqrt{s} = 1.96$ TeV with the CDF detector at the Fermilab Tevatron. The measurement is based on the identification of both leptons in the decay chain $t\bar{t} \rightarrow (W^+b)(W^-\bar{b}) \rightarrow (l^+\bar{\nu}_l)(l^-\nu_l\bar{b})$. Therefore it selects decays with two high transverse energy leptons, high missing transverse energy ($\not{E}_T$) and at least two jets in the final state. The excess of events selected in the data over the background expectation from the other known Standard Model sources is taken as a measurement of the production of $t\bar{t}$ events.

This measurement provides a test of the QCD calculations of the $t\bar{t}$ cross section [1] in a channel which is independent and complementary to other measurements of the $t\bar{t}$ cross section using higher statistics final states in which at least one W boson from the top quark is reconstructed via its hadronic decay, $W \rightarrow q\bar{q}$. It is also the only final state with a favorable signal to background ratio even before requiring the identification of one of the jets in the final state as a jet from a b quark.

The CDF detector is described in detail in [2].

II. DATA SAMPLE & EVENT SELECTION

This analysis is based on an integrated luminosity of 1.2 fb^{-1} collected with the CDFII detector between March 2002 and February 2006. The data are collected with an inclusive lepton trigger that requires an electron or muon with E_T >18 GeV (P_T >18 GeV/c for the muon). From this inclusive lepton dataset, events with an offline reconstructed isolated electron of E_T, or muon with P_T, greater than 20 GeV are selected. A second electron of E_T, of muon of P_T, greater than 20 GeV is also required using looser identification cuts and no requirement on isolation. Events with more than two leptons in the final state are rejected.

This "dilepton" dataset is cleaned of other known Standard Model decays with two leptons in the final states by requiring $\not{E}_T$ >25 GeV (or >50 GeV if any lepton or jet is closer than 20^o from the missing E_T direction) and high missing E_T significance for ee and $\mu\mu$ events with dilepton invariant mass in the Z peak region.

At this point of the selection, events reconstructed with 0 or 1 jets of E_T >15 GeV are used as a control sample for the background estimation. The $t\bar{t}$ candidate region is obtained by requiring at least 2 jets with E_T >15 GeV, summed transverse energy H_T >200 GeV and the two leptons to be of opposite charge.

A. Total $t\bar{t}$ Acceptance

The acceptance for candidate events is measured using the PYTHIA Monte Carlo program [3] simulating $t\bar{t}$ events with an assumed M_{top}=175 GeV. The Monte Carlo selection is restricted to events with both W's from top decaying to a lepton plus neutrino, where the lepton can be any of e, μ or τ. Of the $t\bar{t}$ events with a reconstructed vertex along the z-direction inside ±60 cm of the nominal CDF detector origin (corresponding to 96.0% of full CDF luminous region), the acceptance for the candidate dilepton events is $\mathcal{A}$=0.808%.

The $t\bar{t}$ Monte Carlo prediction is corrected by taking into account any difference observed between data and Monte Carlo efficiencies for identifying high transverse energy electrons and muons. These corrections, in the form of data to Monte Carlo scale factors, are measured using the unbiased leg in Z-boson decays. Another correction comes from the efficiency of the inclusive lepton trigger which is measured in data samples selected with independent sets of triggers.

III. BACKGROUNDS

The sources of background processes considered for this selection are diboson (WW, WZ and ZZ) events, $W\gamma$ events in which the photon is misidentified as lepton, $q\bar{q} \rightarrow Z/\gamma^*$ and QCD production of W boson with multiple jets in which one jet is misidentified as lepton.

The two dominant sources of background dilepton events come from $Z/\gamma^* \rightarrow ee/\mu\mu$ with fake $\not{E}_T$ and from W+jets with a fake lepton. They are estimated using data-based methods. The acceptance for the remaining backgrounds, diboson, $W\gamma$ and $Z/\gamma^* \rightarrow \tau\tau$, is based on Monte Carlo predictions.

The diboson decays are simulated with PYTHIA. Their production cross section is taken from the latest NLO MCFM version [5] and CTEQ6 [6] PDF predictions to be $\sigma_{WW} = 3.65 \pm 0.8$pb, $\sigma_{WZ} = 3.67 \pm 0.1$pb. For the ZZ events the cross section is taken to be $\sigma_{ZZ} = 3.8$pb with an uncertainty of 20%. $Z/\gamma^* \rightarrow \tau\tau$ decays are also simulated with PYTHIA. The production cross-section is taken to be $\sigma_{Z_{\tau\tau}} = 238 \pm 3$pb, multiplied by a K-Factor of 1.4. $W\gamma$ decays are simulated with Baur Monte Carlo. The production cross section is taken to be $\sigma_{ZZ} = 32 \pm 3.2$pb, multiplied by a K-Factor of 1.36. A conversion inefficiency scale factor SF=1.2 $\pm$ 0.12 is applied to the electrons of E_T <40GeV. The WW, Wγ and $Z/\gamma^* \to \tau\tau$ jet multiplicity spectra are corrected to account for discrepancies observed between data and Monte Carlo in the Z-boson decays using jet bin dependent, or Njet, scale factor.The uncertainty of the Monte Carlo based backgrounds comes from the convolution of the Monte Carlo statistics, uncertainties on the N_{jet} scale factor, lepton identification and jet energy scale (JES) correction.

The contamination from $Z/\gamma^* \to ee/\mu\mu$ decays is estimated by selecting a sample of Z boson decays with high $\not\!\!E_T$ inside the 76-106 GeV/c^2 window, after correcting for the presence of non DY/Z events. The remaining DY/Z contamination is calculated as two separate contributions, events outside the Z window and events inside the Z window, using Monte Carlo to predict the ratio of events in different kinematic regions. For this analysis we use separate data estimates for the different jet bin multiplicities. The uncertainty on this background comes mostly from the limited statistics of Z$\gamma*$ data events with high $\not\!\!E_T$ used to normalize the overall prediction, from the statistics of the Monte Carlo and from the uncertainty in the jet energy scale correction. There is a small contribution to the backgrounds from $Z/\gamma^* \to e\mu$ events that originate from the $Z/\gamma^* \to \mu\mu$ process where the one electron is associated with photon conversion and is identified as electron. These events are predicted using $Z/\gamma^* \to \mu\mu$ Monte Carlo sample.

The background from fake lepton source is calculated by using a large sample of generic jets triggered by the presence of at least one jet with E_T >50 GeV. This sample was used to calculate the probability, or fake rate, that a generic electron-like object passes all of the identification cuts used to select a high E_T electron, or P_T muon. This probability is parameterized in terms of lepton transverse energy and isolation, and applied to events with only one high transverse energy reconstructed lepton plus a second electron-like or muon-like object. These events are required to pass all of the candidate events selection cuts treating the electron or muon-like objects as the second lepton in the event. The uncertainty for the fake background is dominated by the differences observed between fake rates calculated in the jet sample triggered by at least one jet with E_T >50 GeV and similar samples requiring at least one jet with E_T >20, 70 and 100 GeV.

IV. SYSTEMATIC UNCERTAINTIES

A common systematics to signal and background Monte Carlo estimates comes from the uncertainty of the lepton identification scale factors, measured in Z events which have a limited jet activity. The systematic associated to this source is conservatively taken to be 2%. Another common systematics is related to the jet energy uncertainties. This is estimated by varying the jet corrections $\pm 1\sigma$ of their systematic uncertainty and measuring the shift of the acceptance. MC-based backgrounds have uncertainty because of the N_{jet} scale factor. These three sources are considered as correlated systematic uncertainties.

Uncorrelated sources of systematic uncertainties are the jet fake systematics, the cross section uncertainties and a 30% systematic uncertainty on the conversion rejection scale factor. For the signal acceptance the systematic uncertainties are due to multiple effects: MC generator, ISF/FSR variation and PDF's uncertainty. The first two components are calculated by comparing the raw Monte Carlo acceptance of the default $t\bar{t}$ PYTHIA sample to specialized Monte Carlo samples. Table I summarizes the systematic uncertainties that affect the $t\bar{t}$ acceptance.

Source	Systematic Error (%)
MC Generator	1.5
ISR	1.7
FSR	1.1
PDF's	0.8
Jet Corrections	3.2
Total	4.2

TABLE I: Summary of systematic uncertainties affecting the $t\bar{t}$ acceptance. The total error is the sum in quadrature of each contribution

V. RESULTS

Table II shows a summary of the background estimates for each jet bin after all cuts but before the H_T and Opposite Charge requirements are applied and in the 2 jet bin after applying only the H_T cut. The last column contains the candidate events with all cuts applied. This table also show the total background expectation for the cross section of the 6.7pb, their sum (labelled as "Total SM expectation") and the number of candidate events in 1.2fb^{-1} of data.

Events per 1200 pb^{-1} vs Njet bins					
Source	0j	1j	$\geq$ 2j	H$_T$	H$_T$, OS
WW	60.82±5.69	16.77±1.87	6.60±1.15	4.37±0.73	4.19±0.71
WZ	5.02±0.44	5.02±0.37	1.91±0.27	1.45±0.23	0.95±0.15
ZZ	3.77±2.92	1.75±1.35	0.88±0.68	0.76±0.60	0.65±0.51
Wγ	11.82±3.02	3.35±0.94	0.97±0.37	0.11±0.11	0.11±0.11
DY$\rightarrow \tau\tau$	1.54±0.36	7.01±1.18	5.81±1.98	2.97±0.89	2.97±0.89
DY$\rightarrow ee + \mu\mu$	17.99±4.29	8.85±2.31	10.72±3.46	7.89±1.60	7.89±2.14
Fakes	38.38±11.33	25.80±8.16	17.98±6.08	12.76±4.91	8.80±3.86
Total background	139.34±17.96	68.55±10.74	44.87±9.49	30.30±5.89	25.56±5.54
$t\bar{t}$ (σ = 6.7 pb)	0.29±0.03	7.49±0.58	59.53±4.53	57.41±4.37	55.95±4.26
Total SM expectation	139.63±17.98	76.04±11.10	104.40±13.17	87.71±8.85	81.52±8.92
GEN6 DATA	143	84	114	88	77

TABLE II: Summary table of background estimates, $\mathrm{t\bar{t}}$ predictions and events in 1.2 fb^{-1} of data for each jet bin after all cuts but before the $\mathrm{H_T}$ and Opposite Charge requirements are applied and in the 2 jet bin after applying only the $\mathrm{H_T}$ cut. The last column contains the candidate events with all cuts applied. The quoted uncertainties are the sum of the statistical and systematics uncertainty.

Table III shows the total number of background, Standard Model expectation and 1.2fb^{-1} data candidate events, divided by lepton flavor contribution.

Events per 1200 pb^{-1} after all cuts				
Source	ee	$\mu\mu$	eμ	$\ell\ell$
WW	1.01±0.19	1.04±0.19	2.14±0.37	4.19±0.71
WZ	0.40±0.07	0.30±0.05	0.25±0.05	0.95±0.15
ZZ	0.28±0.22	0.28±0.22	0.10±0.08	0.65±0.51
Wγ	0.11±0.11	0.00±0.00	0.00±0.00	0.11±0.11
DY$\rightarrow \tau\tau$	0.48±0.23	0.87±0.33	1.61±0.54	2.97±0.89
DY$\rightarrow ee + \mu\mu$	2.46±0.88	4.89±1.25	0.54±0.27	7.89±2.14
Fakes	2.12±1.25	3.10±1.36	3.58±2.02	8.80±3.86
Total background	6.86±1.70	10.47±2.06	8.23±2.30	25.56±5.54
$t\bar{t}$ (σ = 6.7 pb)	12.18±0.94	13.60±1.04	30.17±2.30	55.95±4.26
Total SM expectation	19.04±2.26	24.08±2.68	38.40±3.90	81.52±8.92
GEN6 DATA	16	26	35	77

TABLE III: Summary table by lepton flavor content of background estimates, $\mathrm{t\bar{t}}$ predictions and final candidate events in 1.2fb^{-1} of data.The quoted uncertainties is the sum of the statistical and systematics uncertainty.

The cross section is calculated as:

$$\sigma_{\mathrm{t\bar{t}}} = \frac{\mathrm{N_{obs}} - \mathrm{N_{bkg}}}{\sum_{\mathrm{i}} \mathcal{A}_{\mathrm{i}} \mathrm{x} \mathcal{L}_{\mathrm{i}}} \tag{1}$$

where $\mathrm{N_{obs}}$ is the number of dilepton candidate events, $\mathrm{N_{bkg}}$ is the total background and the dominator is the weighted sum of the corrected acceptance for each dilepton category $\mathcal{A}_i$ multiplied by the luminosity relative to that category $\mathcal{L}_i$. Different luminosity is used as the single leptons in a given category require CDF subdetectors to be fully functional. The total denominator is 8.351 ± 0.047 pb.

For $\mathrm{t\bar{t}}$ events in the dilepton channel, we find a cross section of:

$$\sigma_{\mathrm{t\bar{t}}} = 6.16 \pm 1.05_{\mathrm{stat}} \pm 0.72_{\mathrm{syst}} \pm 0.37_{\mathrm{lumi}} \text{ pb},$$

where the first uncertainty is statistical, the second is the convolution of the acceptance and background systematics and the third comes from the 6% uncertainty in the luminosity measurement.

Figure 1 shows the number of candidate events in 0, 1, $\geq$2jet events together with a histogram representing the component of the background. The yellow band gives the $\mathrm{t\bar{t}}$ contribution for a cross section of 6.7 pb. The red hatched area is the uncertainty in the total background estimate.

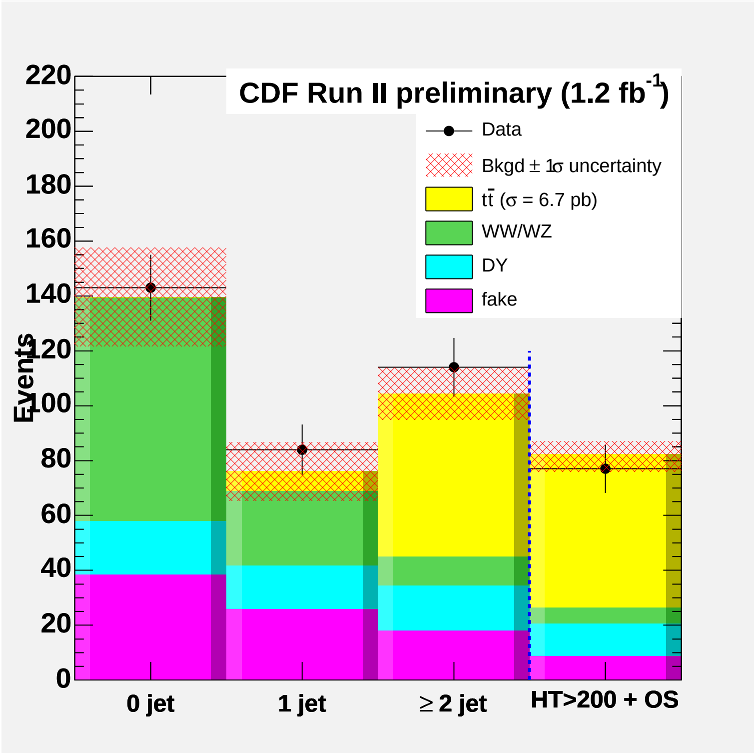

FIG. 1: Dilepton candidate events (black point) by jet multiplicity. The colored histogram represents the background contribution for an assumed $\sigma_{t\bar{t}} = 6.7$pb. The red hatched area is the uncertainty in the total background estimate.

Acknowledgments

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium fuer Bildung und Forschung, Germany; the Korean Science and Engineering Foundation and the Korean Research Foundation; the Particle Physics and Astronomy Research Council and the Royal Society, UK; the Russian Foundation for Basic Research; the Comision Interministerial de Ciencia y Tecnologia, Spain; and in part by the European Community's Human Potential Programme under contract HPRN-CT-20002, Probe for New Physics.

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