

Search for New Physics in the Exclusive Delayed Photon + MET Final State at CDF

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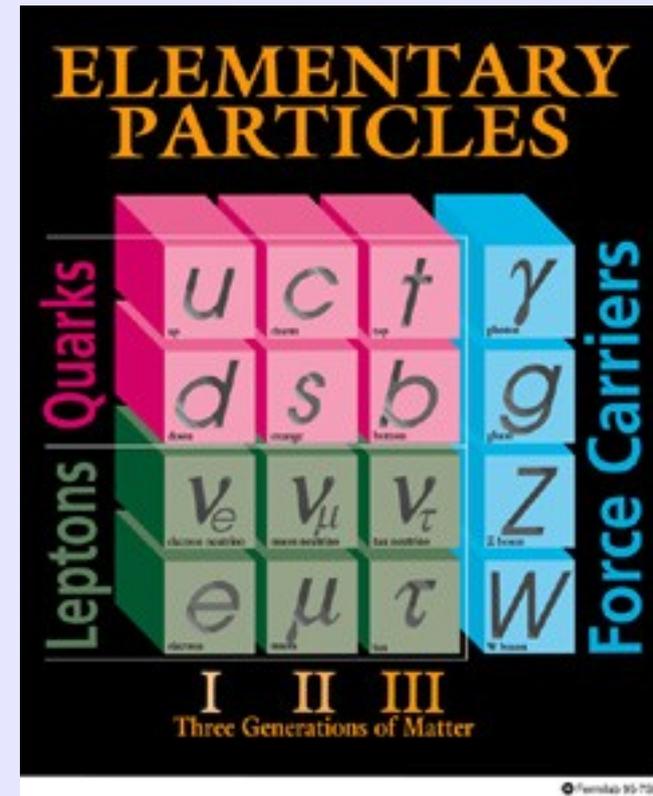


Outline

- Introduction
- Motivation
- Tools
- Overview of the Delayed Photon Analysis
- Backgrounds with Large Times and Cuts to Get Rid of Them
- Background Estimation
- Results
- Conclusions

Standard Model

- The Standard Model (SM) describes all currently known particles and interactions
- Decades of experimental verification have confirmed many of its predictions
- Despite extraordinary success, the Standard Model has problems
 - The “hierarchy problem” - the Higgs mass has divergences that must be canceled with fine tuning
 - Dark matter and dark energy make up a substantial portion of the universe



Supersymmetry

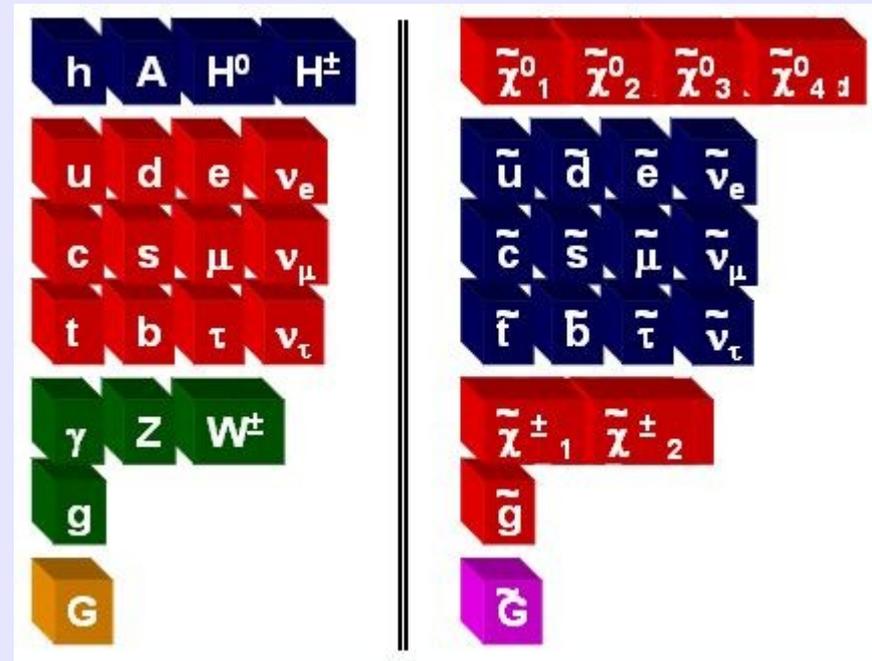
Supersymmetry (SUSY) proposes a symmetry between fermions and bosons – roughly doubles the particle count

The new particles cancel the divergence in the Higgs mass

If “R-parity” is conserved, SUSY could provide a dark matter candidate

This isn't an exact symmetry → SUSY particles must be heavy

Various breaking mechanisms lead to different phenomenology



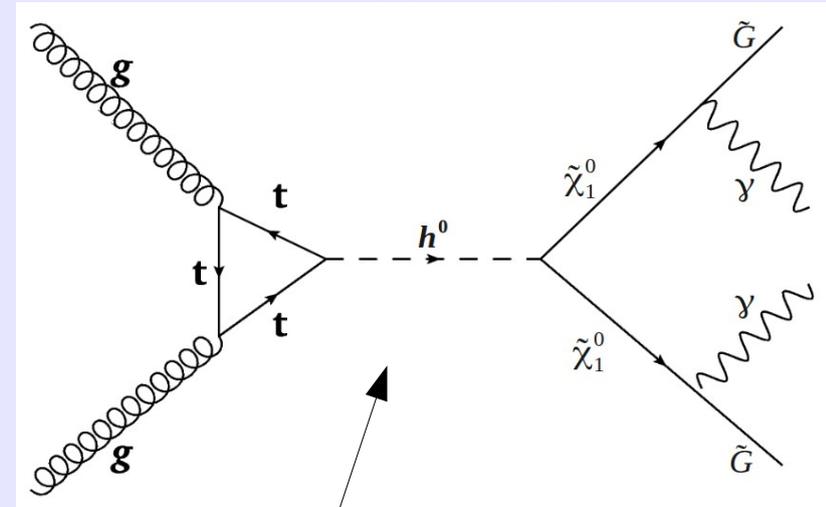
Search for SUSY decays of the Higgs in Gauge Mediated Supersymmetry Breaking Models

In GMSB, the \tilde{G} , the SUSY partner of the graviton, is typically the lightest supersymmetric particle (LSP)

In general GMSB models, it is possible that only the $\tilde{\chi}_1^0$ and \tilde{G} are accessible at the Tevatron

These models are not constrained by current limits \rightarrow worth going after!

The NLSP, $\tilde{\chi}_1^0$, is often long-lived. This is favored in low-scale SUSY breaking models. We look at cases where it has a lifetime of a few nanoseconds



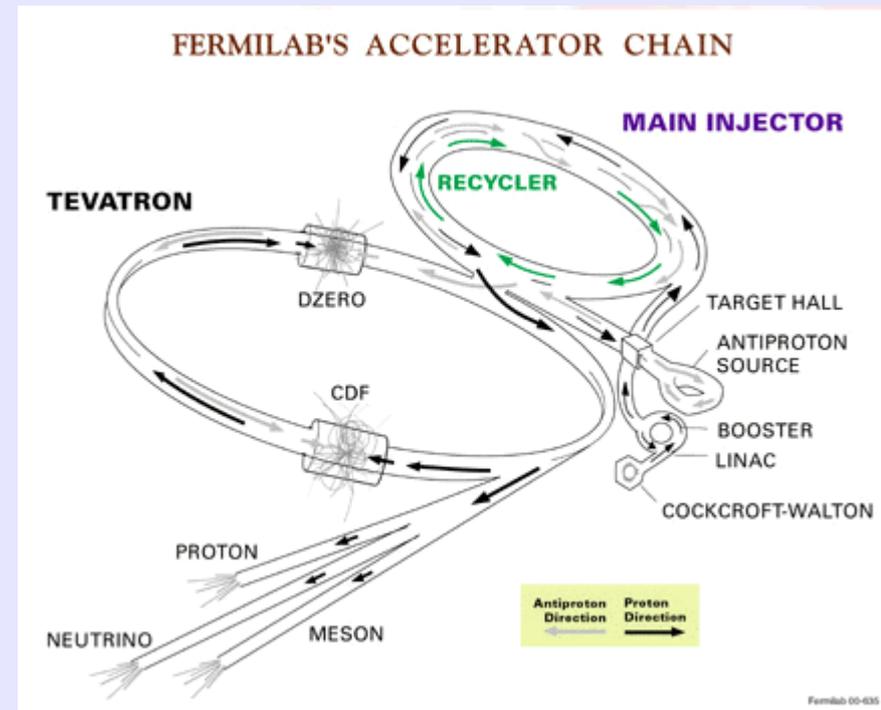
Observation of Higgs at 125 GeV means this should be visible at the Tevatron

Long lifetime means only one $\tilde{\chi}_1^0$ decays in the detector, leading to the exclusive $\gamma + \cancel{E}_T$ final state

Tevatron

The Tevatron, with a center of mass energy of 1.96 TeV, was the highest energy accelerator in the world. It collided protons with anti-protons every 396 ns.

Two detectors, CDF and D0 each collected nearly 10 fb^{-1} of data.



Collider Detector at Fermilab (CDF)

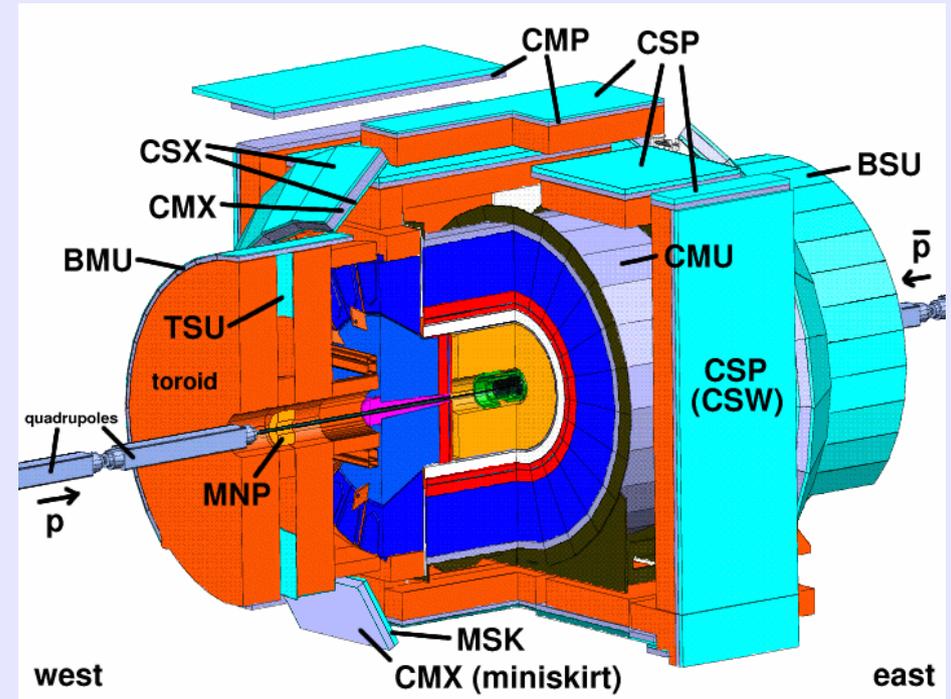
CDF is one of two multi-purpose detectors built to study collisions at the Tevatron.

Components heavily used in this analysis:

Central outer tracker – records the path taken by charged particles.

Electromagnetic calorimeter - records energy deposits from particles that interact electromagnetically

EMTiming system – converts output of the EM calorimeter into the time of arrival of the incident particle. In the central region, it is fully efficient for energies > 6 GeV (resolution ~ 0.6 ns)



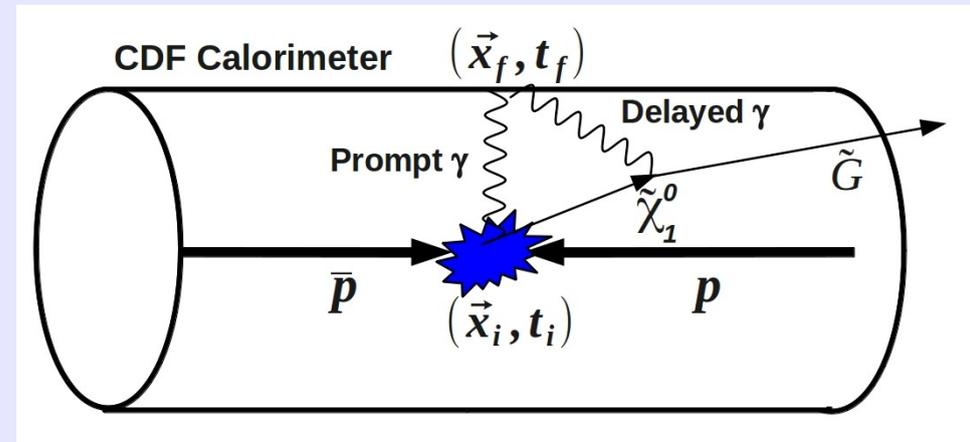
Delayed Photons

Photons from long-lived $\tilde{\chi}_1^0$ arrive at the calorimeter late compared to expectations from prompt photons (“delayed photons”).

This gives provides a distinct search signature.

Our primary analysis variable is the time of arrival of the photon at the EM calorimeter minus the expected time of arrival.

$$t_{corr} = t_f - t_i - \frac{|\vec{x}_f - \vec{x}_i|}{c}$$



A preliminary, internal analysis found a large excess of delayed photons in this final state; therefore, the goals of this analysis are:

- 1) Look for previously unknown biases
- 2) Reexamine the background model

Final State and Backgrounds

Final State

Require

(all E_T relative to $Z = 0$)

- Photon with $E_T > 45$ GeV
- MET > 45 GeV
- At least one space-time vertex with $|Z| < 60$ cm

Veto

- Extra calorimeter clusters with $E_T > 15$ GeV
- Tracks with $P_T > 10$ GeV
- Tracks close to the photon
- Vertices with $|Z| > 60$ cm
- Additional cosmics and beam halo cuts

Backgrounds

Standard Model Sources

$$W \rightarrow e\nu \rightarrow \gamma_{fake} + \cancel{E}_T$$

$$\gamma + jet \rightarrow \gamma + jet_{lost} \rightarrow \gamma + \cancel{E}_T_{fake}$$

$$W \rightarrow \tau\nu \rightarrow \gamma_{fake} + \cancel{E}_T$$

$$W\gamma \rightarrow l\nu\gamma \rightarrow \gamma + l_{lost} + \cancel{E}_T$$

$$Z\gamma \rightarrow \nu\nu\gamma \rightarrow \gamma + \cancel{E}_T$$

Non-Collision

- Cosmics
- Beam Halo

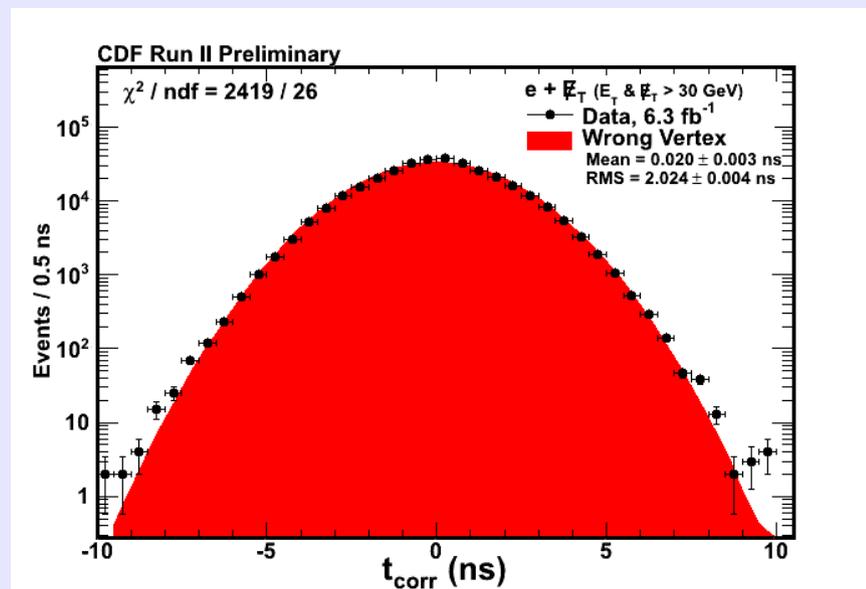
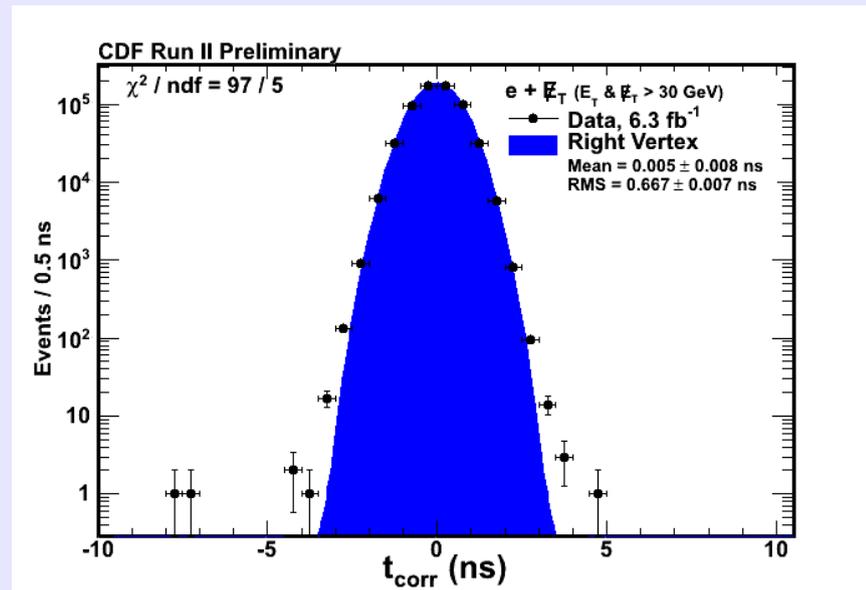
Right and Wrong Vertices

To construct the corrected time, we pick the highest ΣP_T vertex.

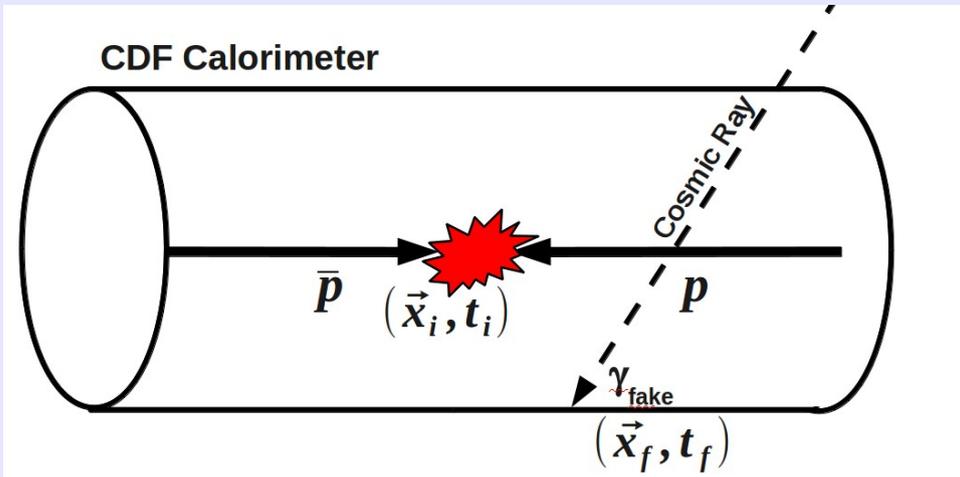
If this vertex is the origin of the particle that created the deposit in the calorimeter, it is a **Right Vertex** event - by definition, the mean is zero but with an RMS of ~ 0.64 ns.

There often multiple vertices per event. Sometimes the wrong vertex has a higher ΣP_T than the right vertex, and sometimes the right vertex is not reconstructed at all.

Wrong vertex events have an RMS of ~ 2 ns and generally a non-zero mean.

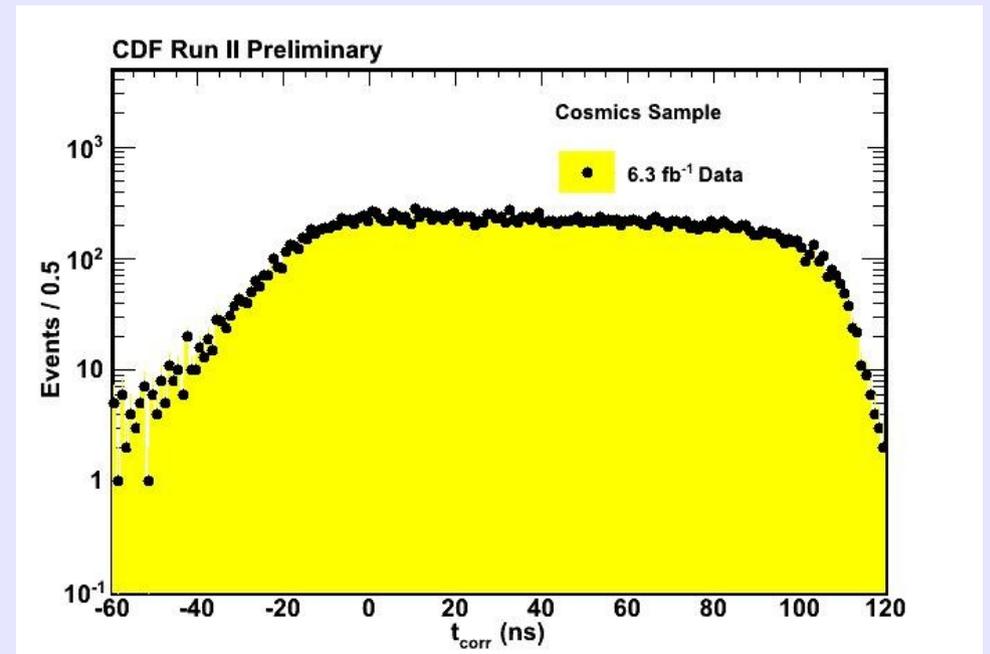


Cosmic Rays

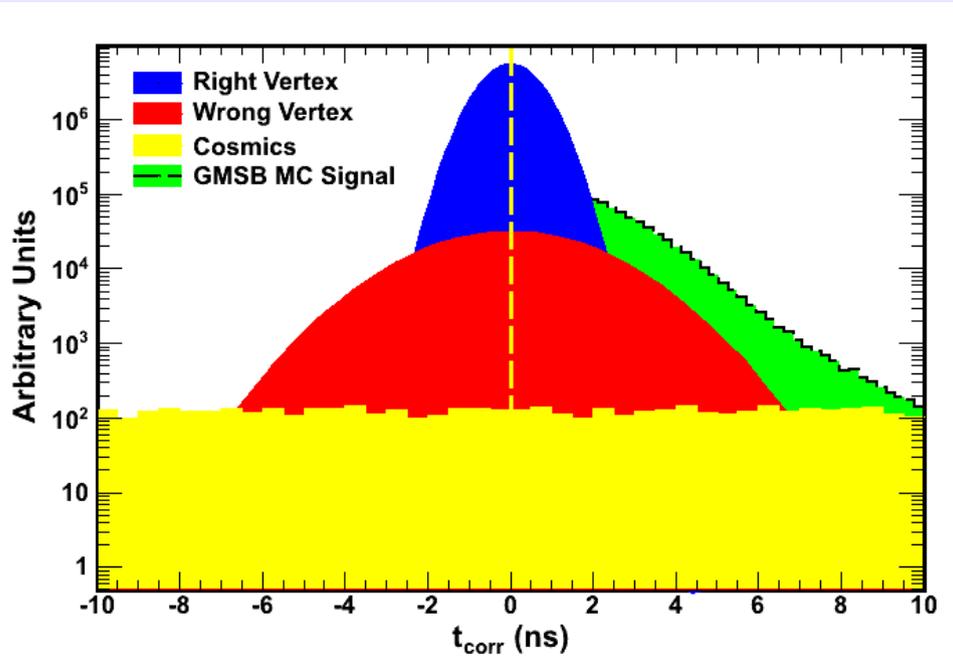


Cosmic rays occasionally reach the detector and leave an energy deposit which is reconstructed as a photon

This is uncorrelated with the bunch structure of the beam, so the rate of recording such events is flat in time, except near the opening and closing of the energy integration window



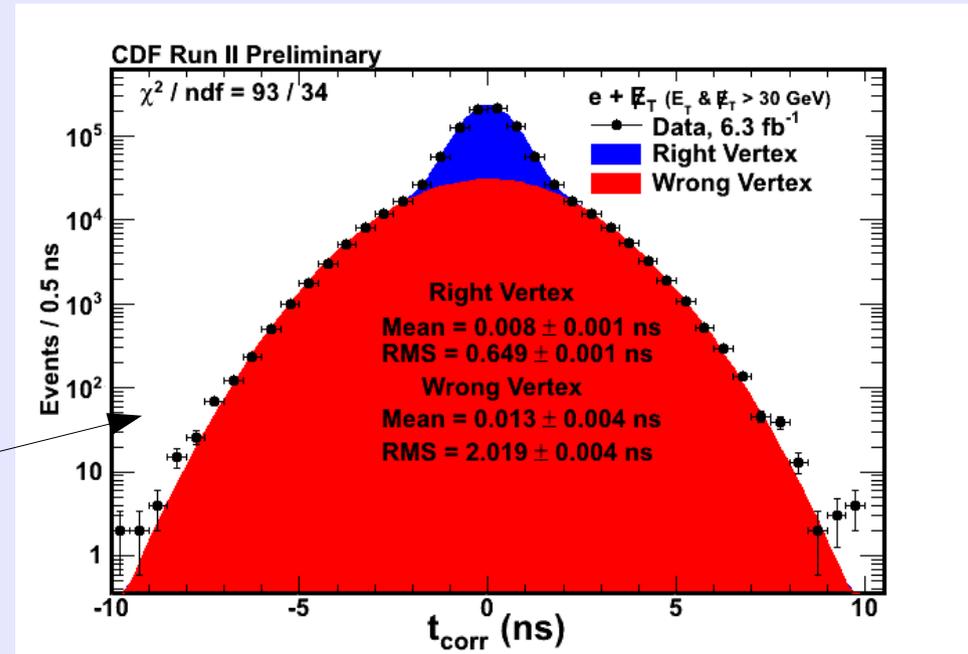
Timing Distributions



The distribution of photons from GMSB decays are expected to be a decaying exponential smeared by the detector resolution

$W \rightarrow e\nu$ where we ignore the track for the purposes of selecting a vertex acts as a control region for $\gamma + \cancel{E}_T$

Real collision data with electrons is well modeled by a double Gaussian description



Sources of Large Times from SM Backgrounds

The following effects can bias the timing distribution:

1) E_T Threshold Effect:

A distortion caused by events entering or leaving our sample due to mis-measured E_T near the cut.

Topology Biases:

2) Fake photons: Fake photons tend to be biased to larger times due to being more likely at large path lengths.

3) Lost jets: Losing an object tends to happen at more extreme vertex Z positions (to allow the object to point out of the detector).

Sources of Large Time Events:

1) E_T Threshold Effect

Promotion Effect

Wrong vertex gives shorter apparent path length

- Longer apparent time
- Larger measured E_T

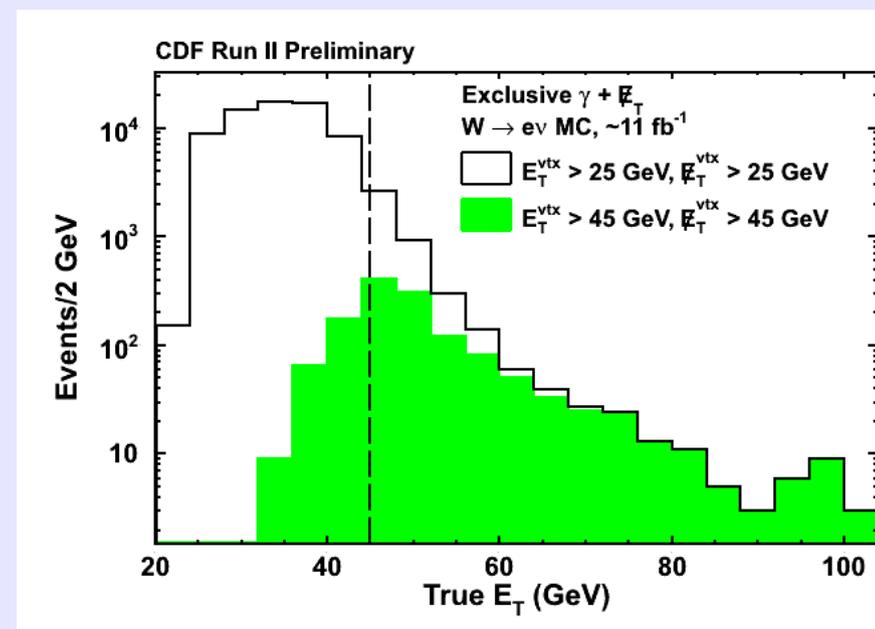
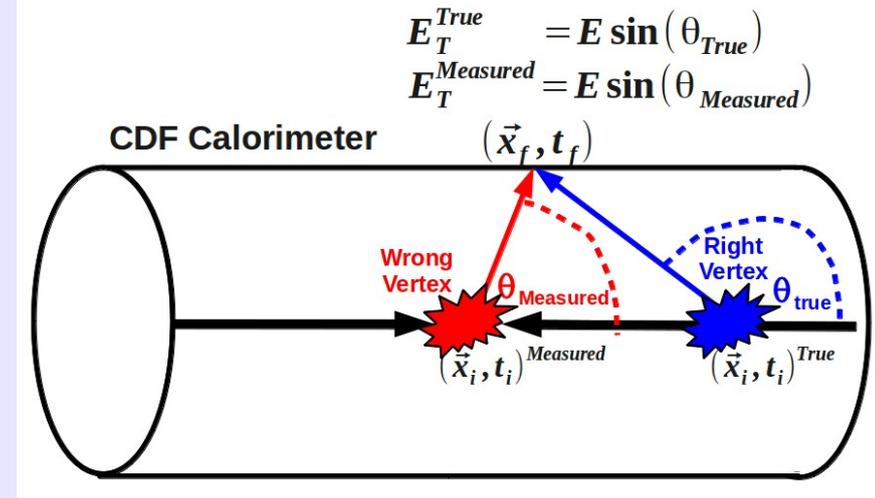
Events below the E_T threshold enter the sample and **increase** the positive time bias.

Demotion Effect

Wrong vertex gives larger apparent path length

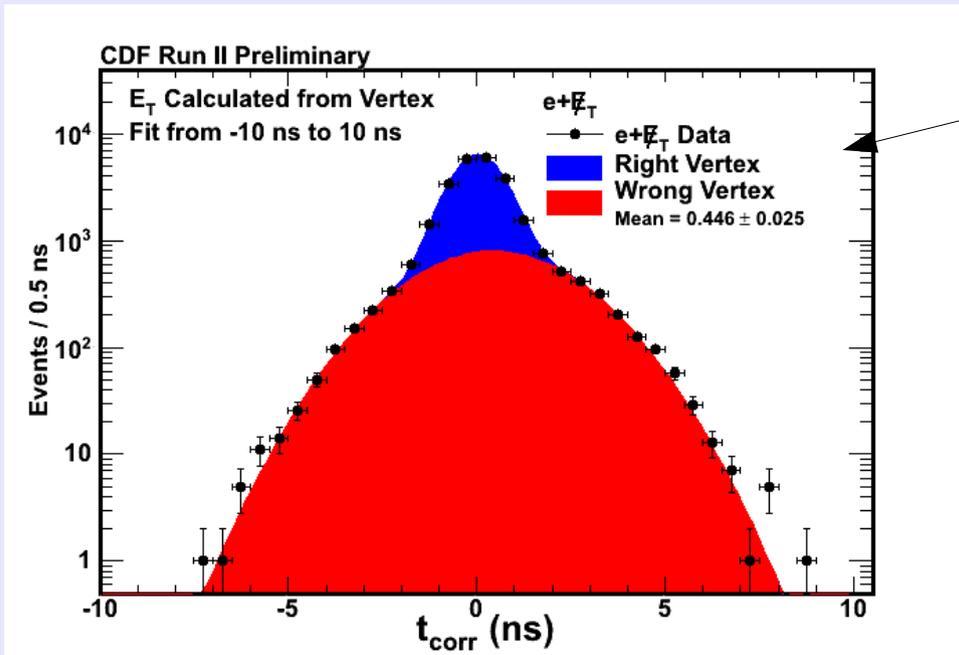
- Shorter apparent time
- Smaller measured E_T

Events above the E_T threshold exit the sample and **decrease** the negative time bias.

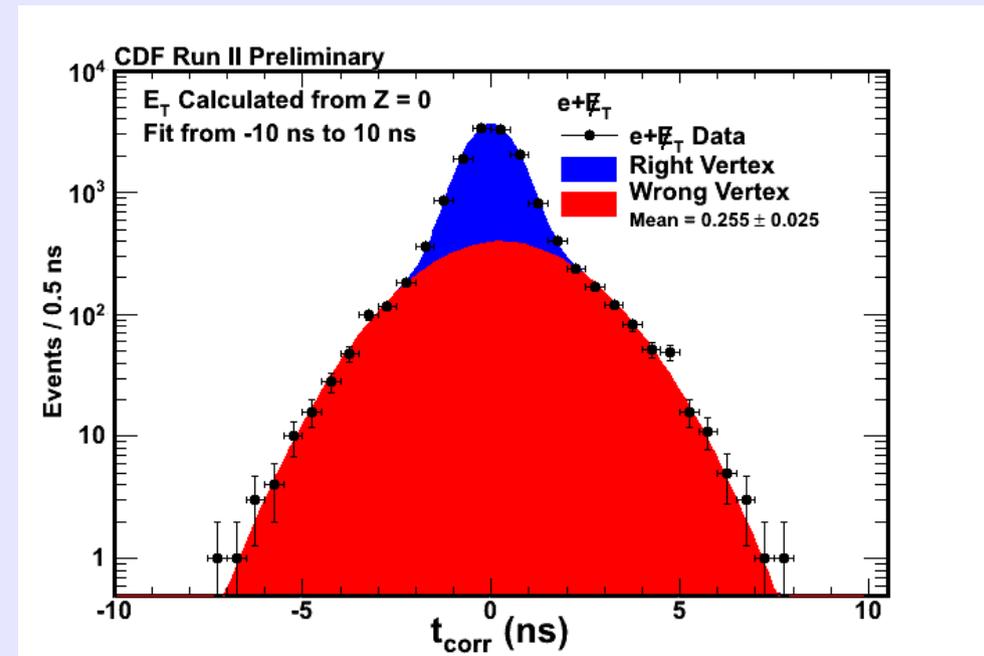


1) Solution: E_T^0 Cut

Decouple the timing measurement from the E_T measurement by calculating E_T relative to $Z = 0$

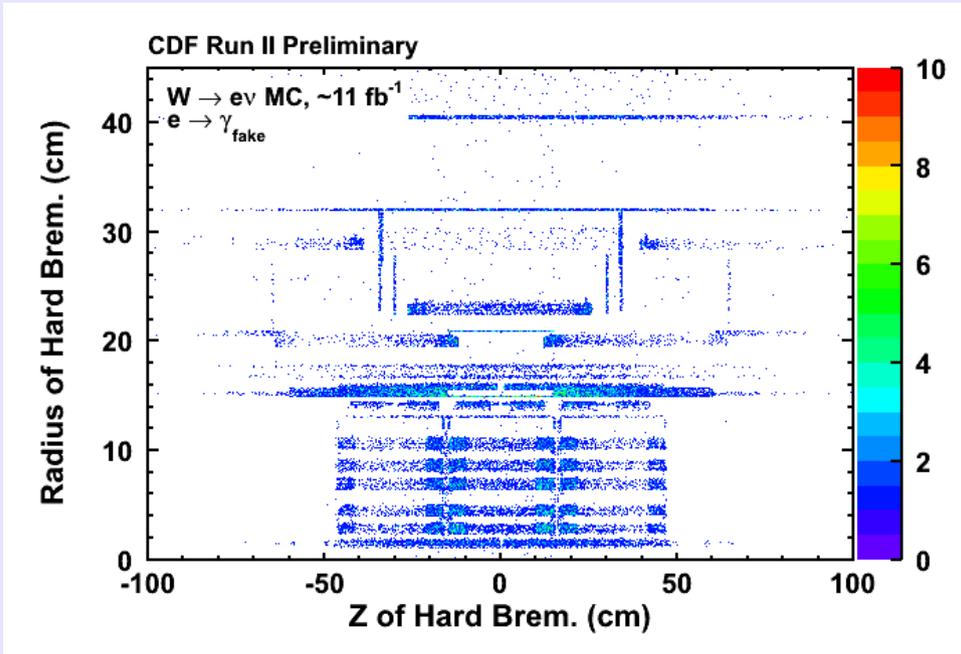


Real data with electrons using E_T relative to the selected vertex



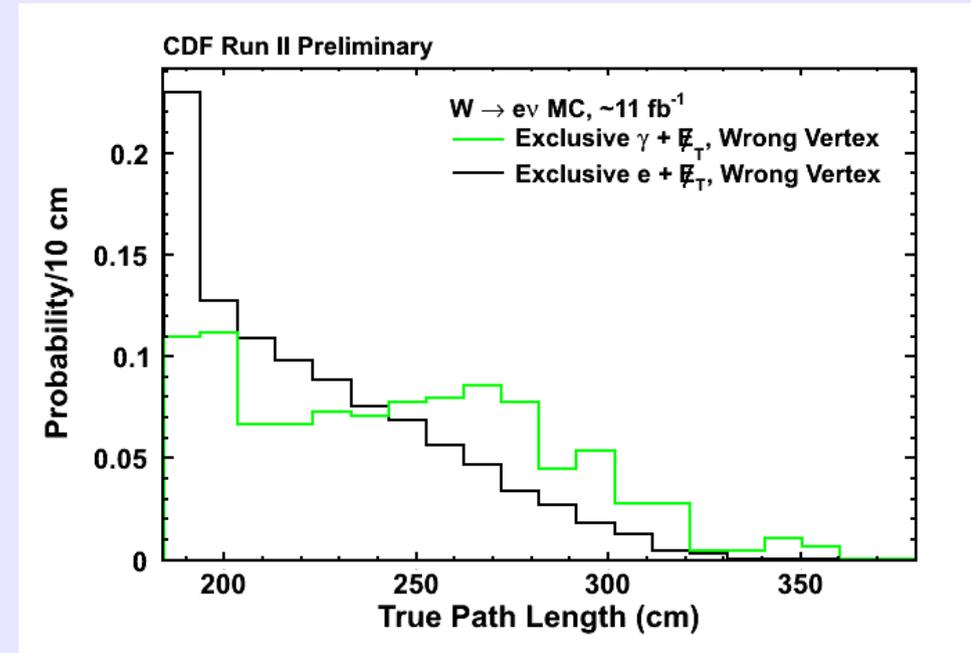
The same data using $E_T^0 \rightarrow$ the wrong-vertex mean decreases by \sim half!

Sources of Large Time Events: 2) Fake Photons



Most electrons that fake photons are due to hard interactions with detector material

This makes them have longer path lengths on average \rightarrow larger apparent times with a wrong vertex



2) Solution: ΔR_{pull} Cut

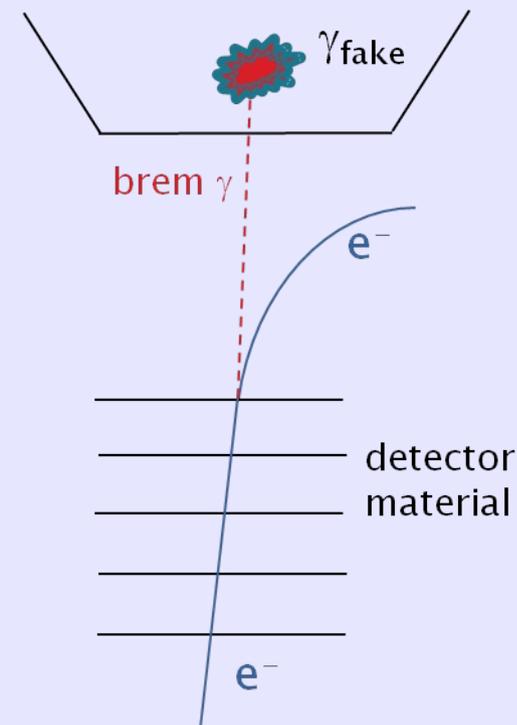
Develop a new fake rejection technique:

Electrons faking photons start off pointing towards the calorimeter deposit, but due to the hard interaction, the path has a “kink” that ruins track extrapolation

Create a ΔR between the track and the calorimeter deposit based on standardized versions of the initial η and ϕ of the track

~73% rejection of fake photons

~90% efficiency



Sources of Large Time Events:

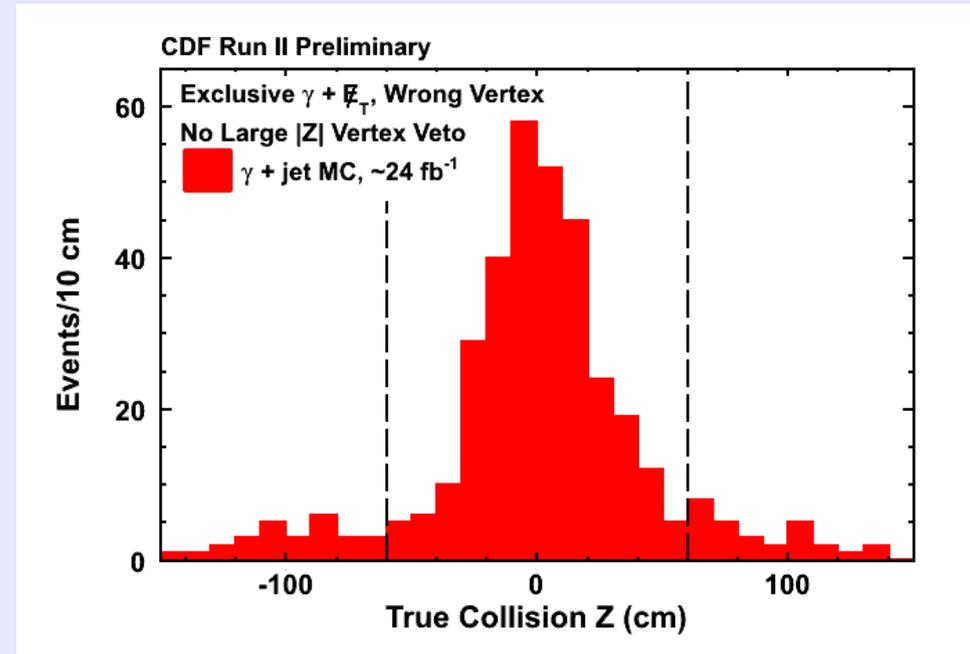
3) Large $|Z|$ Production

γ +jet events tend to occur unusually often at large $|Z|$ positions

Jets are messy objects – to lose one, it usually has to be pointed into an uninstrumented region

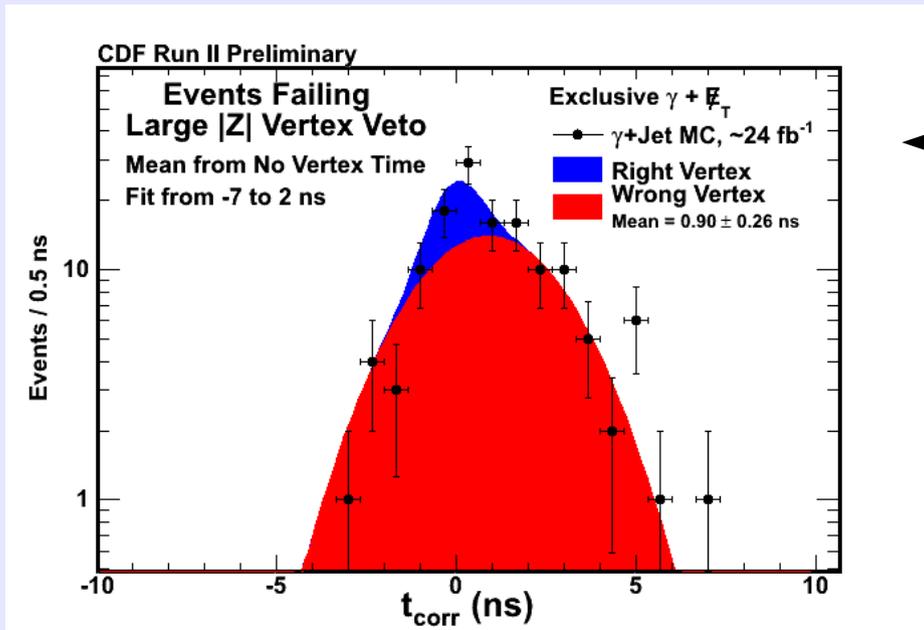
Events with large $|Z|$ are more likely to lose a jet due to it being oriented out of detector

Large $|Z|$ events have large times \rightarrow the true time-of-flight is large compared to any possible time-of-flight correction



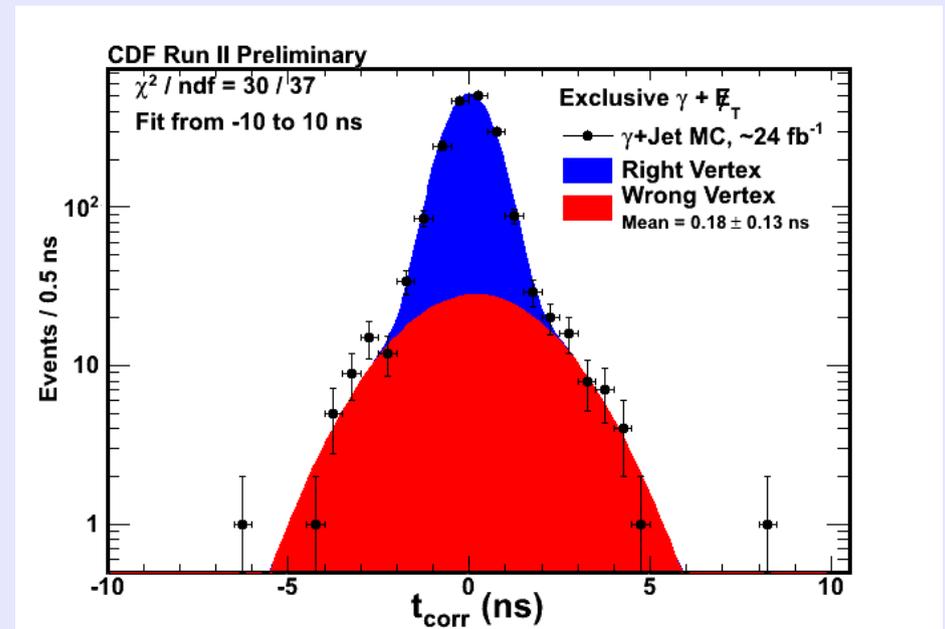
3) Solution: Large $|Z|$ Veto

Reject any event with a vertex with 3 or more tracks and $|Z| > 60$ cm (~95% efficient for right vertex events)

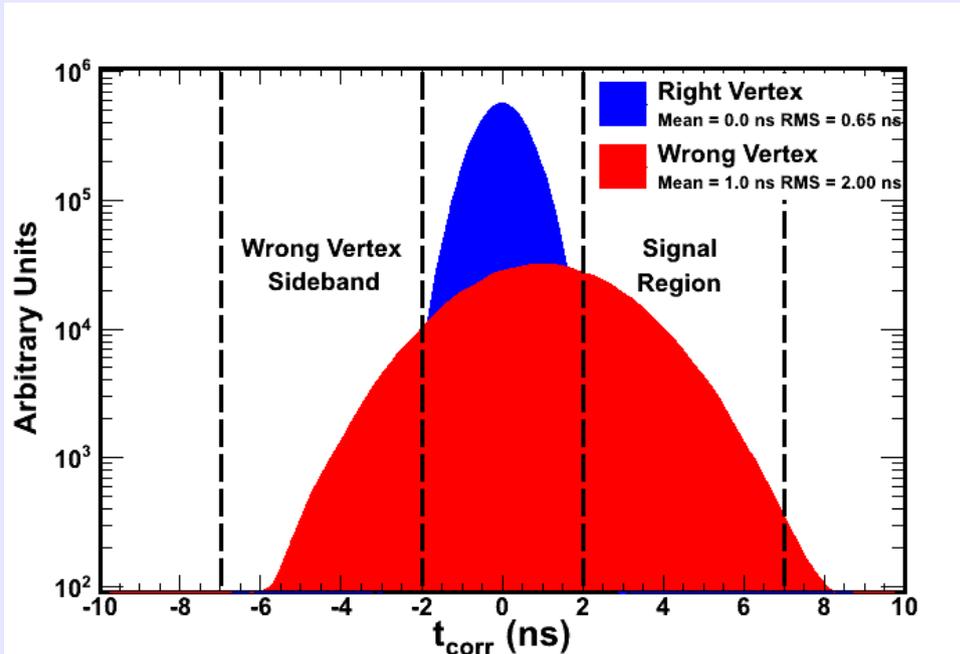


γ +jet events failing the large $|Z|$ veto are highly shifted

After the veto, the distribution is well behaved with a small wrong-vertex mean



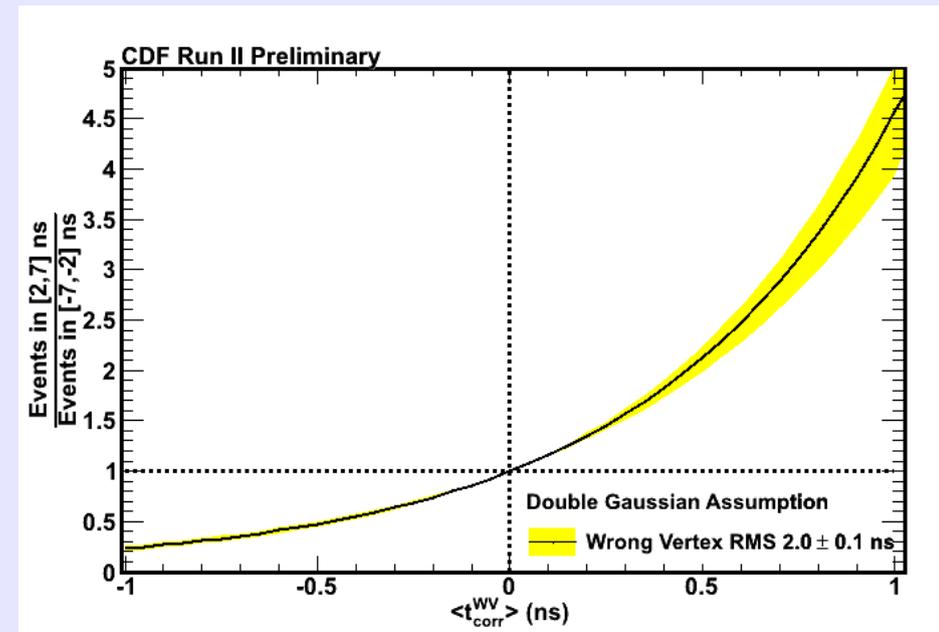
Predicting Background Events in the Signal Region From the Wrong-Vertex Mean



We want to be able to predict the number of background events in (2,7) ns using a data-driven method

Note: right-vertex events are largely irrelevant in the signal region

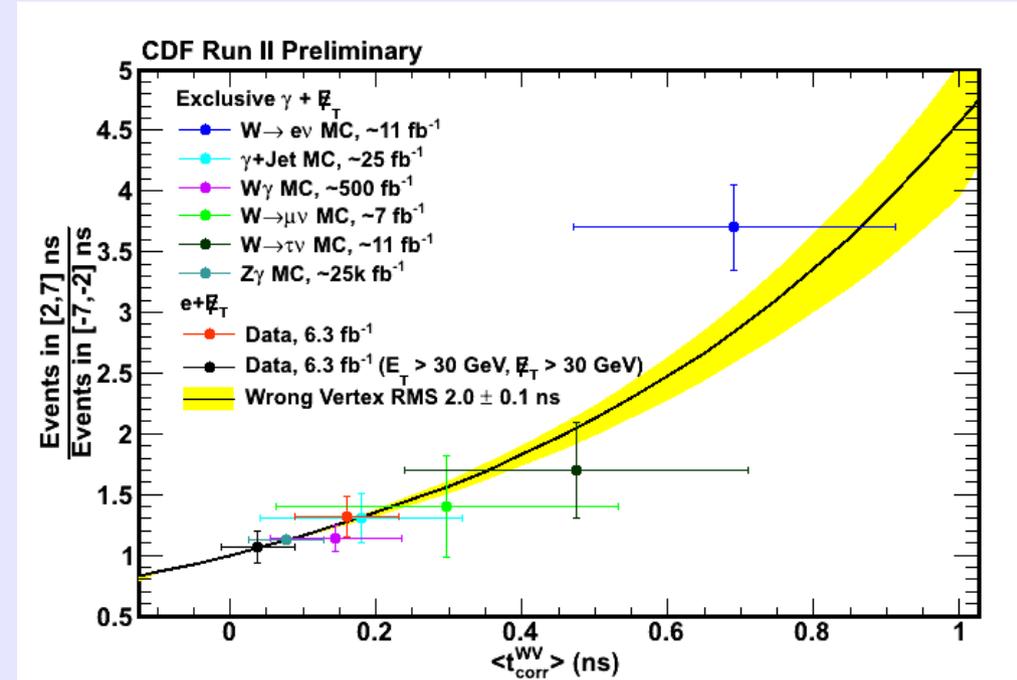
The number of wrong-vertex background events in the signal region depends directly on its normalization which we can get from (-7,-2) ns, and the wrong-vertex mean which we get from a second sample



Checking the Double Gaussian Approximation with Lots of Datasets

We isolate wrong vertex events in Monte Carlo and fit to find the wrong-vertex mean and RMS

For real data, we use electrons so we can use the electron track to identify wrong vertex events



Our data after all cuts is at ~ 0.2 ns

The ratio of events in (2,7) ns to events in (-7,-2) ns follows our predictions according to the double Gaussian approximation.
(Not a fit!)

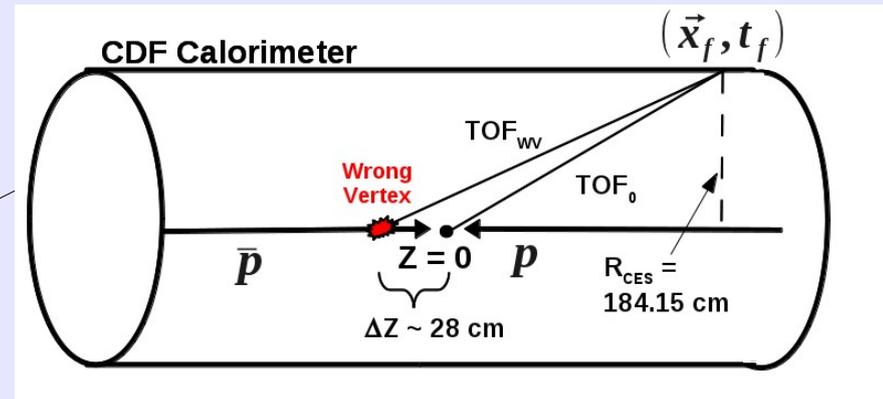
Estimating the Wrong-Vertex Mean From the No-Vertex Sample

Create orthogonal sample of events passing all cuts but good vertex requirement. Create the corrected time relative to the center of the detector: $t_{corr}^0 = t_{arrival} - 0 - TOF_0$

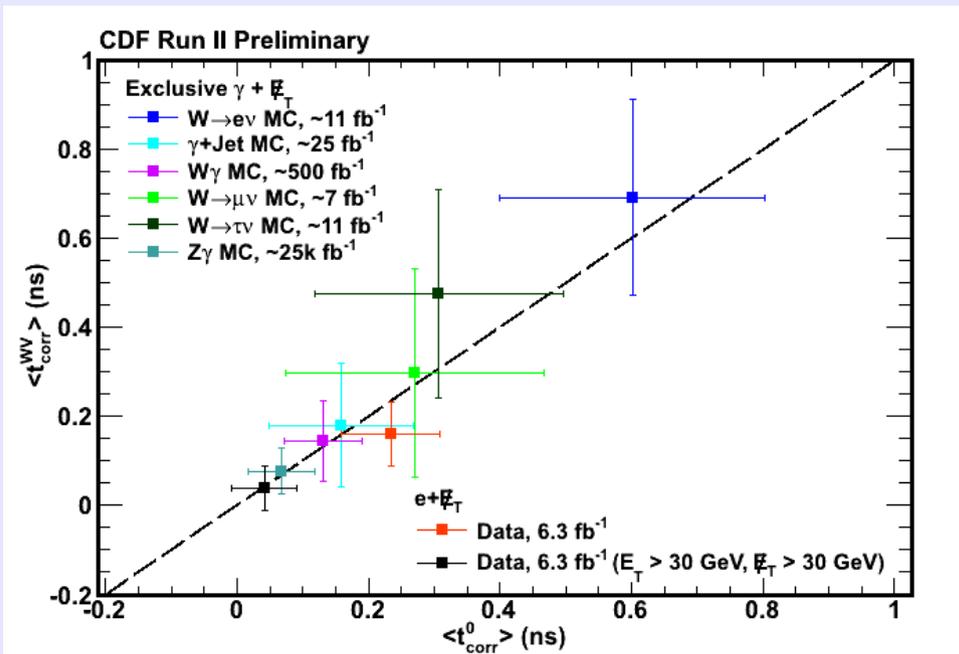
Substituting into wrong-vertex time:

$$t_{corr}^{WV} = t_{corr}^0 - t_{WV} + (TOF_0 - TOF_{WV})$$

Zero on average

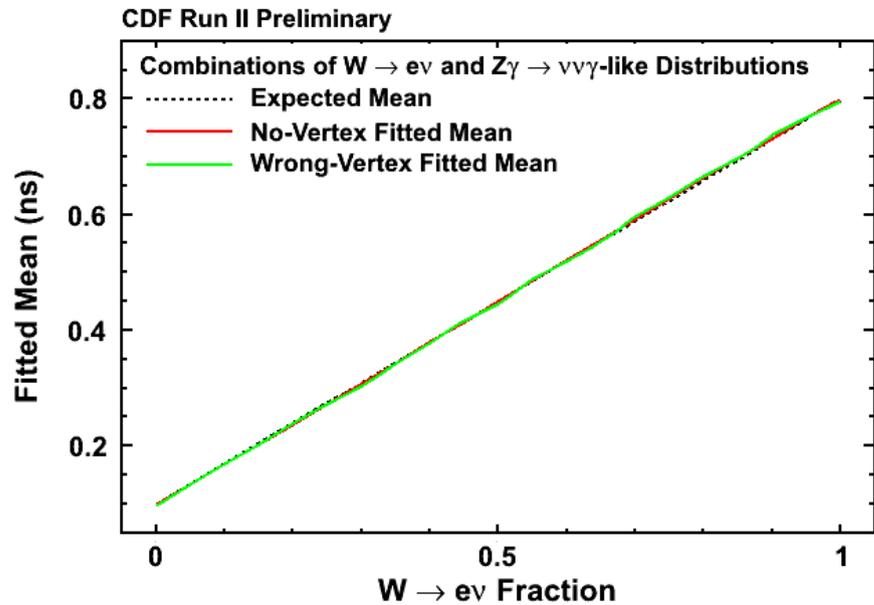


Typical $\Delta Z \ll$ than radius of detector \rightarrow average \sim zero



The mean no-vertex time is approximately equal to the mean wrong-vertex time for all control samples!

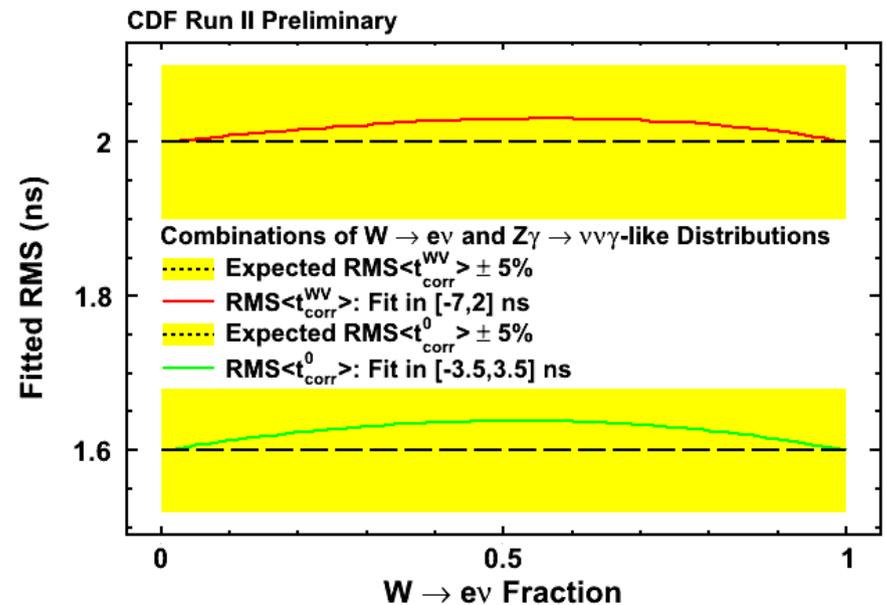
Effect of Combining Collision Background Sources



We generate Gaussians with means of 0.1 ns and 0.7 ns. We combine them in various fractions.

The fitted RMS increases slightly as we approach a 50% combination. We cover this with a 5% systematic.

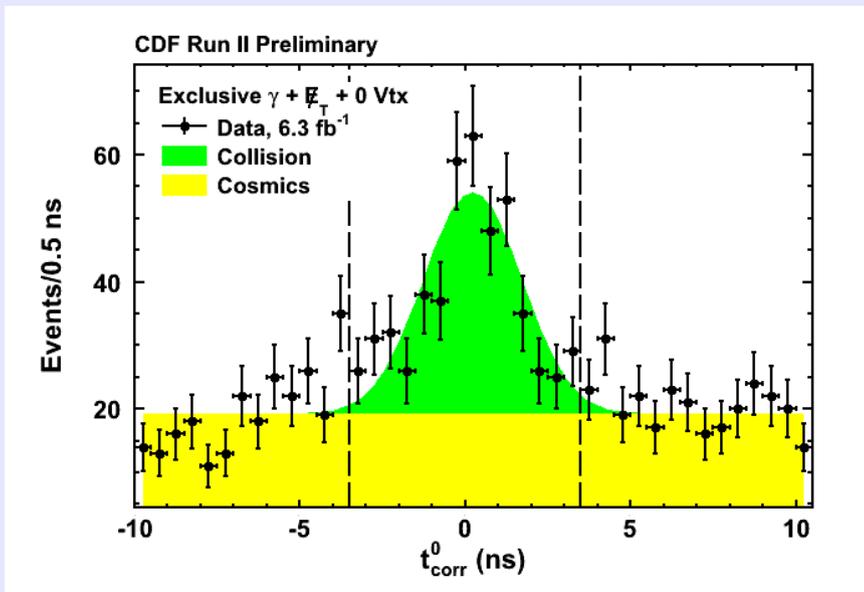
Up to this point, we considered single Standard Model sources. Does the double Gaussian description apply with combinations of sources?



Putting It All Together: Likelihood Fit

- Estimate the number of background events in the signal region using a combined likelihood fit to the sideband regions extrapolated to the signal region
 - Good vertex: (-7,2) ns and (20,80) ns
 - No vertex: (-3.5, 3.5) ns and (20,80) ns
- Include systematic uncertainties as constraint terms:
 - Right-vertex mean = 0.0 ± 0.05 ns
 - Right-vertex RMS = 0.64 ± 0.05 ns
 - Wrong-vertex mean = No-vertex mean ± 0.08 ns
 - Wrong-vertex RMS = 2.0 ± 0.1 ns

Sideband Regions



No Vertex:

Collision Events = 260 ± 30

Collision Mean = 0.2 ± 0.1

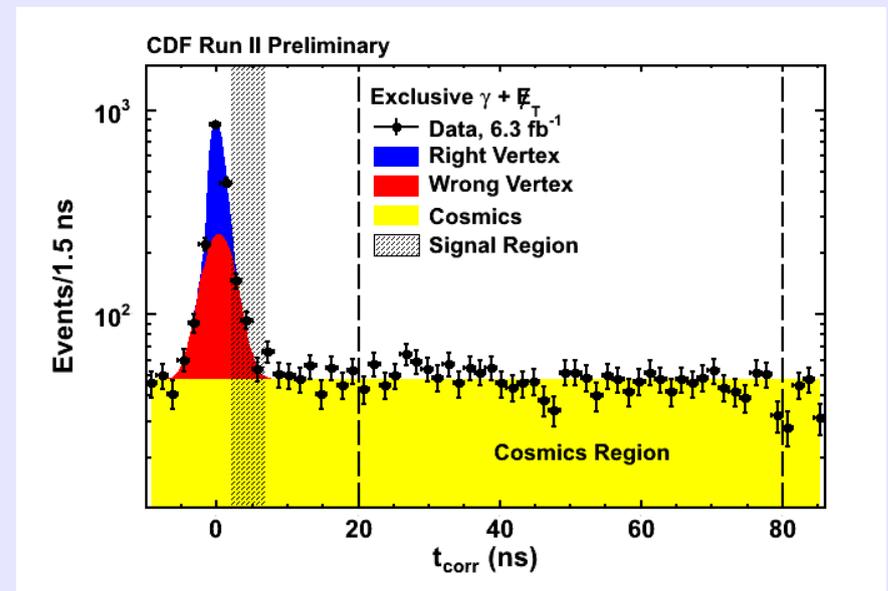
Cosmics/ns = 38.1 ± 0.8

Good Vertex:

Right-Vertex Events = 870 ± 70

Wrong-Vertex Events = 680 ± 80

Cosmics/ns = 31.9 ± 0.7

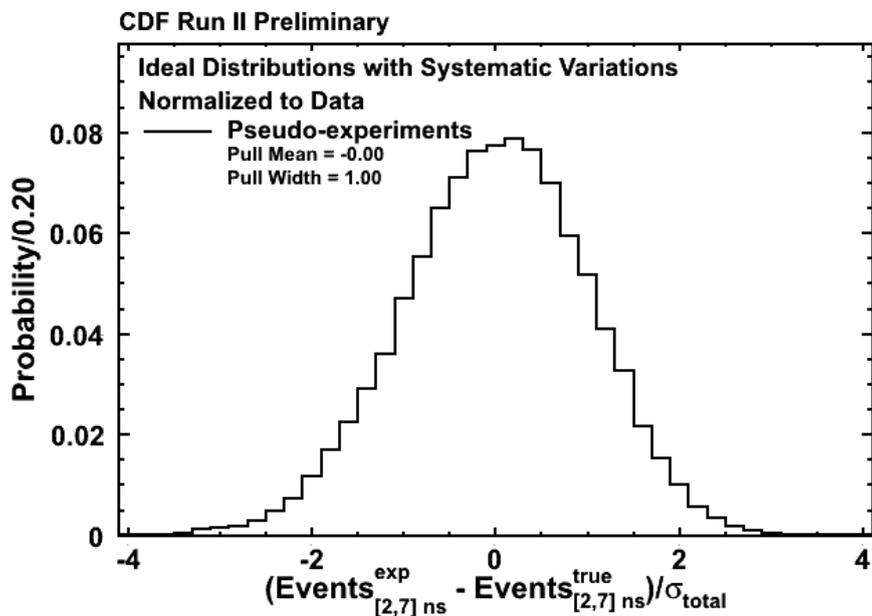


Next: use the numbers to validate the fit

Validating the Likelihood Fit

- Generate ideal pseudo-experiments varying parameters within their systematic uncertainties
- Generate more realistic pseudo-experiments from full MC of the three largest SM backgrounds
- Sample at the statistics level seen in data
- Add the expected level of cosmics to the good and no vertex distributions

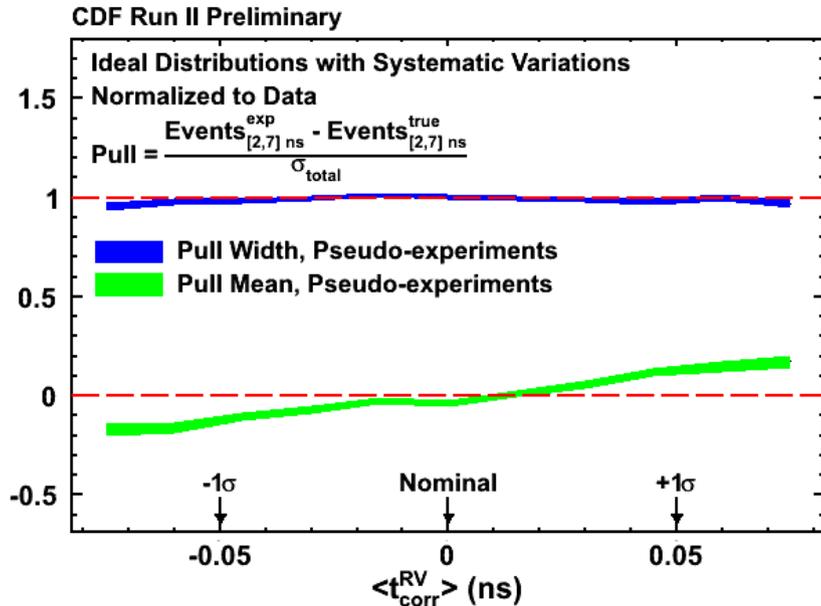
Ideal Distributions: How Well Do We Do?



All parameters with systematic uncertainties are allowed to vary within those uncertainties.

The pull distribution shows that with full variation of the systematics, the fit is unbiased (mean ~ 0) and the errors are well estimated (RMS ~ 1).

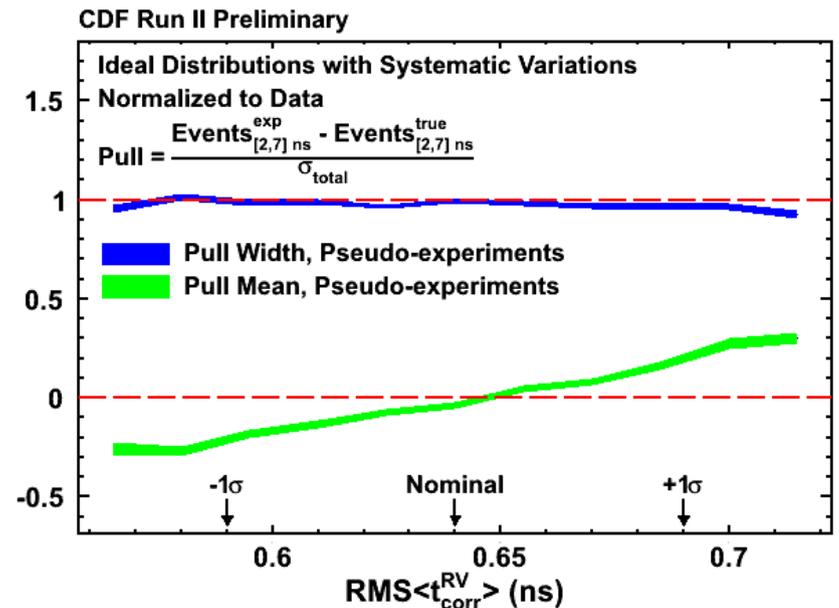
Ideal Distributions: Pulls vs. Systematic parameters



In both cases, the pull width indicates that the uncertainties are well estimated over the entire range

Figures range from -1.5σ to 1.5σ in systematic uncertainty

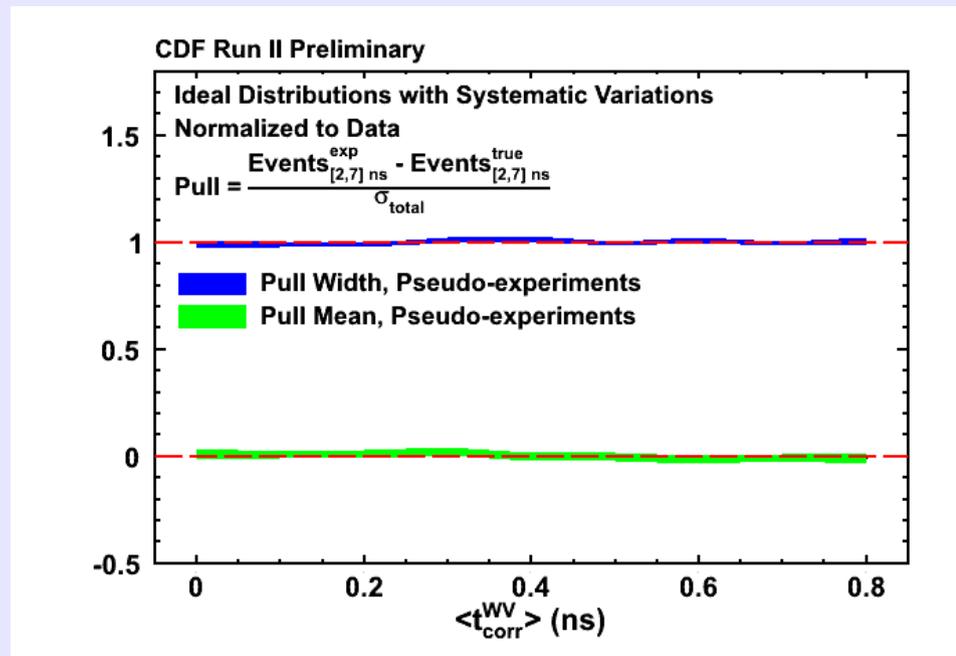
The fit remains largely unbiased over this range



How well does the fitter do for different wrong vertex means?

The wrong-vertex mean is not known a priori.

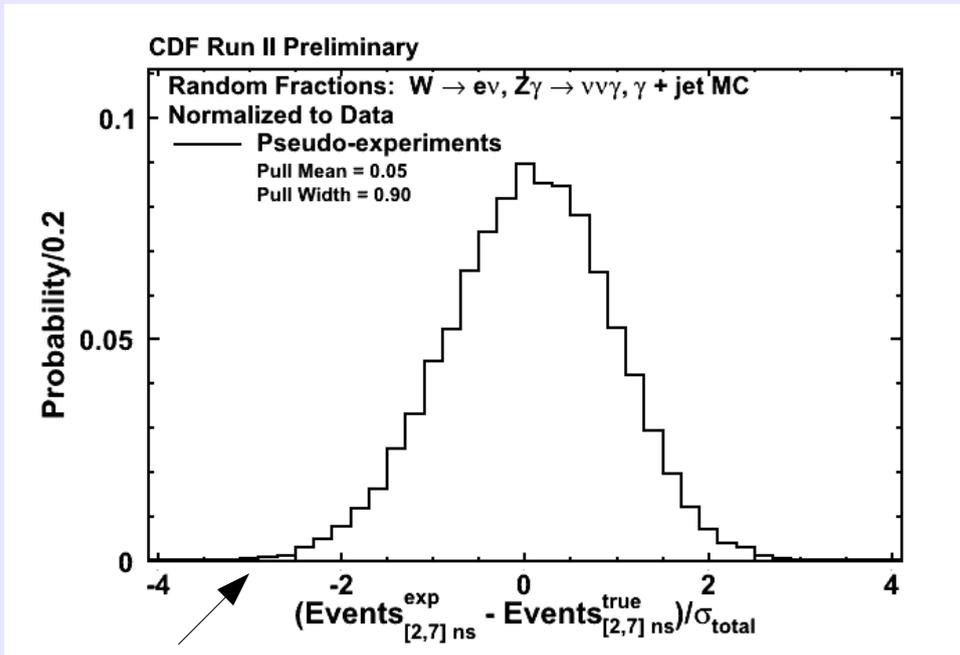
We vary wrong-vertex mean between 0.0 ns and 0.8 ns to see how well the fitter responds.



The quality of the estimation of number of events in the signal region is largely not affected by the particular wrong vertex mean chosen.

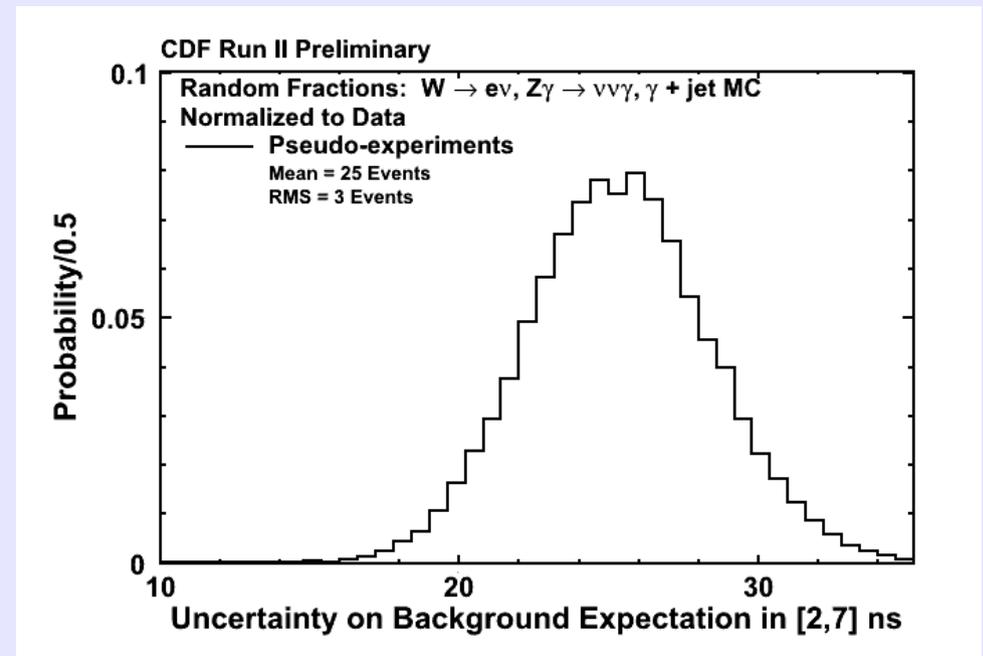
How well do we do when we combine fully simulated MC samples?

We take $Z\gamma$, $W \rightarrow e\nu$, and γ +jet MC in random fractions.



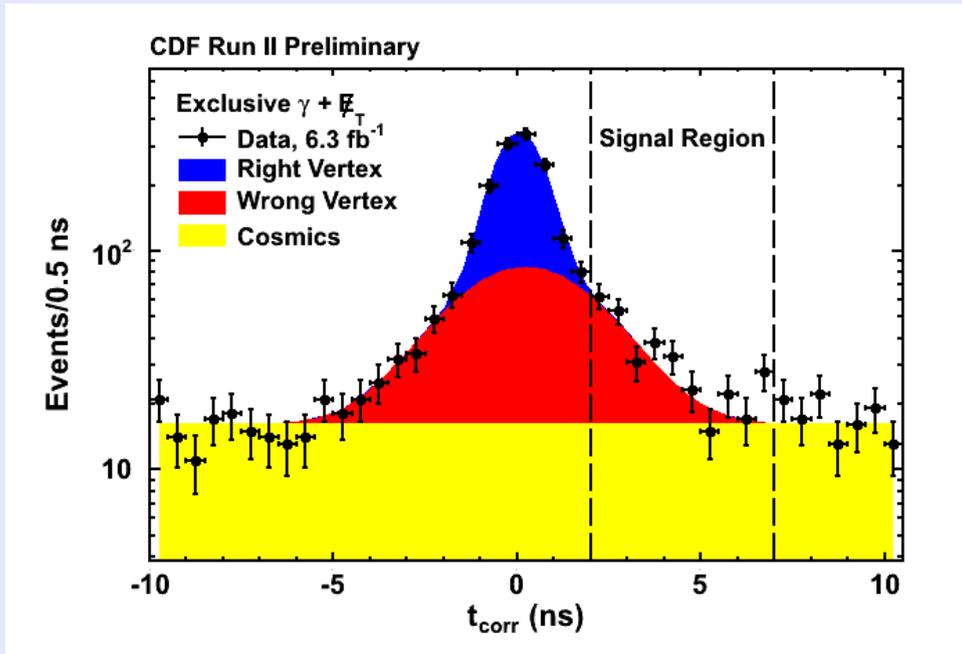
Pull distribution: largely unbiased and the errors well estimated.

Double Gaussian approximation is very successful, even under worse case combinations.



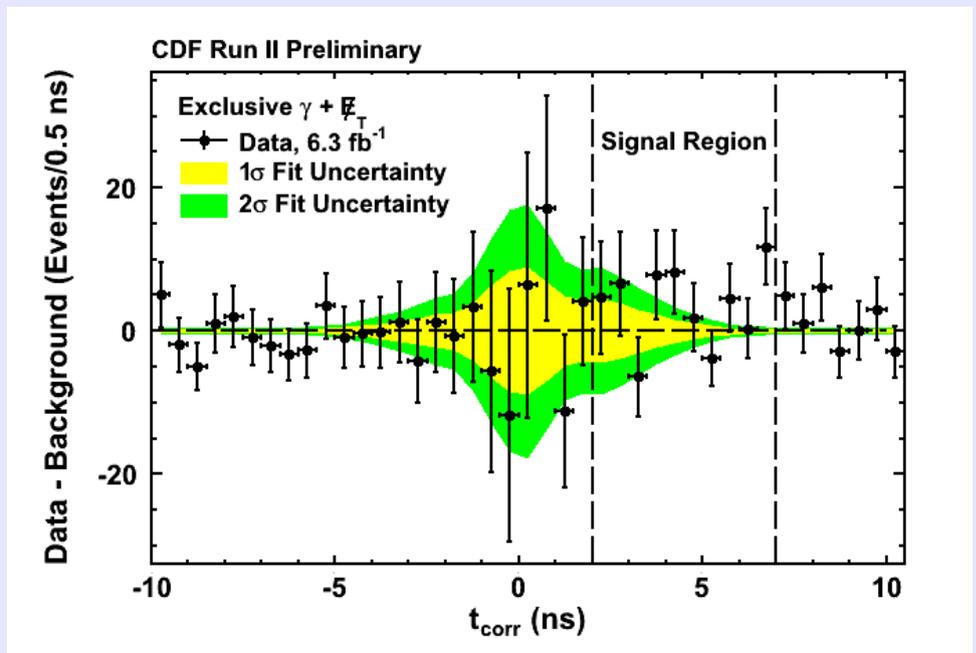
Fit uncertainty ~ 25 counts.

Results



$N(\text{SR})$ expected = 286 ± 24
 $N(\text{SR})$ observed = 322

The counting experiment significance is only 1.2σ , but preliminary results suggest that a shape significance could be much larger



Conclusions

- First attempt at understanding this final state
- Uncovered previously unknown timing biases
- Created new requirements to minimize those biases in an efficient way for signal
- Developed a data driven method to estimate background contributions
- Found a modest but interesting excess → if real, could be the first observation of the Higgs in a SUSY mode

Backups

Event Reduction Table for 6.3 fb^{-1}

Cut	# of Events
Preselect a sample with a Photon w/ $E_T > 45 \text{ GeV}$ & $\text{MET} > 45 \text{ GeV}$	38,291
Reject Beam Halo Events	36,764
Reject Cosmic Events	24,462
Track Veto	16,831
Jet Veto	12,708
Large $ Z $ Vertex Veto	11,702
$e \rightarrow \gamma_{\text{fake}}$ Rejection	10,363
Good Vertex Events/No Vertex Events	5,421/4,942

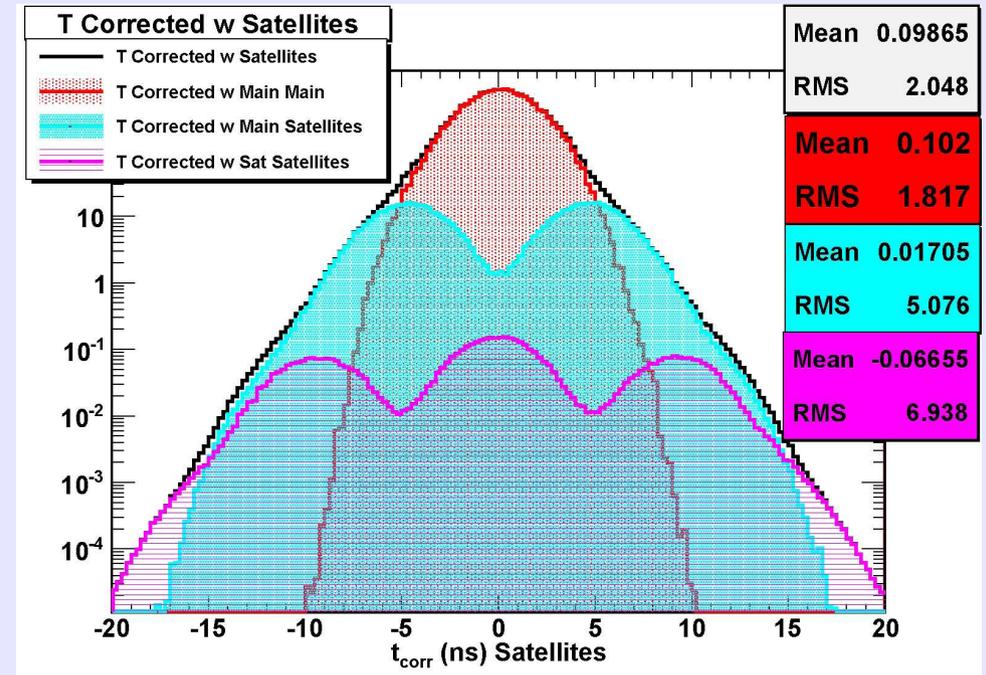
Overview of the Delayed Photon Analysis: Satellite Bunches

Satellite bunches occur 18.8 ns before and after the primary bunches

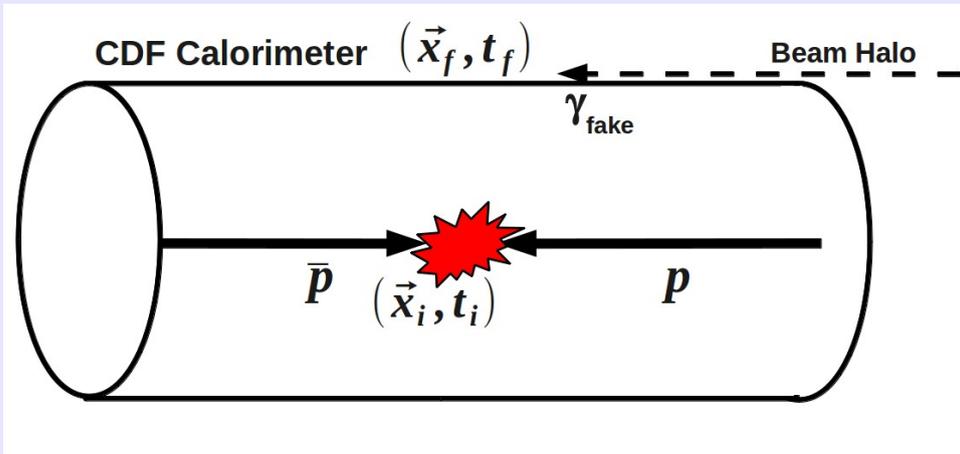
Satellite bunches contain ~1% as many particles as the main bunches do

Satellite-satellite and satellite-main collisions contribute heavily suppressed peaks to the corrected time distribution

These contributions are negligible in this analysis



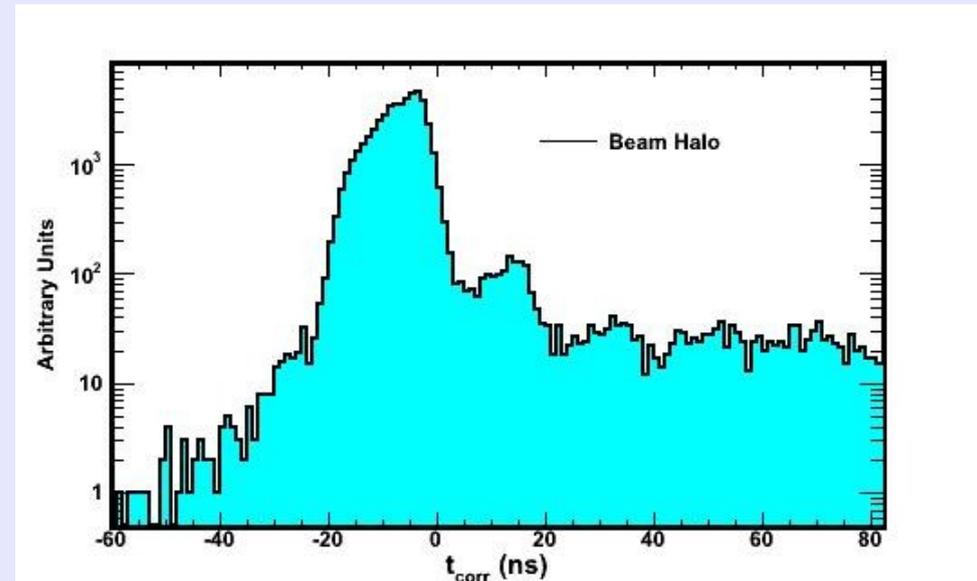
Overview of the Delayed Photon Analysis: Beam Halo



Beam halo particles are typically muons produced beam interactions upstream of the detector

These particles travel parallel to the beam. If they interact in the calorimeter, they predominantly appear as photons arriving earlier than expected.

Our cuts are efficient at removing beam halo

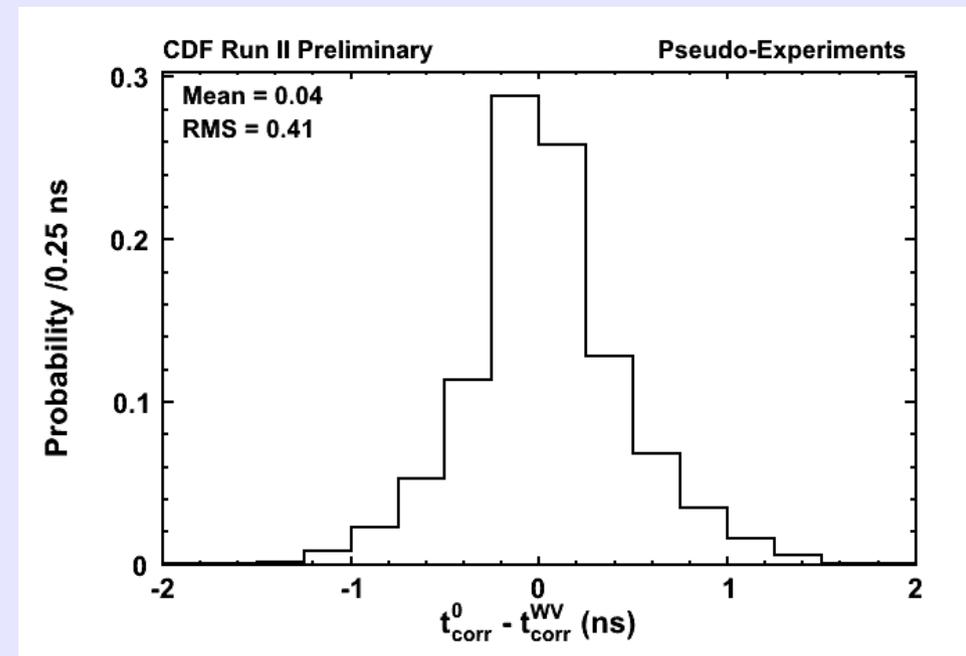


No-Vertex Time and Wrong-Vertex Time Toy MC

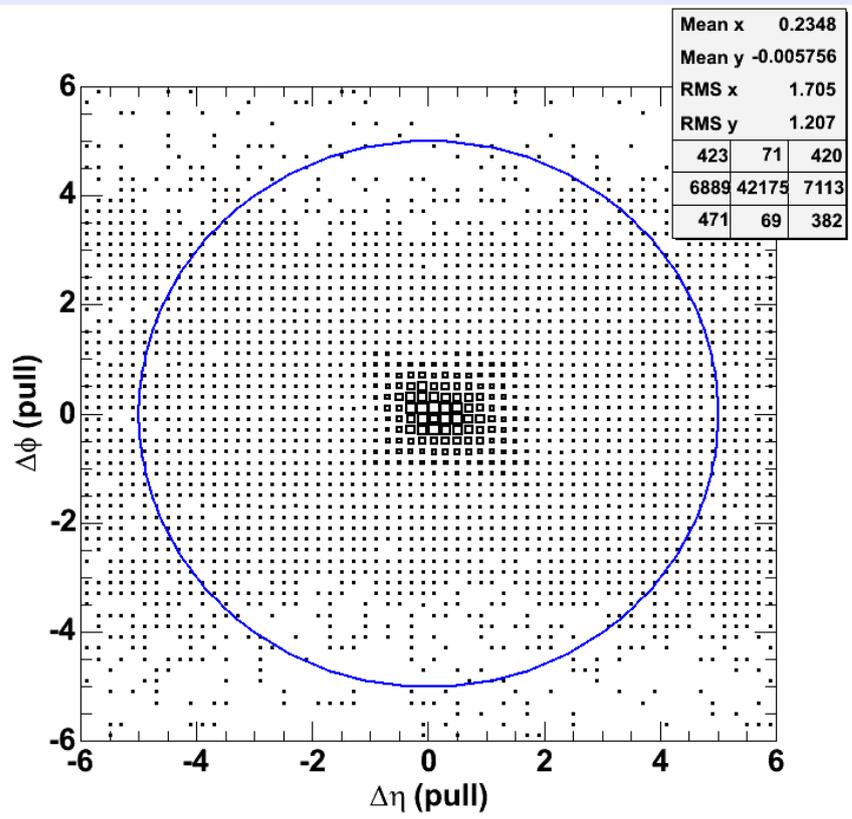
Consider pseudo-experiments where vertices are generated according to the Z and T profiles of the beam spot (Z RMS \sim 28 cm, T RMS \sim 1.28 ns).

Assume spherically symmetric production to determine CES Z.

Shows that if the process dependent geometric time of flight difference is the same for no-vertex and wrong-vertex events, the means of the two distributions will be very close.

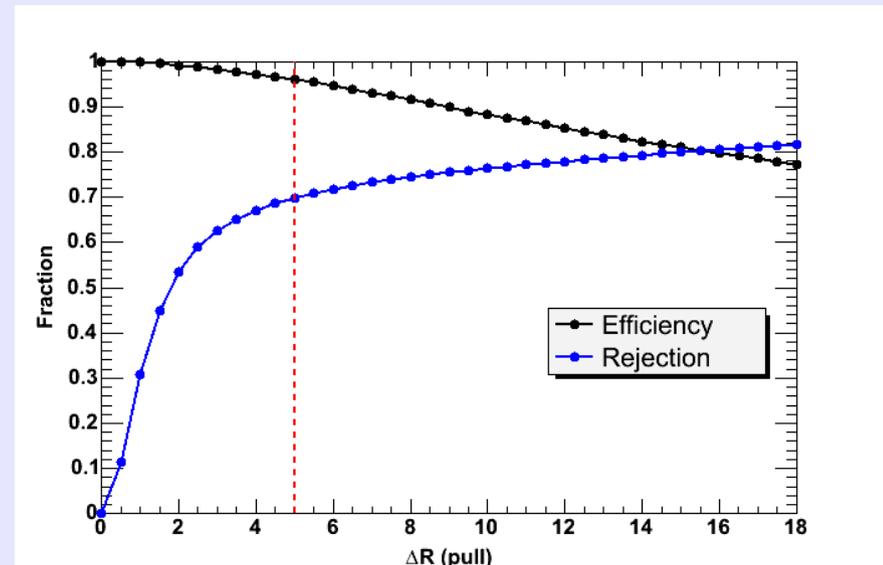


$\Delta R(\text{pull})$



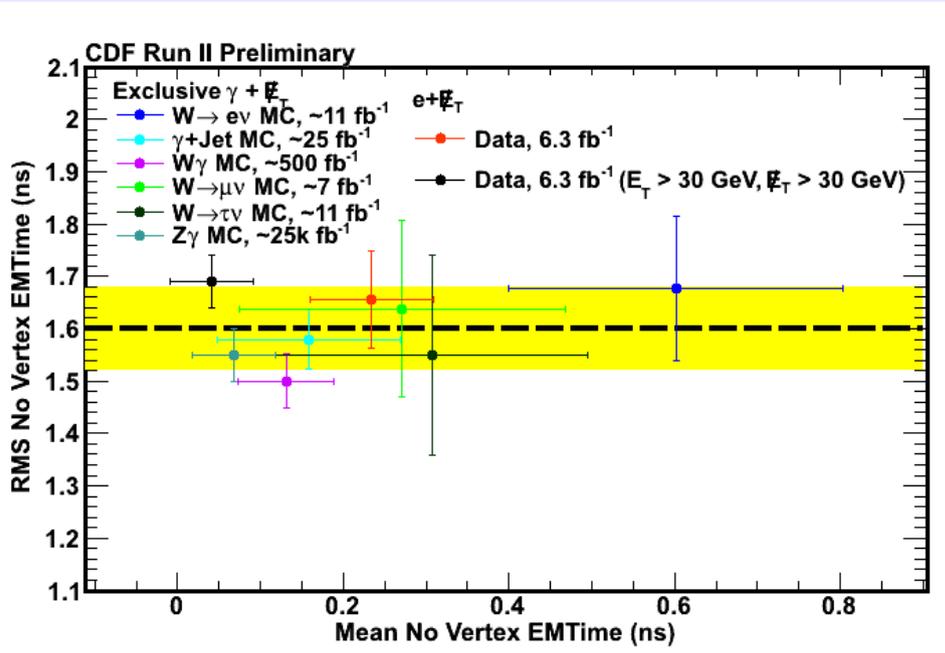
-Find the track with Φ_0 and η closest to the reconstructed photon.

-Standardize the variables to account for worse resolution in Φ_0 due to the “kink” in the track from the hard interaction.



Vetoing reconstructed photons with a track with $\Delta R(\text{pull}) < 5$ removes 73% of fake photons while accepting 95% of real photons.

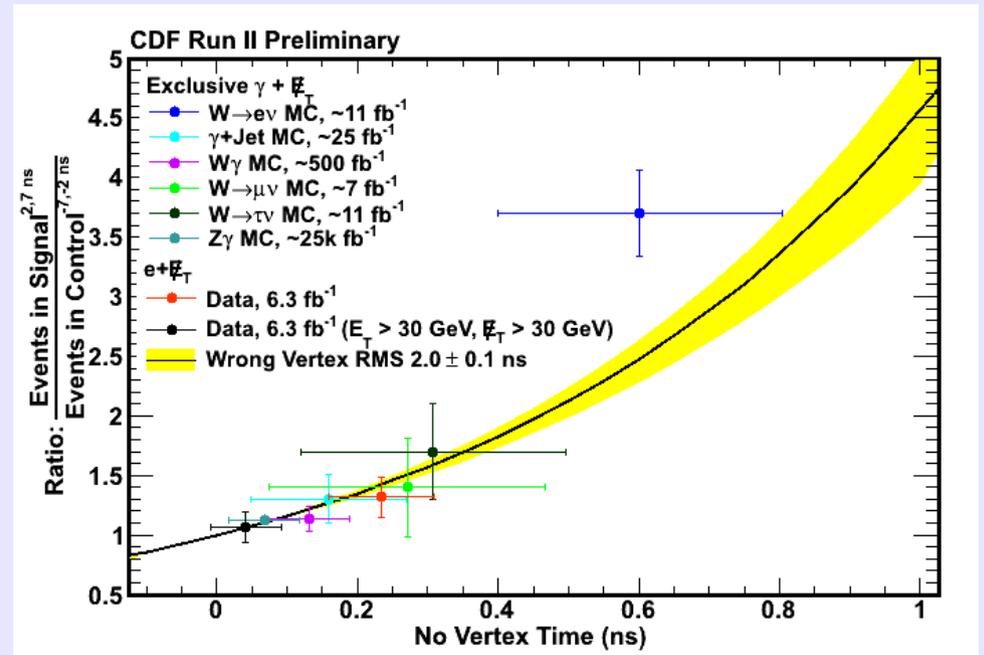
Predicting $N(\text{SR})/N(\text{CR})$ From No Vertex Mean



We isolate no vertex events in Monte Carlo and electron data and fit to find the no vertex mean.

The RMS of the no-vertex distribution does not depend on the mean of the distribution.

$N(\text{SR})/N(\text{CR})$ follows the prediction from the no-vertex mean as well as for the wrong-vertex mean \rightarrow we can use the no-vertex mean as proxy for the wrong-vertex mean.



Combined Likelihood Function

$$-\ln L = \sum_i^{Nbins(GV)} \nu_i^{GV} - n_i^{GV} \ln \nu_i^{GV} + \sum_j^{Nbins(NV)} \nu_j^{NV} - n_j^{NV} \ln \nu_j^{NV} + \sum_k^{Nconstraints} \frac{(\theta_k - \theta_k^0)^2}{2\sigma_k^2}$$

Good vertex portion includes bins between (-7,2) ns and (20,80) ns

No vertex portion includes bins between (-3.5, 3.5) ns and (20,80) ns

ν is the number of expected events in a bin

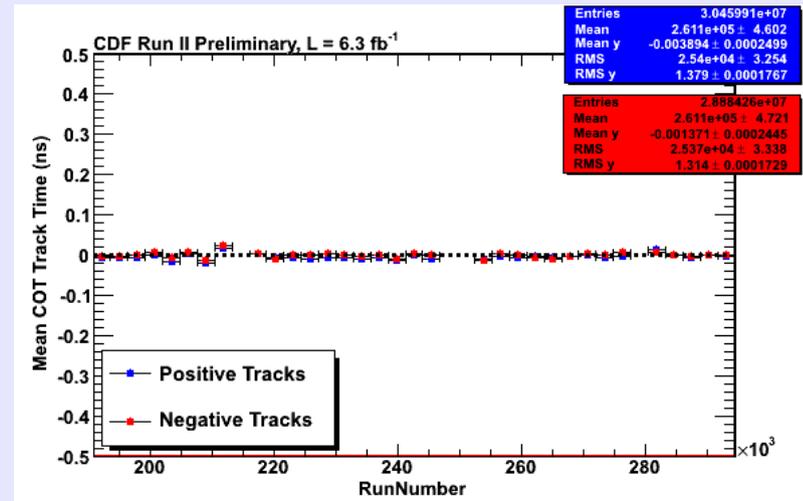
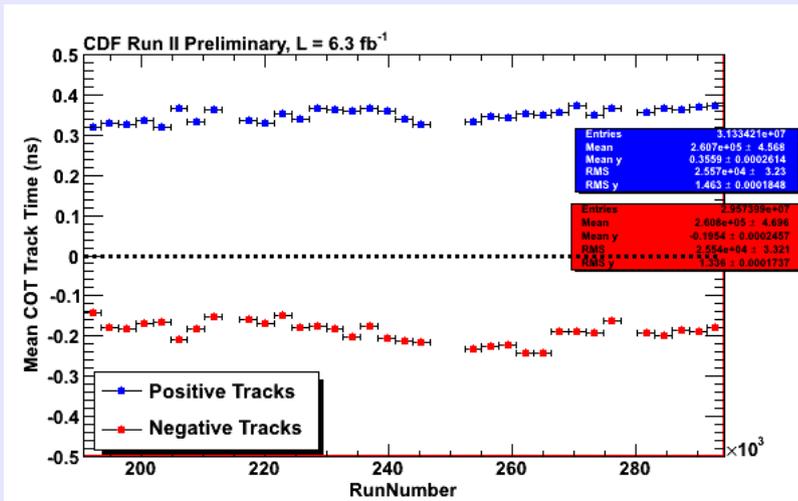
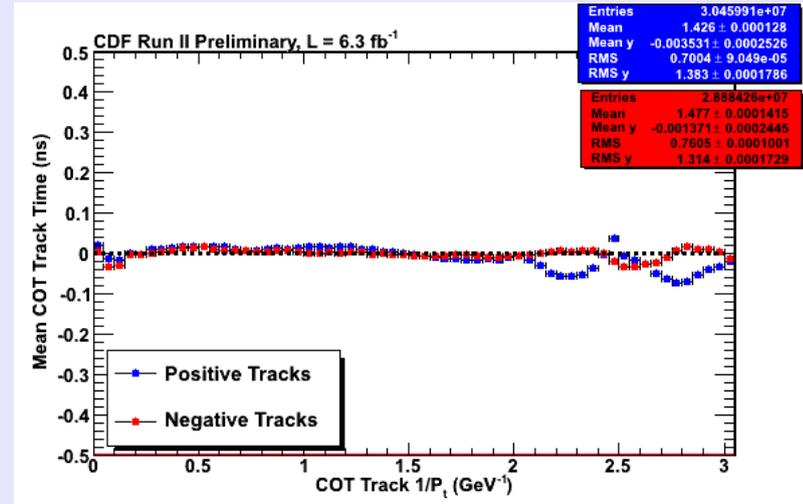
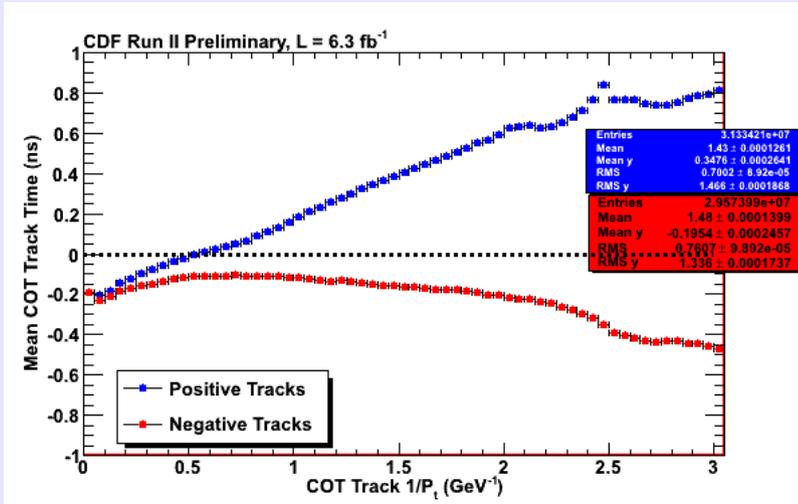
n is the number of observed events in a bin

θ_k is the parameter being constrained

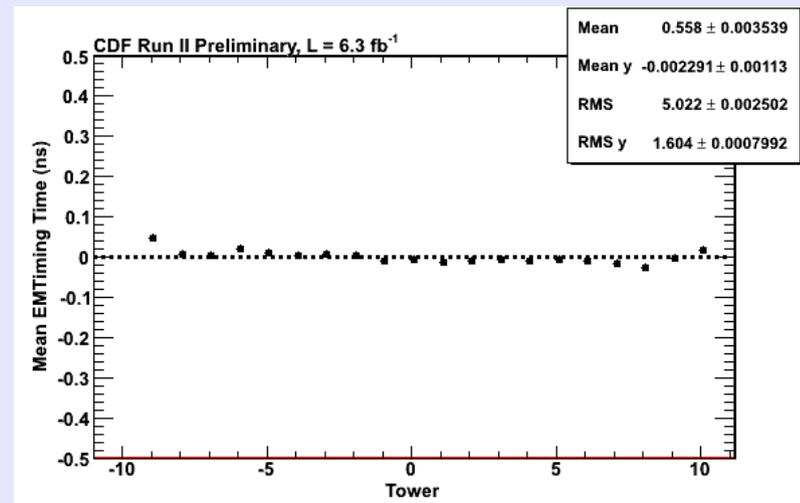
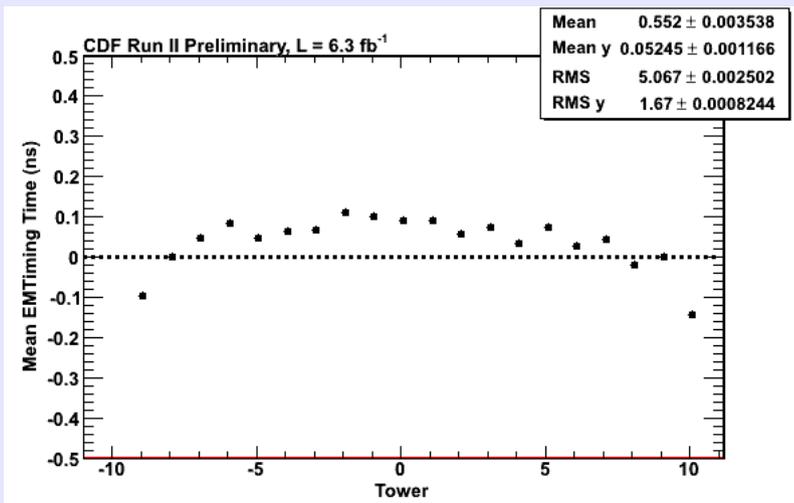
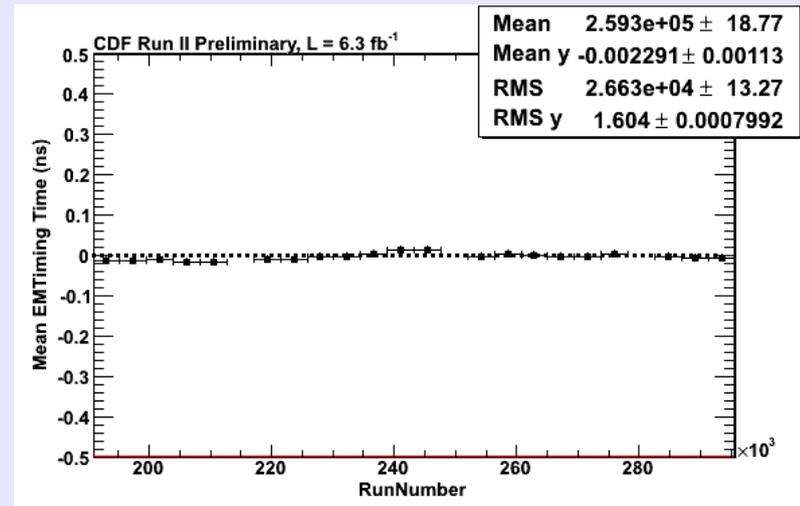
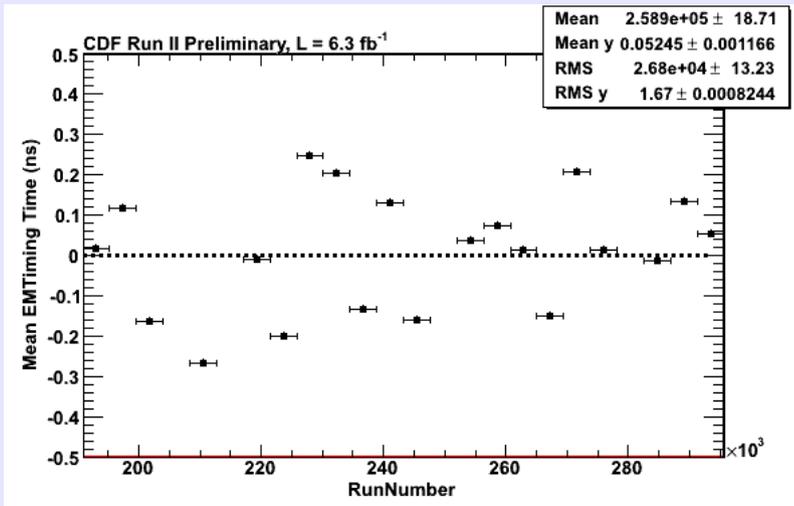
θ_k^0 is the nominal value of the constrained parameter

σ_k is the systematic uncertainty on θ_k

COT Track t_0 Corrections



EMTiming Corrections



No Vertex Distribution

If no good vertex reconstructed, we can still construct the raw time variable: the corrected time, around a vertex with $Z = 0$ and $T = 0$.

The raw time distribution is Gaussian with RMS ~ 1.6 ns.

We will show that the mean of the no vertex distribution is always close to that of the wrong vertex distribution.

